Gauge Theory Bootstrap:

Pion amplitudes and low energy parameters

Yifei He

Ecole Normale Supérieure, Paris

Seminar at University of Vienna, 17/12/2024

Based on:

[YH and Kruczenski, Phys. Rev. Lett. 133, 191601, Phys. Rev. D. 110, 096001]

[YH and Kruczenski, arXiv: 2403.10772]

Low energy physics of asymptotically free gauge theory

asymptotically free gauge theory $SU(N_c)$ with N_f massive quarks $m_q \ll \Lambda_{\mathrm{QCD}}$

confinement & chiral symmetry breaking

Low energy physics of asymptotically free gauge theory

asymptotically free gauge theory $SU(N_c)$ with N_f massive quarks $m_q \ll \Lambda_{\rm QCD}$

confinement & chiral symmetry breaking

$$\mathcal{L} = i\sum_j^{N_f} \bar{q}_j \not{D} q_j - \sum_j^{N_f} m_q \bar{q}_j q_j - \frac{1}{4} G_a^{\mu\nu} G^a_{\mu\nu} + \text{gauge fixing} + \text{ghost}$$

gauge theory parameters: N_c N_f m_q $\Lambda_{\rm QCD}$

Low energy physics of asymptotically free gauge theory

asymptotically free gauge theory $SU(N_c)$ with N_f massive quarks $m_q \ll \Lambda_{\text{QCD}}$

confinement & chiral symmetry breaking

$$\mathcal{L} = i\sum_j^{N_f} \bar{q}_j \not{D} q_j - \sum_j^{N_f} m_q \bar{q}_j q_j - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a + \text{gauge fixing} + \text{ghost}$$

gauge theory parameters: N_c N_f m_q Λ_{QCD}

What is the low energy physics?

Physics of Goldstone bosons

chiral symmetry breaking

explicitly — *dynamically*

quark mass $m_q \bar{q}_j q_j$

$$SU(N_f)_L \times SU(N_f)_R \to SU(N_f)_V$$

condensate $\langle \bar{q}_j q_j \rangle$

pseudo-Goldstone bosons dominate the low energy physics

Physics of Goldstone bosons

chiral symmetry breaking

explicitly — *dynamically*

quark mass $m_q \bar{q}_j q_j$ $SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V$ *condensate* $\langle \bar{q}_j q_j \rangle$

pseudo-Goldstone bosons dominate the low energy physics

e.g. $N_f = 2$ pions $\pi_0 = \pi^3$ $\pi_\pm = \frac{1}{\sqrt{2}}(\pi^1 \pm i\pi^2)$

very low energy effective Lagrangian (lowest order):

$$\mathcal{L} = \frac{f_\pi^2}{4} \left\{ \mathrm{Tr}\left(\partial_\mu U \partial^\mu U^\dagger\right) + m_\pi^2 \mathrm{Tr}\left(U + U^\dagger\right) \right\} \qquad U = e^{i\frac{\vec{\tau}\cdot\vec{\pi}}{f_\pi}}$$

$$\mathcal{L}_2^{2\pi} = \frac{1}{2}\partial_\mu \vec{\pi} \cdot \partial^\mu \vec{\pi} - \frac{1}{2} m_\pi^2 \vec{\pi}^2 \qquad \mathcal{L}_2^{4\pi} = \frac{1}{6 f_\pi^2}\Big((\vec{\pi}\cdot\partial_\mu\vec{\pi})^2 - \vec{\pi}^2(\partial_\mu\vec{\pi}\cdot\partial^\mu\vec{\pi}) \Big) + \frac{m_\pi^2}{24 f_\pi^2}(\vec{\pi}^2)^2 \quad \ldots$$

The EFT approach

non-renormalizable, add new terms with unknown coefficients:

e.g.

$$
\begin{aligned}
\mathcal{L}_4 \;=\; & \frac{l_1}{4}\left\{\mathrm{Tr}[D_\mu U(D^\mu U)^\dagger]\right\}^2 + \frac{l_2}{4}\mathrm{Tr}[D_\mu U(D_\nu U)^\dagger]\mathrm{Tr}[D^\mu U(D^\nu U)^\dagger] \\
& + \frac{l_3}{16}[\mathrm{Tr}(\chi U^\dagger + U\chi^\dagger)]^2 + \frac{l_4}{4}\mathrm{Tr}[D_\mu U(D^\mu \chi)^\dagger + D_\mu \chi(D^\mu U)^\dagger] \\
& + l_5\left[\mathrm{Tr}(f^R_{\mu\nu} U f^{\mu\nu}_L U^\dagger) - \frac{1}{2}\mathrm{Tr}(f^L_{\mu\nu} f^{\mu\nu}_L + f^R_{\mu\nu} f^{\mu\nu}_R)\right] \\
& + i\frac{l_6}{2}\mathrm{Tr}[f^R_{\mu\nu} D^\mu U(D^\nu U)^\dagger + f^L_{\mu\nu}(D^\mu U)^\dagger D^\nu U] \\
& - \frac{l_7}{16}[\mathrm{Tr}(\chi U^\dagger - U\chi^\dagger)]^2
\end{aligned}
$$

The EFT approach

non-renormalizable, add new terms with unknown coefficients:

e.g.

$$\begin{aligned}\mathcal{L}_4 &= \frac{l_1}{4}\left\{\mathrm{Tr}[D_\mu U(D^\mu U)^\dagger]\right\}^2 + \frac{l_2}{4}\mathrm{Tr}[D_\mu U(D_\nu U)^\dagger]\mathrm{Tr}[D^\mu U(D^\nu U)^\dagger] \\ &+\frac{l_3}{16}[\mathrm{Tr}(\chi U^\dagger + U\chi^\dagger)]^2 + \frac{l_4}{4}\mathrm{Tr}[D_\mu U(D^\mu \chi)^\dagger + D_\mu \chi (D^\mu U)^\dagger] \\ &+l_5\left[\mathrm{Tr}(f^R_{\mu\nu} U f^{\mu\nu}_L U^\dagger) - \frac{1}{2}\mathrm{Tr}(f^L_{\mu\nu} f^{\mu\nu}_L + f^R_{\mu\nu} f^{\mu\nu}_R)\right] \\ &+i\frac{l_6}{2}\mathrm{Tr}[f^R_{\mu\nu} D^\mu U (D^\nu U)^\dagger + f^L_{\mu\nu}(D^\mu U)^\dagger D^\nu U] \\ &-\frac{l_7}{16}[\mathrm{Tr}(\chi U^\dagger - U\chi^\dagger)]^2\end{aligned}$$

χPT: unknown coefficients determined from fitting with experimental data

The EFT approach

non-renormalizable, add new terms with unknown coefficients:

e.g.

$$\begin{aligned}\mathcal{L}_4 = & \frac{l_1}{4}\{\mathrm{Tr}[D_\mu U(D^\mu U)^\dagger]\}^2 + \frac{l_2}{4}\mathrm{Tr}[D_\mu U(D_\nu U)^\dagger]\mathrm{Tr}[D^\mu U(D^\nu U)^\dagger] \\ & + \frac{l_3}{16}[\mathrm{Tr}(\chi U^\dagger + U\chi^\dagger)]^2 + \frac{l_4}{4}\mathrm{Tr}[D_\mu U(D^\mu \chi)^\dagger + D_\mu \chi (D^\mu U)^\dagger] \\ & + l_5\left[\mathrm{Tr}(f^R_{\mu\nu} U f^{\mu\nu}_L U^\dagger) - \frac{1}{2}\mathrm{Tr}(f^L_{\mu\nu} f^{\mu\nu}_L + f^R_{\mu\nu} f^{\mu\nu}_R)\right] \\ & + i\frac{l_6}{2}\mathrm{Tr}[f^R_{\mu\nu} D^\mu U (D^\nu U)^\dagger + f^L_{\mu\nu}(D^\mu U)^\dagger D^\nu U] \\ & - \frac{l_7}{16}[\mathrm{Tr}(\chi U^\dagger - U\chi^\dagger)]^2\end{aligned}$$

