Particle Physics after the Higgs Discovery: Searching for Physics beyond the Standard Model



Thomas Mannel, University of Siegen Particle Physics after the Higgs Discovery: Searching for ...

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Contents







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Introduction

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Constructing the Standard Theory of Particle Physics:

Observed Matter: Quarks and Leptons:

U C up charm top Quarks S down strange bottom v_{e} V_{μ} Elektron-Myon-Tau-Leptonen Neutrino Neutrino Neutrino e μ _{Myon} τ Elektron Tau

Materie (Fermionen)

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Add in the oberved interaction (except gravity):



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Image: A matrix

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Theoretical Problems before the Higgs Discovery:

- The weak bosons W[±] and Z⁰ have to be very massive: Interactions of longitudinal modes generate problems Model without (something like) the Higgs becomes invalid at the TeV scale
- Quark and Leptons have masses Incompatible with a consistent description of parity violation
- This theory fails at high energy scales of about $\mathcal{O}(1)$ TeV

Rationale for LHC (and btw. SSC): something has to happen at the TeV scale!

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The Higgs Discovery



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- The Higgs has specific couplings
- The Higgs couplings generate all masses of quarks and leptons as well as of the massive gauge bosons
- ... with a possible exception of the neutrinos ...

... this has generated a new problem:

This model (with the Higgs particle) is extremely successful!



- It explains all particle physics phenomenology up to energy scales of O(100) GeV
- Various observables have been tested even at the quantum level (such as the Lamb shift, *g* factor of the muon, ...)

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From the theory point of view:

• Including the Higgs this becomes a renormalizable theory: This model can be valid up to extremely high scales $\mu \sim 10^{12}$ GeV!



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(Too) many open questions:

- The model has 27 free parameters,
 22 of these are due to the fact that there are three families
- Many parameters turn out to be "hierarchical" (instead of being $\mathcal{O}(1)$)
 - $m_e = 511 \text{ keV}$ compared to $m_{\tau} = 1.777 \text{ GeV}$
 - $m_u = \text{few MeV}$ compared to $m_{top} = 175 \text{ GeV}$
 - All masses turn (except of the top-quark mass) out to be small compared to the electroweak scale
- Why do we have three families?
- Violation of the CP Symmetry is very small! Despite of the possibility to have large CPV from the strong sector: Strong CP problem
- What are "Dark Matter" and "Dark Energy"?

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A theoretical problem, or: Why many people believe/believed in TeV scale "new physics"?

- Theoretical argument: Related to the mass of the Higgs
- Stabilization of the electroweak scale:



• Quadratic Dependence on the cut-off

$$\Delta m_{H}^2 = -rac{\lambda_f^2}{8\pi^2}\Lambda_{
m UV}^2$$

 $\bullet\,$ Drives the Higgs mass up to the UV cut off $\Lambda_{UV}\sim \textit{M}_{PL}$

 Stabilization at the TeV scale: Needs to introduce something at the TeV scale (e.g. SuperSymmetry):



• Only logarithmic divergence

$$\Delta m_{H}^{2} = m_{\rm soft}^{2} \frac{\lambda}{16\pi^{2}} \ln\left(\frac{\Lambda_{\rm UV}}{m_{\rm soft}}\right)$$

• $m_{\rm soft} \sim O({\rm TeV})$: Scale of "new physics"

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• How strong are these arguments?

• Could there something be wrong with our understanding of

- electroweak symmetry breaking?
- scale and conformal invariance?
 - (c.f. Lee Wick Model)
- ...
- Maybe it is "just so"

Currently there is no evidence for any effects beyond the SM from particle physics

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Beyond the Standard Theory

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Search Strategies for new Effects



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General Tool: Effective (Quantum) Field Theories

Integrating out heavy degrees of freedom

- ϕ : light fields, Φ : heavy fields with mass Λ
- Generating functional as a functional integral Integration over the heavy degrees of freedom

$$Z[j] = \int [d\phi][d\Phi] \exp\left(\int d^4x \left[\mathcal{L}(\phi, \Phi) + j\phi\right]\right)$$

= $\int [d\phi] \exp\left(\int d^4x \left[\mathcal{L}_{eff}(\phi) + j\phi\right]\right)$ with
 $\exp\left(\int d^4x \mathcal{L}_{eff}(\phi)\right) = \int [d\Phi] \exp\left(\int d^4x \mathcal{L}(\phi, \Phi)\right)$

