

## 16. Two-point function of an interacting theory

### Heisenberg picture

basis of eigenstates of  $P^\mu$

$|\Omega\rangle$  vacuum of the interacting theory  
(to be distinguished from the vacuum of the free theory  $|0\rangle$ )

$$P^\mu |\Omega\rangle = 0$$

$$P^\mu |p_1, \dots, p_n \text{ in out}\rangle = (p_1^\mu + \dots + p_n^\mu) |p_1, \dots, p_n \text{ in out}\rangle$$

$$p_i^0 = \sqrt{\vec{p}_i^2 + m_{ph}^2}$$

discussion for scalar theory (extension to higher spin straightforward, but more indices!)

↙ Heisenberg operators!

$$\langle \Omega | T \phi(x) \phi(0) | \Omega \rangle =$$

$$= \theta(x^0) \langle \Omega | \phi(x) \phi(0) | \Omega \rangle$$

$$+ \theta(-x^0) \langle \Omega | \phi(0) \phi(x) | \Omega \rangle$$

$$= \Theta(x^0) \sum_{\alpha} \langle \Omega | \phi(x) | \alpha \rangle \langle \alpha | \phi(0) | \Omega \rangle$$

← insert complete ONS of momentum eigenstates

$$+ \Theta(-x^0) \sum_{\alpha} \langle \Omega | \phi(0) | \alpha \rangle \langle \alpha | \phi(x) | \Omega \rangle$$

$$= \Theta(x^0) \sum_{\alpha} \langle \Omega | e^{+iP_x} \phi(0) e^{-iP_x} | \alpha \rangle \langle \alpha | \phi(0) | \Omega \rangle$$

$$+ \Theta(-x^0) \sum_{\alpha} \langle \Omega | \phi(0) | \alpha \rangle \langle \alpha | e^{+iP_x} \phi(0) e^{-iP_x} | \Omega \rangle$$

$$= \Theta(x^0) \sum_{\alpha} e^{-ip_{\alpha} x} |\langle \Omega | \phi(0) | \alpha \rangle|^2$$

$$+ \Theta(-x^0) \sum_{\alpha} e^{ip_{\alpha} x} |\langle \Omega | \phi(0) | \alpha \rangle|^2$$

$$\sum_{\alpha} |\alpha\rangle \langle \alpha| = |\Omega\rangle \langle \Omega| + \int d\mu(p) |p\rangle \langle p|$$

$$+ \sum_{n=2}^{\infty} \frac{1}{n!} \int d\mu(p_1) \dots d\mu(p_n) |p_1, \dots, p_n^{in}\rangle \langle p_1, \dots, p_n^{in}|$$

assume:  $\langle \Omega | \phi(0) | \Omega \rangle = 0$  (in theories

with  $\langle \Omega | \phi(0) | \Omega \rangle \neq 0$  take

$\phi(x) - \langle \Omega | \phi(0) | \Omega \rangle$  instead of  $\phi(x)$ )

$$\Rightarrow \langle \Omega | T \phi(x) \phi(0) | \Omega \rangle$$

$$= \Theta(x^0) \int d\mu(p) e^{-ipx} |\langle \Omega | \phi(0) | p \rangle|^2$$

$$+ \Theta(-x^0) \int d\mu(p) e^{+ipx} |\langle \Omega | \phi(0) | p \rangle|^2$$

+ contributions from intermediate states  
with  $n \geq 2$

$\langle \Omega | \phi(0) | p \rangle$  is independent of  $p$  because  
of Lorentz invariance:

$|p\rangle \equiv |\sqrt{m_{ph}^2 + \vec{p}^2}, \vec{p}\rangle$  can be obtained by  
a Lorentz transformation acting on  $|m_{ph}, \vec{0}\rangle$ :

$$|p\rangle = |\sqrt{m_{ph}^2 + \vec{p}^2}, \vec{p}\rangle = U(L) |m_{ph}, \vec{0}\rangle$$

where  $L(m_{ph}, \vec{0}) = (\sqrt{m_{ph}^2 + \vec{p}^2}, \vec{p})$

$$\Rightarrow \langle \Omega | \phi(0) | p \rangle = \langle \Omega | \phi(0) U(L) | m_{ph}, \vec{0} \rangle$$

$$= \langle \Omega | U(L)^{-1} \phi(0) U(L) | m_{ph}, \vec{0} \rangle$$

$$\uparrow$$

$$U(L) | \Omega \rangle = | \Omega \rangle$$

$$= \langle \Omega | \phi(0) | m_{ph}, \vec{0} \rangle \quad \text{independent of } p$$

$$\uparrow$$

$\phi$  is a scalar field:  $U(L)^{-1} \phi(x) U(L) = \phi(L^{-1}x)$

$$Z := |\langle \Omega | \phi(0) | p \rangle|^2$$

with a suitable definition of the phase of  $|p\rangle$ :

$$\sqrt{Z} = \langle \Omega | \phi(0) | p \rangle$$

$$\Rightarrow \langle \Omega | T \phi(x) \phi(0) | \Omega \rangle =$$

$$= Z \left\{ \underbrace{\theta(x^0) \int d\mu(p) e^{-ipx} + \theta(-x^0) \int d\mu(p) e^{ipx}}_{\frac{1}{i} \Delta(x; m_{ph})} \right\} + \dots$$

$$= \frac{Z}{i} \int \frac{d^4k}{(2\pi)^4} \frac{e^{-ikx}}{m_{ph}^2 - k^2 - i\varepsilon} + \dots$$

pole of the two-point function in momentum space determines the physical mass  $m_{ph}$

fermions:

$$\langle \Omega | T \psi(x) \bar{\psi}(0) | \Omega \rangle = \frac{1}{i} S(x; m_{ph}) + \dots$$

↑ ↑  
Heisenberg operators

$$= \frac{1}{i} \int \frac{d^4 k}{(2\pi)^4} \frac{e^{-ikx}}{m_{ph} - \not{k} - i\varepsilon} + \dots$$

⏟  
continuum contributions

$$\langle \Omega | \psi(0) | p, s \rangle = \sqrt{Z} u(p, s; m_{ph})$$

$$\langle p, s | \bar{\psi}(0) | \Omega \rangle = \sqrt{Z} \bar{u}(p, s; m_{ph})$$

$$\langle \bar{p}, s | \psi(0) | \Omega \rangle = \sqrt{Z} v(p, s; m_{ph})$$

↑  
anti-particle  
state

$$\langle \Omega | \bar{\psi}(0) | \bar{p}, s \rangle = \sqrt{Z} \bar{v}(p, s; m_{ph})$$