χPT: unknown coefficients determined from fitting with experimental data

in principle should be computed from UV gauge theory

Strongly coupled physics

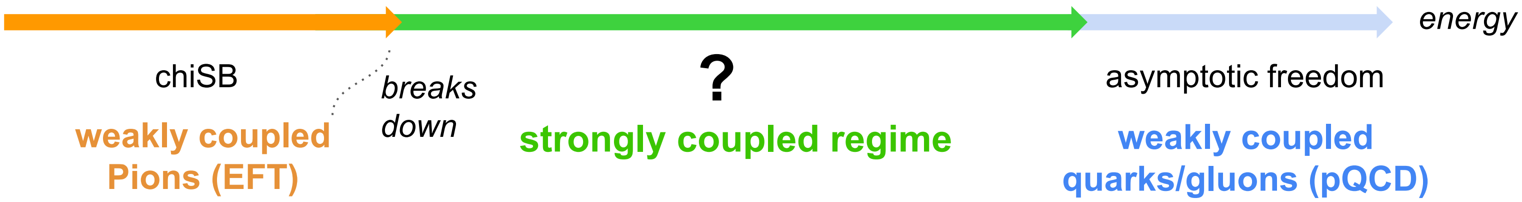
chiSB

weakly coupled
Pions (EFT)

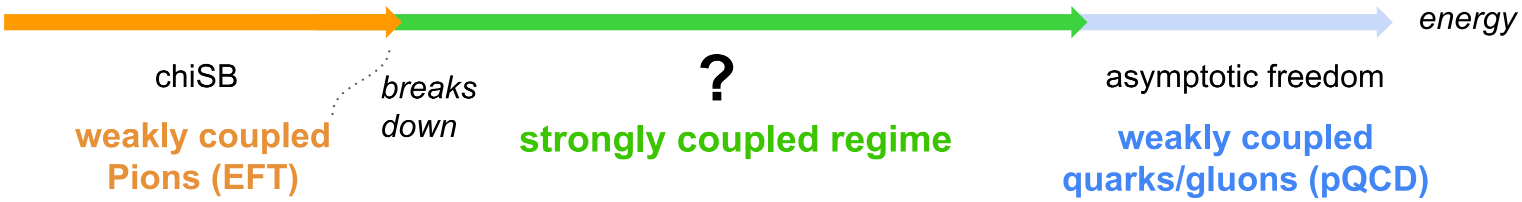
asymptotic freedom

weakly coupled
quarks/gluons (pQCD)

Strongly coupled physics



Strongly coupled physics → Gauge Theory Bootstrap



compute the strongly coupled hadron dynamics:
amplitudes, form factors, correlation functions, spectrum/couplings

Gauge Theory Bootstrap

Strongly coupled physics → Gauge Theory Bootstrap

energy

chiSB

weakly coupled Pions (EFT)

breaks down

?

strongly coupled regime

asymptotic freedom

weakly coupled quarks/gluons (pQCD)

compute the strongly coupled hadron dynamics:
amplitudes, form factors, correlation functions, spectrum/couplings

Gauge Theory Bootstrap

to remove in the future

rules of the game:

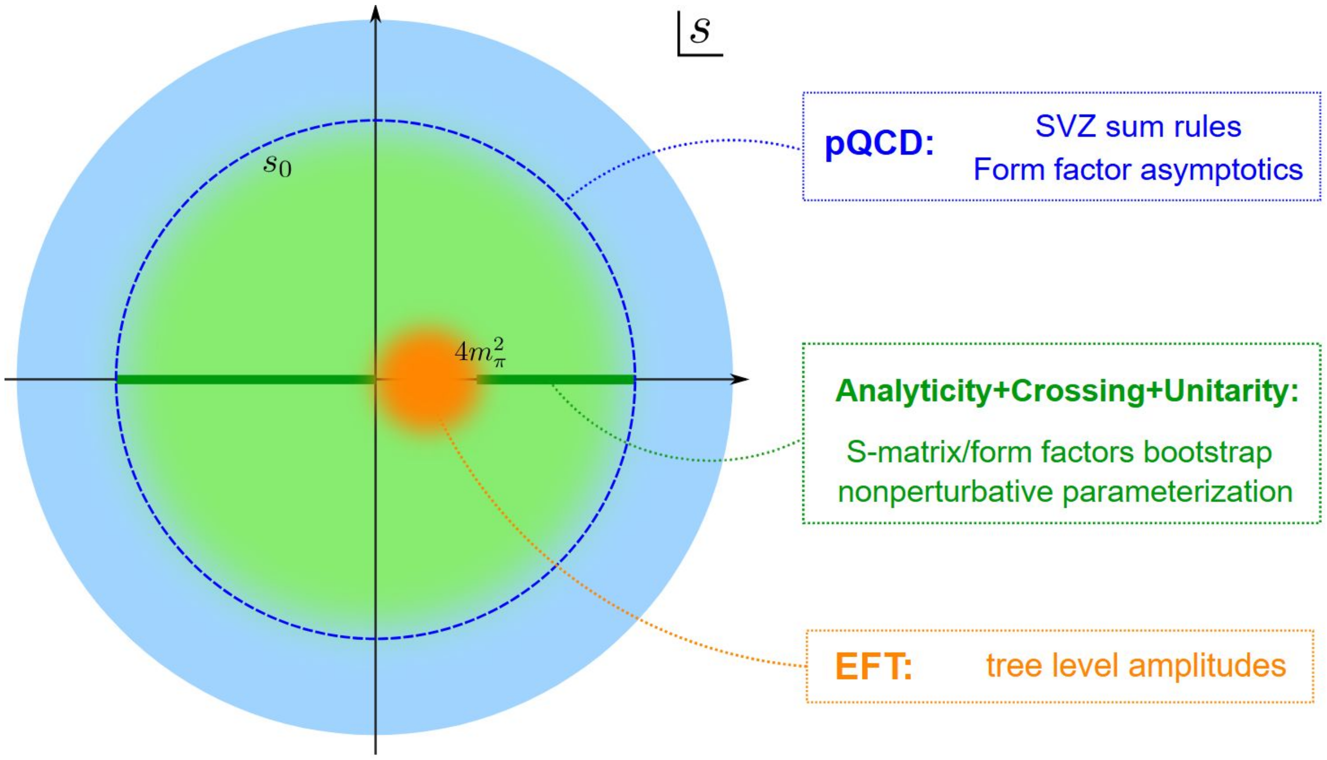
assume — *chiral symmetry breaking & confinement*

input — $\underbrace{N_c \quad N_f \quad m_q \quad \alpha_s}_{\text{defining gauge theory}} \qquad \underbrace{f_\pi \quad m_\pi}_{\text{universal low energy parameters}}$