- For length scales $x \gg 1/\Lambda$: local effective Lagrangian
- Technically: (Operator Product) Expansion in inverse powers of Λ

$$\mathcal{L}_{\mathrm{eff}}(\phi) = \mathcal{L}_{\mathrm{eff}}^{(4)}(\phi) + \frac{1}{\Lambda} \mathcal{L}_{\mathrm{eff}}^{(5)}(\phi) + \frac{1}{\Lambda^2} \mathcal{L}_{\mathrm{eff}}^{(6)}(\phi) + \cdots$$

- \mathcal{L}_{eff} is in general non-renormalizable, but ...
- $\mathcal{L}_{eff}^{(4)}$ is the renormalizable piece
- For a fixed order in $1/\Lambda$: Only a finite number of insertions of $\mathcal{L}_{eff}^{(4)}$ is needed!
- ullet ightarrow can be renormalized
- Renormalizability is not an issue here

Effective field Theory as used by the practitioneer

Effective Field Theories are everywhere:

• Effective (local) Lagrangian

$$\mathcal{L} = \mathcal{L}_{\dim 4}^{SM} + \mathcal{L}_{\dim 5} + \mathcal{L}_{\dim 6} + \cdots$$

• $\mathcal{L}_{dim\,n}$ are suppressed by large mass scales

$$\mathcal{L}_{\dim n} = \frac{1}{\Lambda^{n-4}} \sum_{i} C_n^{(i)} O_n^{(i)}$$

- $O_n^{(i)}$: Operators of dimension *n*,
- $C_n^{(i)}$: (dimensionless) couplings
- Λ scale of "new physics"

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Towers of effective field theories

Depending on the Scale:

- Λ_{BSM} ≥ μ ≥ ν: Standard Model Effective Theory (SMEFT) matching to an unknown theory BSM
- v ≥ µ ≥ m_b: Weak Effective Theory (WET) matching to the SM or SMEFT
- *m_b* ≥ μΛ_{QCD}: Heavy Quark Effective Theory matching to WET
- Λ_{χSB} ≥ μ ≥ 0: Chiral Perturbation Theory matching to ...

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Excursion: It worked in the past: The Top Quark Story

- First indirect hint to a heavy top quark: $B - \overline{B}$ Oscillation of ARGUS (1987)
- The world in 1987 ("PETRA Days"): The top was believed to be at ~ 25 GeV
 ... based on "good theoretical arguments"!!
- ARGUS could not have seen anything with a 25 GeV Top (within SM)
- Clear indication if a very heavy top quark!



• The consequences:

Mass of the top quark must be well above 100 GeV!

- (–) No Toponium
- (-) No Top quark discovery at LEP and SLC
- (-) No "New Physcis $\mathcal{O}(30 \text{ GeV})$ " just around the corner
- (+) CP violation in the B sector may become observable large asymmetries, however, for decays with small branching fractions
- This was actually good for Flavour Physics (*B* mesons and Kaons)

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What can we say today about the "new physics" scale?

Pick some processes that are very sensitive to "new physics"

- = well measurable and well under theoretical control
 - Anomalous magnetic moment of the muon $(g-2)_{\mu}$, looks now SM like, lmits based on EFT arguments have to be re-evaluated
 - Mixing phenomena in the neutral flavoured meson systems: ($B-\overline{B}$ -Mixing, $D-\overline{D}$ -Mixing, $K-\overline{K}$ -Mixing)
 - Processes forbidden at leading order: "Flavour Changing Neutral Currents": $b \rightarrow s$ transitions and alike
 - Flavour Violation in the leptonic sector: $\mu \rightarrow e\gamma$ and alike
 - Patterns of the violations of the CP symmetry, Unitarity Triangel fits etc.

