

29. The real space RG calculation for the Ising model in  $d = 2$  on a triangular lattice in a nonzero magnetic field leads to the approximate result (see exercise 25)

$$K' = \frac{K}{2} \left[ \phi\left(K + \frac{1}{2}h\right) + \phi\left(K - \frac{1}{2}h\right) \right]^2$$

$$h' = h + \frac{1}{2} \log \frac{1 + 3e^{-4K-2h}}{1 + 3e^{-4K+2h}} + 3K \left[ \phi\left(K + \frac{1}{2}h\right)^2 - \phi\left(K - \frac{1}{2}h\right)^2 \right]$$

where  $\phi(x) = \frac{1 + e^{-4x}}{1 + 3e^{-4x}}$ . Find the 5 fixed points of the renormalization group transformations, linearize around them (study one after the other) and discuss the nature of the fixed points. Perform a numerical simulation.

30. We study now the mapping of quantum Ising models to classical Ising models. We consider first the case of the single spin quantum Ising model

$$\hat{H} = \Gamma \sigma^x$$

and demonstrate that its partition function  $Z = \text{Tr} e^{-\beta \hat{H}}$  is represented as the partition function of the classical 1d Ising chain.

In the spirit of the Feynman path integral formulation write  $e^{-\beta \hat{H}} = e^{-\frac{\beta}{N} \hat{H}} e^{-\frac{\beta}{N} \hat{H}} \dots e^{-\frac{\beta}{N} \hat{H}}$  and insert complete sets of  $|S_z\rangle$  eigenstates of  $\sigma_z^z$ . Insert additionally also complete sets of  $|S_x\rangle$  eigenstates of  $\sigma_x^x$ , noting that  $\langle S_x | S_z \rangle = \frac{1}{\sqrt{2}} e^{i\pi \frac{1-S_x}{2} \frac{1-S_z}{2}}$ . Verify that  $\langle S'_z | e^{\frac{\beta}{N} \Gamma \sigma^x} | S_z \rangle = c e^{J_\tau S'_z S_z}$ . Here  $J_\tau = \text{artanh}(e^{-2\delta\tau\Gamma})$ ,  $\delta\tau = \frac{\beta}{N}$  and  $c$  is an unimportant constant. Demonstrate the equivalence of the single spin quantum Ising model to the Ising chain.

31. Generalize the above arguments to the quantum Ising chain

$$\hat{H} = -J \sum_i \sigma_i^z \sigma_{i+1}^z - \Gamma \sum_i \sigma_i^x$$

showing that it is mapped to the 2d classical Ising model of size  $\beta = \frac{1}{T}$  in the imaginary time direction. In the limit  $T \rightarrow 0$  the 2d classical Ising model has infinite spatial and temporal directions. The coupling constants in the classical model are anisotropic and depend on the discretization step  $\delta\tau = \frac{\beta}{N}$

$$J_{\text{spat}} = J \delta\tau, \quad J_\tau = \text{arctanh}(e^{-2\delta\tau\Gamma})$$