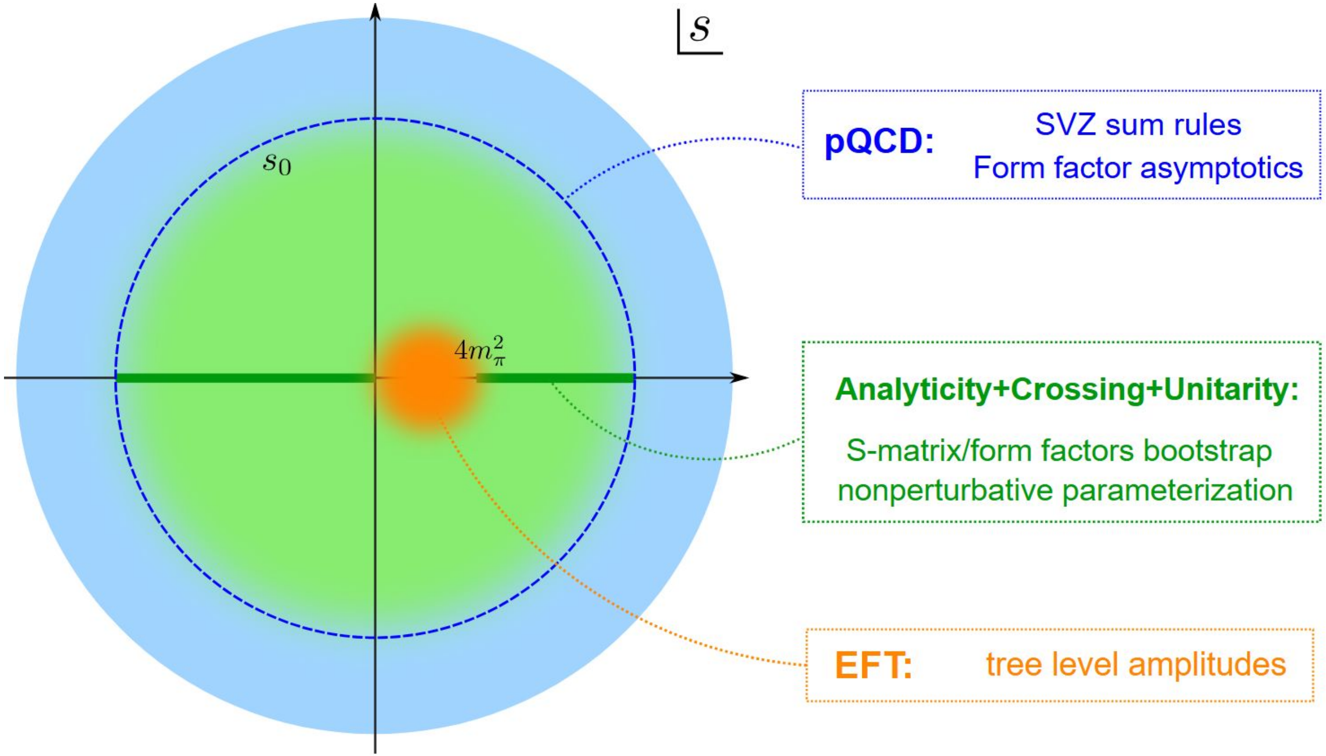
defining gauge theory *universal low energy parameters*

theoretical/numerical computation, not using experimental scattering data as input

Gauge Theory Bootstrap: summary

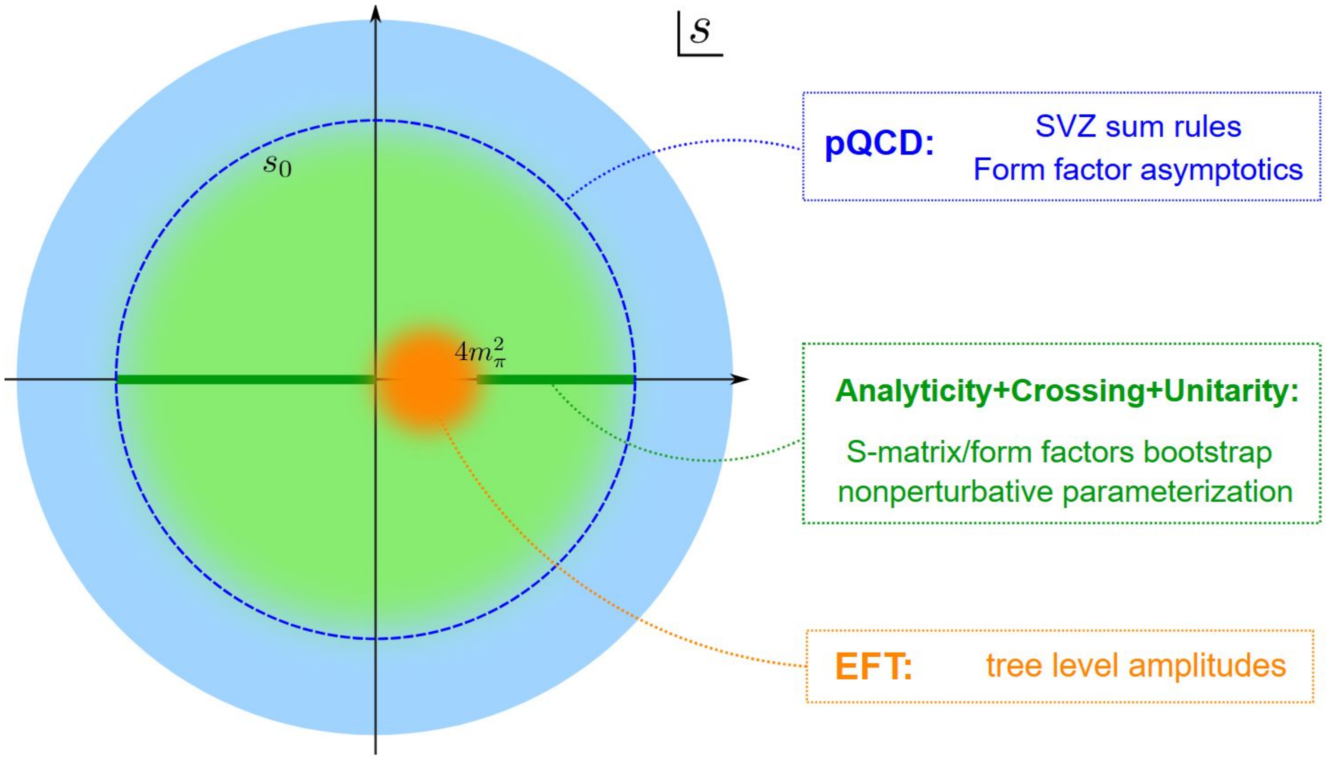


Gauge Theory Bootstrap: summary



look for amplitudes/form factors that: 1, satisfy generic consistency conditions (analyticity, crossing, unitarity)
2, match low energy behavior (chiSB) and high energy (pQCD)