From lepton physics:

• Majorana masses for the ν 's are generated by a unique dim-5 operator:

$$\mathcal{L}_{\dim 5} = rac{1}{\Lambda_{\mathrm{LNV}}} \sum_{ij} C_5^{ij} (\bar{L}_j H^c)^c (H^{c,\dagger} L_i)$$

- Generates a mixing matrix for the leptons (PMNS Matrix), analogous to the CKM Matrix
- This term is Lepton Number Violating, related to the scale Λ_{LNV}
- We know that Neutrinos have small masses: Δm_{ij}^2 are measured! We know that there is a family mixing for leptons.
- $\bullet\,$ Use these inputs: $\Lambda_{LNV} \sim 10^{14} \; GeV!$

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Good News:

- Neutrino masses induces non-trivial charged lepton flavour physics
- LFV muon decays: $\mu \rightarrow e\gamma$, $\mu \rightarrow 3e$
- LFV τ decays: $\tau \rightarrow e/\mu\gamma$, $\tau \rightarrow 3\mu$, $\tau \rightarrow e2\mu$...

Bad News:

- If the dim-5 Operator is the only source of LFV (and LNV), these effects are super-small!
- Naive counting

$$\mathcal{A}(\ell o \ell' \gamma) \propto rac{G_{\mathcal{F}}}{16 \pi^2} |V_{ ext{PMNS}}|^2 rac{\Delta m_
u^2}{M_W^2} \quad o \quad ext{Br}(au o \mu \gamma) \leq 10^{-54}$$

Hopefully Λ_{QFV} and Λ_{LFV} is not that high!

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From the mixing in the flavoured neutral meson systems:

- For Quarks there is no contribution to $\mathcal{L}_{dim 5}$
- Look at $\Delta F = 2$ flavour transitions:

$$\begin{split} &O_1^{(6)} = (\bar{s}_L \gamma_\mu d) (\bar{s}_L \gamma^\mu d) & (\text{Kaon Mixing}) \\ &O_2^{(6)} = (\bar{b}_L \gamma_\mu d) (\bar{b}_L \gamma^\mu d) & (B_d \text{ Mixing}) \\ &O_3^{(6)} = (\bar{b}_L \gamma_\mu 2) (\bar{b}_L \gamma^\mu s) & (B_s \text{ Mixing}) \\ &O_4^{(6)} = (\bar{c}_L \gamma_\mu u) (\bar{c}_L \gamma^\mu u) & (D \text{ Mixing}) \end{split}$$

- With generic couplings $C_i \sim \mathcal{O}(1)$:
 - $\Lambda \sim 1000 \text{ TeV}$ from Kaon mixing
 - $\Lambda \sim 1000$ TeV from D mixing
 - $\Lambda \sim 400$ TeV from B_d mixing
 - $\Lambda \sim 70$ TeV from B_s mixing

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Current "Tensions"

There are currently a few "tensions" (Tension = Deviation with a significance of less than 5 σ , so normally nothing to worry about!)

- Anomalous Magnetic Moment Moment of the Muon, very likely disappeared
- $b \rightarrow s \ell \ell$ Transitions:
 - P'_5 : Angular distribution in $B \to K^* \ell \ell$
 - Rates for $B \rightarrow K \mu \mu$ and $B_s \rightarrow \phi \mu \mu$
 - Lepton-Universality Violation: Disappeared!
- R(D) and $R(D^*)$: Rates for $B \to D^{(*)} \ell \bar{\nu}$
- V_{xb}^{incl} vs. V_{xb}^{excl} : $b o q \ell \bar{\nu}$ transitions

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 $b \rightarrow s\ell\ell$ anomalies: Angular distribution in $B \rightarrow K^*\ell\ell$

• $B \rightarrow K^* \ell \ell \rightarrow K \pi \ell \ell$ contains a lot of information

- Angular distributions in the final state
- Set up clever ratios to reduce hadronic uncertainties

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Anomalies in the angular distributions:



However, how well can we compute this?

- There are various hadronic uncertainties
- Needs additional scrutiny within the Standard Model!

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$b \rightarrow s \ell \ell$ anomalies: Rates in $B \rightarrow K \mu \mu$ and $B_s \rightarrow \phi \mu \mu$



Rates at low q^2 seem to be lower than the SM prediction!

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Semileptonic Decays with a τ Lepton

Tension in the exclusive semileptonic ${\it B}
ightarrow {\it D}^{(*)} au ar{
u}$ decays

$$R(D) = \frac{\Gamma(B \to D\tau\bar{\nu})}{\Gamma(B \to D\ell\bar{\nu})} \quad R(D^*) = \frac{\Gamma(B \to D^*\tau\bar{\nu})}{\Gamma(B \to D^*\ell\bar{\nu})}$$



- Theory predictions are quite precise:
- Heavy Quark Symmetry fixes the longitudinal form factor f₀
- in Addition, its contribution is supressed by m²_τ/m²_b

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Determinations of CKM Matrix Elements

 V_{cb} and V_{ub} can be determined from various decay modes:



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How do we continue form here?

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Narrative from 2019 (Start of P3H): Particle Physics at the crossroads

LHC finds New Particles

- Find our what it is!
- How does this become compatible with the precision data?
- Why do we have MFV?
- ... and where does it come from?

LHC finds no New Particles

- Era of indirect searches
- Quark and Lepton Flavor Physics
- Indirect searches at highest energies

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 "Precision Colllider Physics" at LHC

Perspective for a direct detection has shrunk!



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Generic searches at highest energies: SMEFT fits (Biekotter et al., 1812.07587)



There are too many operators at dim 6

- Pick some operators to study
- ... and fit them "one by one"
- Results are compatible with $f_i = 0$

Uncertainties are (still) too large:

• Assuming a scale $\Lambda=10$ TeV (which is still too low for LHC!) we get

 $-1000 \leq f_i \leq 1000$

SMEFT is not (yet?) the proper tool to analyse LHC data!

- No sufficient sensitivity (yet?)
- Too many parameters

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Simplified Models



- "Resolve" the dim 6 operators by introducing exchanged (heavy) particles
- Restricts the number of operators in SMEFT

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• "Reverse Engineering"

Quark-Flavor Physics and WET fits

This works much better as for Collider Physics and SMEFT

$$\mathcal{O}_{7}^{(\prime)} = (\bar{s}\sigma_{\mu\nu}P_{R(L)}b)F^{\mu\nu},$$
$$\mathcal{O}_{9\ell}^{(\prime)} = (\bar{s}\gamma_{\mu}P_{L(R)}b)(\bar{\ell}\gamma^{\mu}\ell),$$
$$\mathcal{H}_{\text{eff}} = -\frac{4G_{F}}{\sqrt{2}}V_{tb}V_{ts}^{*}\sum_{i}C_{i}(\mu)\mathcal{O}_{i}(\mu)$$
$$\mathcal{O}_{10\ell}^{(\prime)} = (\bar{s}\gamma_{\mu}P_{L(R)}b)(\bar{\ell}\gamma^{\mu}\gamma_{5}\ell),$$

• Fit the coefficients $C_i(m_b)$ from the data

compare to the values in the Standard Model

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• *B* anomalies seem to follow a simple picture:

$$C_{9\mu}^{(NP)}=C_{9e}^{(NP)}$$
 and $C_{9}^{(NP)}=-C_{10}^{(NP)}$

- Interpretiations in terms of Leptoquarks and/or Z'
- ... or still Standard Model???

more work is needed ...

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Outlook

Discussion on the European Strategy in Particle Physics is in full swing! Should be oriented along physics questions!

- What is the solution of the strong CP problem?
 - Is there a light axion or other feebly interacting light particles?
- What is Dark Energy / Dark Energy?
 - Has Dark Matter a particle-physics explanation?
 - ... if yes, does it consist of light pr heavy pqarticles?
- Where is the missing CP Violation?
 - Is there PMNS CP violation and/or other CP violation in the lepton sector?
 - Are there other sources of CP violation in the quark sector?

... and maybe more ...

Conclusion

100 Years of Quantum Theory: From $i\hbar \frac{\partial}{\partial t}\psi(t, \vec{x}) = H\psi(t, \vec{x})$ to

$$\begin{aligned} \mathcal{Z} &= -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ &+ i \mathcal{F} \mathcal{D} \mathcal{F} \\ &+ \mathcal{F} \mathcal{D} \mathcal{F} \\ &+ \mathcal{F} \mathcal{D}_{ij} \mathcal{F}_{j} \mathcal{P} + h.c. \\ &+ |\mathcal{D}_{\mu} \mathcal{P}|^{2} - \mathcal{V} (\mathcal{P}) \end{aligned}$$

- Quantum mechanics was a dramatic change of paradigm
- ... triggered by experimental evidences
- and so was the (special) theory of relativty





- The gravitational constant *G* is the conversion factor relating energy-momentum to space-time geometry
- it is still missing in the figure! Does this indicate that after 100 years a new paradigm shift is around the corner?

Does this indicate that - after 100 years - a new paradigm shift is around the corner?

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