Gauge Theory Bootstrap: summary



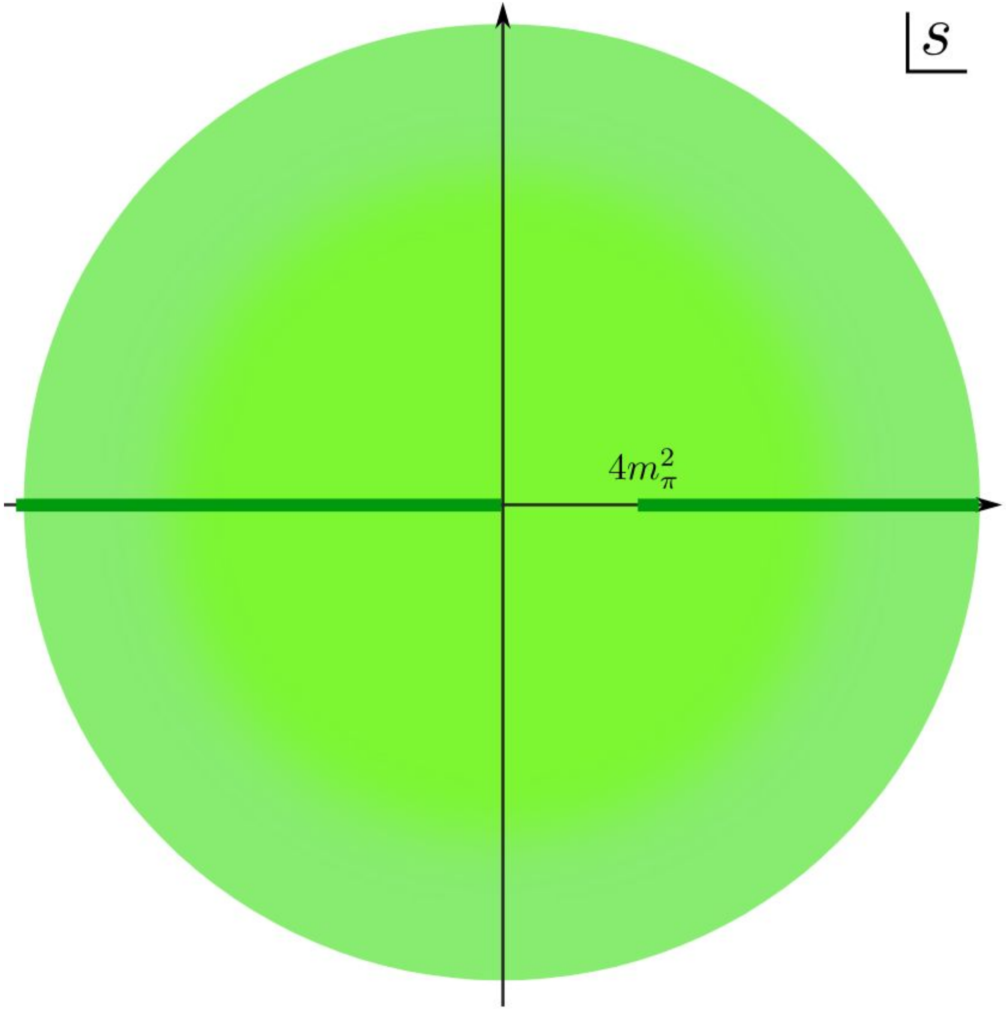
test with

$N_f = 2 \quad N_c = 3$

compare with experiments

changeable in the approach

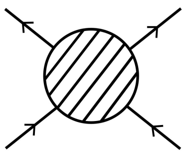
look for amplitudes/form factors that: 1, satisfy generic consistency conditions (analyticity, crossing, unitarity)
2, match low energy behavior (chiSB) and high energy (pQCD)



Analyticity+Crossing+Unitarity:

S-matrix bootstrap
nonperturbative parameterization

S-matrix bootstrap parameterization



modern S-matrix bootstrap: [Paulos, Penedones, Toledo, van Rees, Vieira, 2016&2017]

$$\langle p_1, a; p_2, b|\mathbf{T}|p_3, c; p_4, d\rangle = A(s,t,u)\delta_{ab}\delta_{cd} + A(t,s,u)\delta_{ac}\delta_{bd} + A(u,t,s)\delta_{ad}\delta_{bc}$$

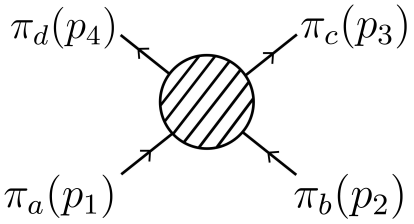
Crossing $A(s,t,u) = A(s,u,t)$

Analyticity cuts $s,t,u > 4$

$$m_\pi = 1$$

$$s+t+u = 4$$

S-matrix bootstrap parameterization



modern S-matrix bootstrap: [Paulos, Penedones, Toledo, van Rees, Vieira, 2016&2017]

$$\langle p_1, a; p_2, b|\mathbf{T}|p_3, c; p_4, d\rangle = A(s,t,u)\delta_{ab}\delta_{cd} + A(t,s,u)\delta_{ac}\delta_{bd} + A(u,t,s)\delta_{ad}\delta_{bc}$$

Crossing $A(s,t,u) = A(s,u,t)$ ***Analyticity*** cuts $s,t,u > 4$

$$m_\pi = 1$$

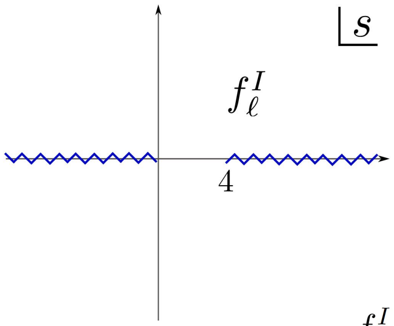
$$s + t + u = 4$$

nonperturbative parameterization encoding **Analyticity** and **Crossing**:

$$A(s,t,u) = \frac{1}{\pi^2}\int_4^\infty dx \int_4^\infty dy \left[\frac{\rho_1(x,y)}{(x-s)(y-t)} + \frac{\rho_1(x,y)}{(x-s)(y-u)} + \frac{\rho_2(x,y)}{(x-t)(y-u)}\right]$$ +subtraction terms

parameters: $\{\rho_{\alpha=1,2}(x,y), \dots\}$ numerics: discretize $\{\rho_{\alpha,ij}, \dots\}$ bootstrap variables

S-matrix bootstrap parameterization



analytic function of *s*

$$\begin{aligned} T^{I=0}(s,t,u) &= 3A(s,t,u) + A(t,s,u) + A(u,t,s) \\ T^{I=1}(s,t,u) &= A(t,s,u) - A(u,t,s) \\ T^{I=2}(s,t,u) &= A(t,s,u) + A(u,t,s) \end{aligned}$$

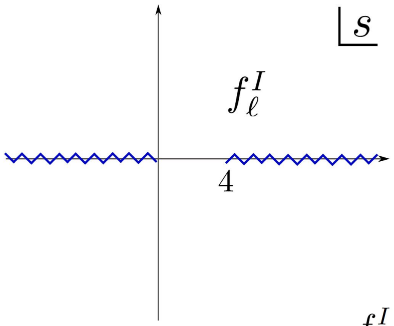
$SU(2)_V$ isospin

Symmetry

$\Downarrow$

$$f_\ell^I(s) = \frac{1}{4}\int_{-1}^{1} d\cos\theta P_\ell(\cos\theta) T^I(s,t)$$

S-matrix bootstrap parameterization



analytic function of *s*

$$
\begin{aligned}
T^{I=0}(s,t,u) &= 3A(s,t,u) + A(t,s,u) + A(u,t,s) \\
T^{I=1}(s,t,u) &= A(t,s,u) - A(u,t,s) \\
T^{I=2}(s,t,u) &= A(t,s,u) + A(u,t,s)
\end{aligned}
$$

$SU(2)_V$ isospin

Symmetry

⇓

$$f_\ell^I(s) = \frac{1}{4}\int_{-1}^{1} d\cos\theta \, P_\ell(\cos\theta) T^I(s,t)$$

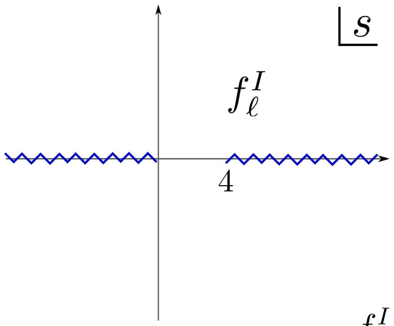
physical kinematic region $s > 4$

$$S_\ell^I(s) = 1 + i\pi\sqrt{\frac{s-4}{s}} f_\ell^I(s) = \eta_\ell^I(s)\, e^{2i\delta_\ell^I(s)}$$

elasticity phase shift

unphysical region $f_\ell^I(0 < s < 4)$ *real linear functionals of bootstrap variables*

S-matrix bootstrap parameterization



analytic function of *s*

$$
\begin{aligned}
T^{I=0}(s,t,u) &= 3A(s,t,u) + A(t,s,u) + A(u,t,s) \\
T^{I=1}(s,t,u) &= A(t,s,u) - A(u,t,s) \\
T^{I=2}(s,t,u) &= A(t,s,u) + A(u,t,s)
\end{aligned}
$$

$SU(2)_V$ isospin

Symmetry

⇓

$$f_\ell^I(s) = \frac{1}{4}\int_{-1}^{1} d\cos\theta \, P_\ell(\cos\theta) T^I(s,t)$$

physical kinematic region $s > 4$

$$S_\ell^I(s) = 1 + i\pi\sqrt{\frac{s-4}{s}} f_\ell^I(s) = \eta_\ell^I(s)\, e^{2i\delta_\ell^I(s)}$$

elasticity phase shift

unphysical region $f_\ell^I(0 < s < 4)$ *real linear functionals of bootstrap variables*

Unitarity

$$|S_\ell^I(s^+)| \leq 1, \; s > 4 \quad \forall \ell, I$$

$$\begin{pmatrix} 1 & S_\ell^I(s) \\ S_\ell^{I*}(s) & 1 \end{pmatrix} \succeq 0$$

positive semidefinite → convex space of amplitudes

convex optimization

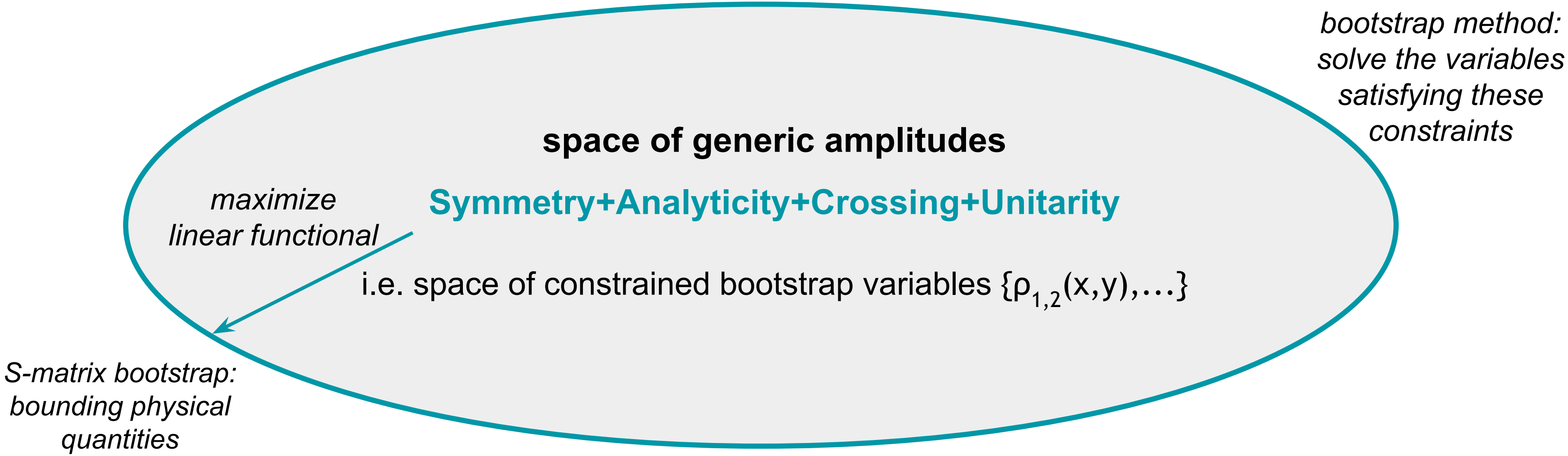
Bootstrap methods for nonperturbative computations

bootstrap method:
solve the variables
satisfying these
constraints

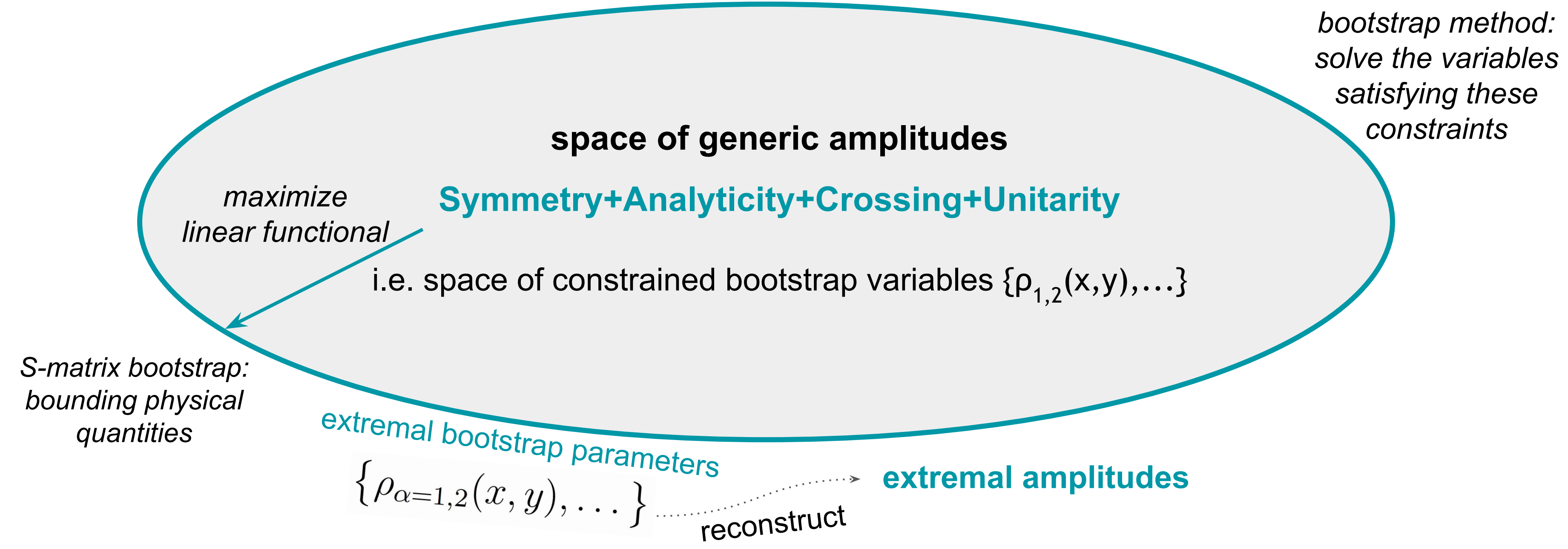
Symmetry+Analyticity+Crossing+Unitarity

bootstrap variables $\{\rho_{1,2}(x,y),\ldots\}$

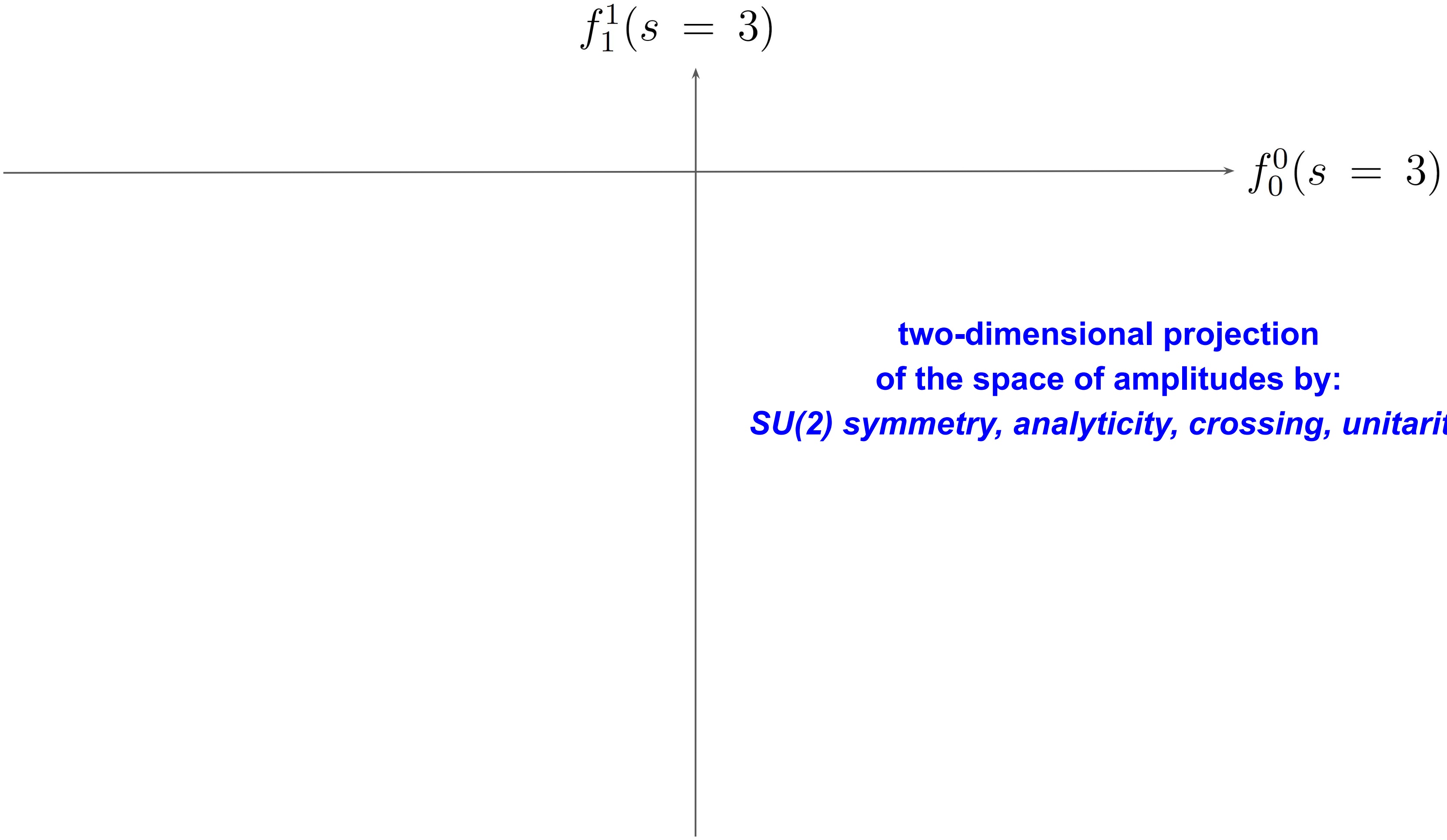
Space of generic bootstrap solutions



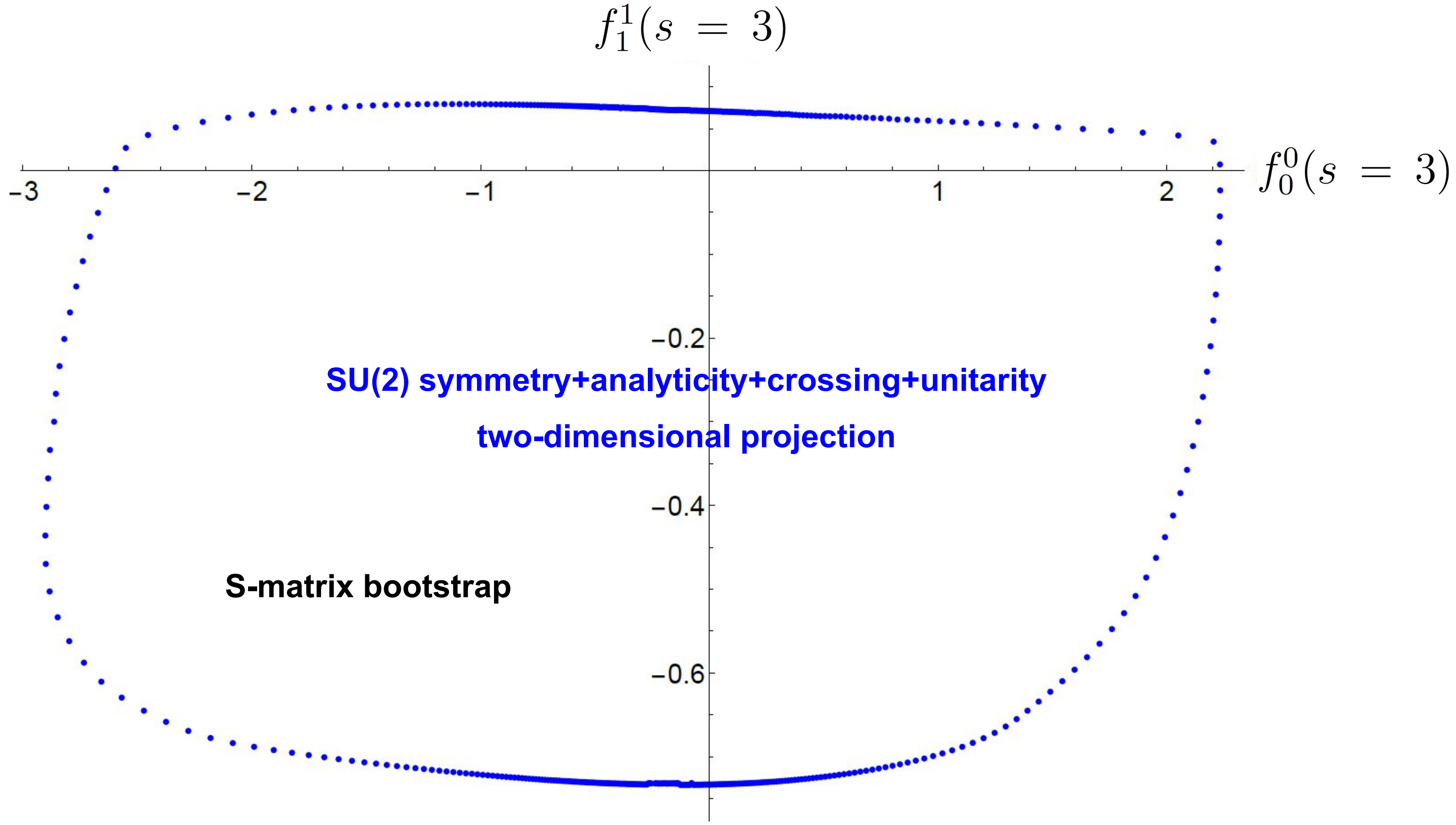
Bootstrap maximization → nonperturbative computations



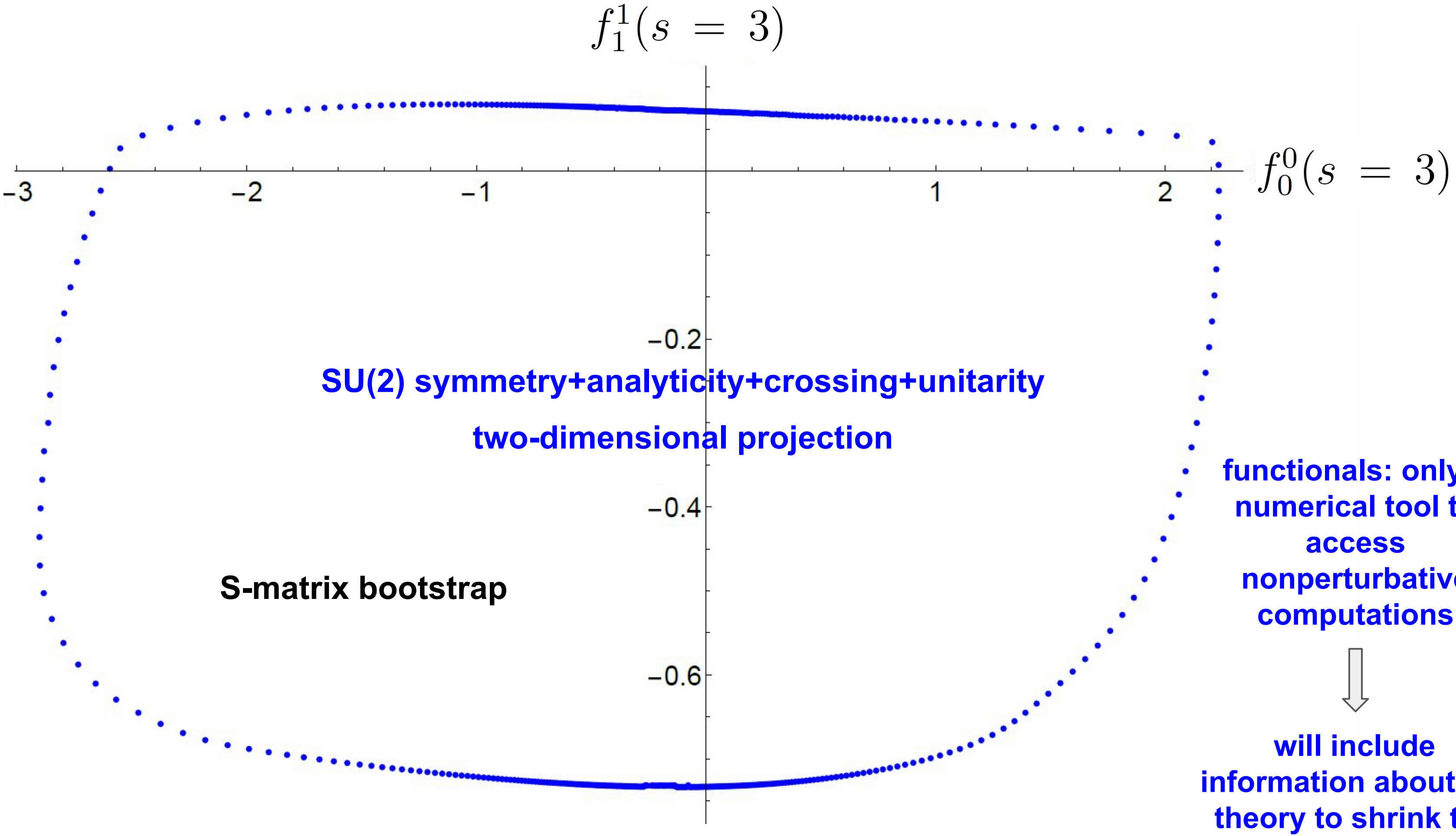
maximization → non-perturbative numerical computation of scattering amplitudes



two-dimensional projection
of the space of amplitudes by:
SU(2) symmetry, analyticity, crossing, unitarity



each boundary point: an extremal numerical amplitude



SU(2) symmetry+analyticity+crossing+unitarity

two-dimensional projection

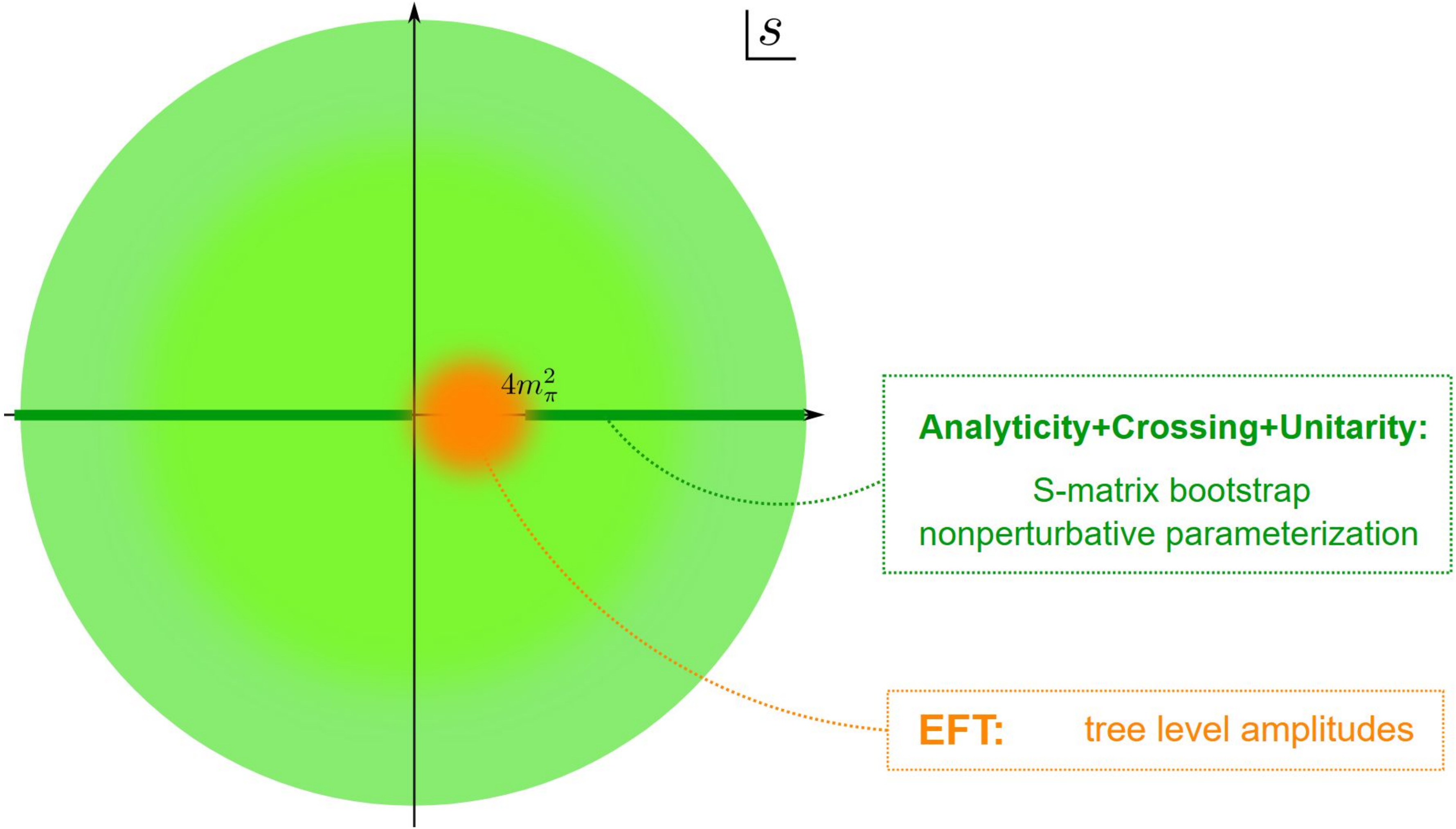
S-matrix bootstrap

functionals: only a numerical tool to access nonperturbative computations

⇩

will include information about the theory to shrink the space

each boundary point: an extremal numerical amplitude



s
$4m_\pi^2$
Analyticity+Crossing+Unitarity:
S-matrix bootstrap
nonperturbative parameterization
EFT:
tree level amplitudes

Weakly coupled Goldstone bosons

chiral symmetry breaking: weakly coupled Goldstone bosons at very low energy

interaction: $$\mathcal{L}_2^{4\pi} = \frac{1}{6f_\pi^2}\Big((\vec{\pi}\cdot\partial_\mu\vec{\pi})^2 - \vec{\pi}^2(\partial_\mu\vec{\pi}\cdot\partial^\mu\vec{\pi})\Big) + \frac{m_\pi^2}{24f_\pi^2}(\vec{\pi}^2)^2$$

tree-level amplitude: $$A_{\text{tree}}(s,t,u) = \frac{4}{\pi}\frac{s-m_\pi^2}{32\pi f_\pi^2}$$ *linear in s* [Weinberg, 1966]

good in the unphysical region (very low energy) $0 < s,t,u < 4m_\pi^2$

Weakly coupled Goldstone bosons

chiral symmetry breaking: weakly coupled Goldstone bosons at very low energy

interaction: $$\mathcal{L}_2^{4\pi} = \frac{1}{6f_\pi^2}\Big((\vec{\pi}\cdot\partial_\mu\vec{\pi})^2 - \vec{\pi}^2(\partial_\mu\vec{\pi}\cdot\partial^\mu\vec{\pi})\Big) + \frac{m_\pi^2}{24f_\pi^2}(\vec{\pi}^2)^2$$

tree-level amplitude: $$A_{\text{tree}}(s,t,u) = \frac{4}{\pi}\frac{s-m_\pi^2}{32\pi f_\pi^2}$$ *linear in s* [Weinberg, 1966]

good in the unphysical region (very low energy) $0 < s,t,u < 4m_\pi^2$

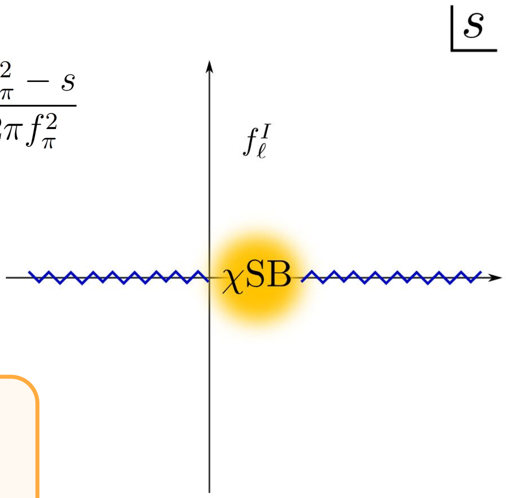
S0: $f^0_{0,\text{tree}}(s) = \frac{2}{\pi}\frac{2s-m_\pi^2}{32\pi f_\pi^2}$ P1: $f^1_{1,\text{tree}}(s) = \frac{2}{\pi}\frac{s-4m_\pi^2}{96\pi f_\pi^2}$ S2: $f^2_{0,\text{tree}}(s) = \frac{2}{\pi}\frac{2m_\pi^2-s}{32\pi f_\pi^2}$

good in unphysical region (very low energy) $0 < s < 4m_\pi^2$

Chiral symmetry breaking input

approximate linear behavior at very low energy: input in gauge theory bootstrap

S0: $f^0_{0,\text{tree}}(s) = \frac{2}{\pi}\frac{2s - m_\pi^2}{32\pi f_\pi^2}$ P1: $f^1_{1,\text{tree}}(s) = \frac{2}{\pi}\frac{s - 4m_\pi^2}{96\pi f_\pi^2}$ S2: $f^2_{0,\text{tree}}(s) = \frac{2}{\pi}\frac{2m_\pi^2 - s}{32\pi f_\pi^2}$



numerically

requires p.w. in the bootstrap match the tree level p.w. in unphysical region

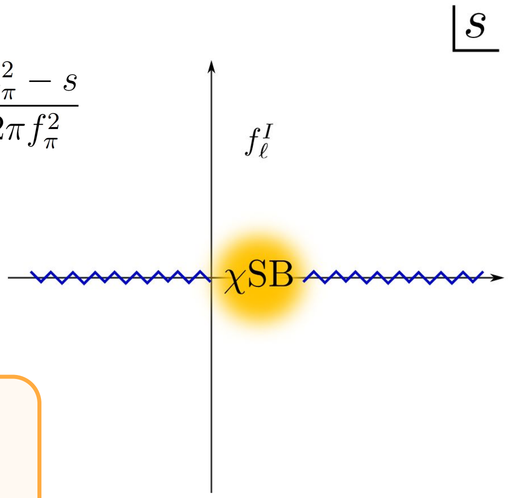
$$f^0_0(s) \simeq f^0_{0,\text{tree}}(s) \qquad f^1_1(s) \simeq f^1_{1,\text{tree}}(s) \qquad f^2_0(s) \simeq f^2_{0,\text{tree}}(s) \qquad 0 < s < 4m_\pi^2$$

constraints on bootstrap variables

Chiral symmetry breaking input

approximate linear behavior at very low energy: input in gauge theory bootstrap

S0: $f^0_{0,\text{tree}}(s) = \frac{2}{\pi}\frac{2s - m_\pi^2}{32\pi f_\pi^2}$ P1: $f^1_{1,\text{tree}}(s) = \frac{2}{\pi}\frac{s - 4m_\pi^2}{96\pi f_\pi^2}$ S2: $f^2_{0,\text{tree}}(s) = \frac{2}{\pi}\frac{2m_\pi^2 - s}{32\pi f_\pi^2}$



numerically

requires p.w. in the bootstrap match the tree level p.w. in unphysical region

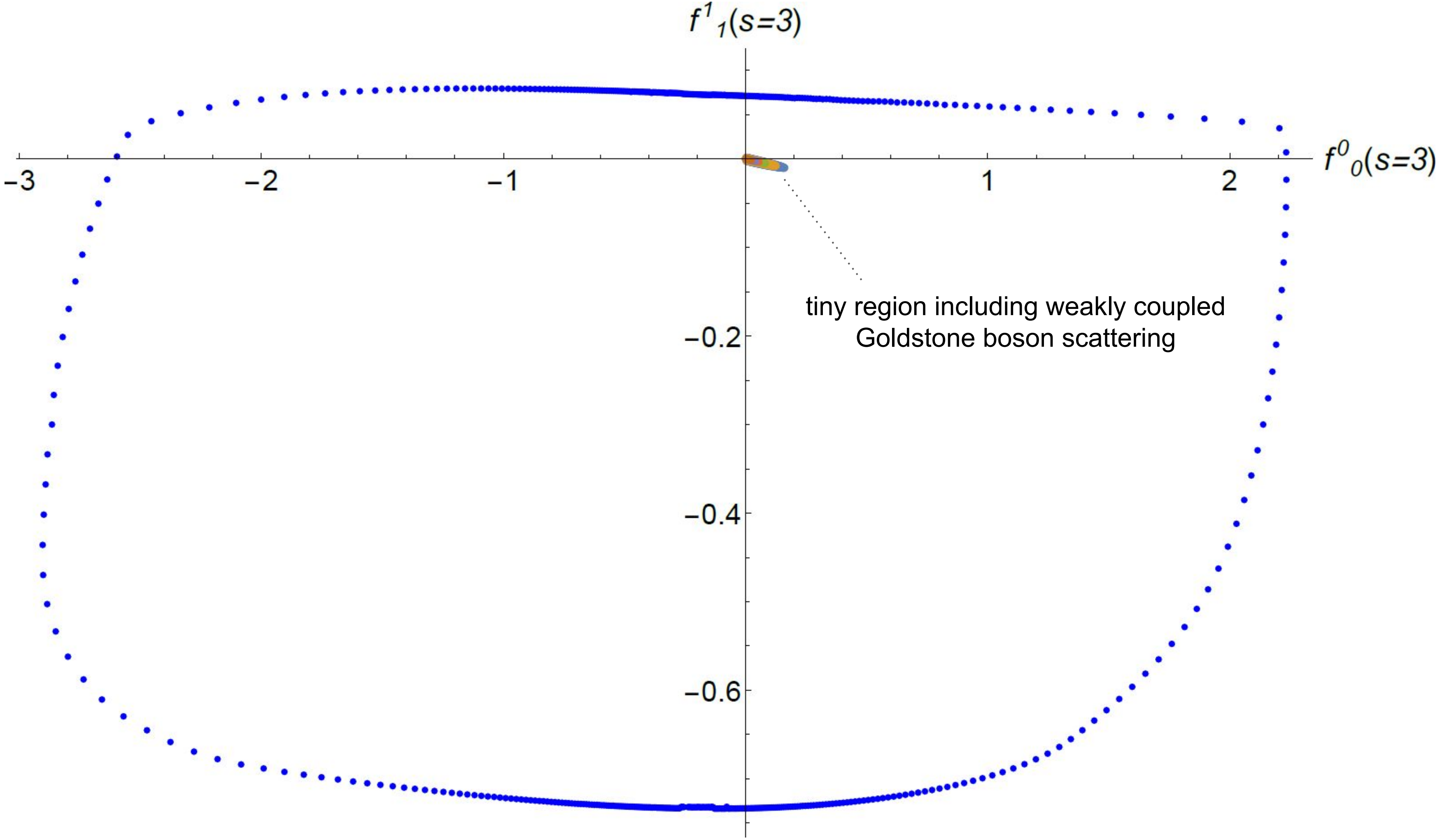
$$f^0_0(s) \simeq f^0_{0,\text{tree}}(s) \qquad f^1_1(s) \simeq f^1_{1,\text{tree}}(s) \qquad f^2_0(s) \simeq f^2_{0,\text{tree}}(s) \qquad 0 < s < 4m_\pi^2$$

constraints on bootstrap variables

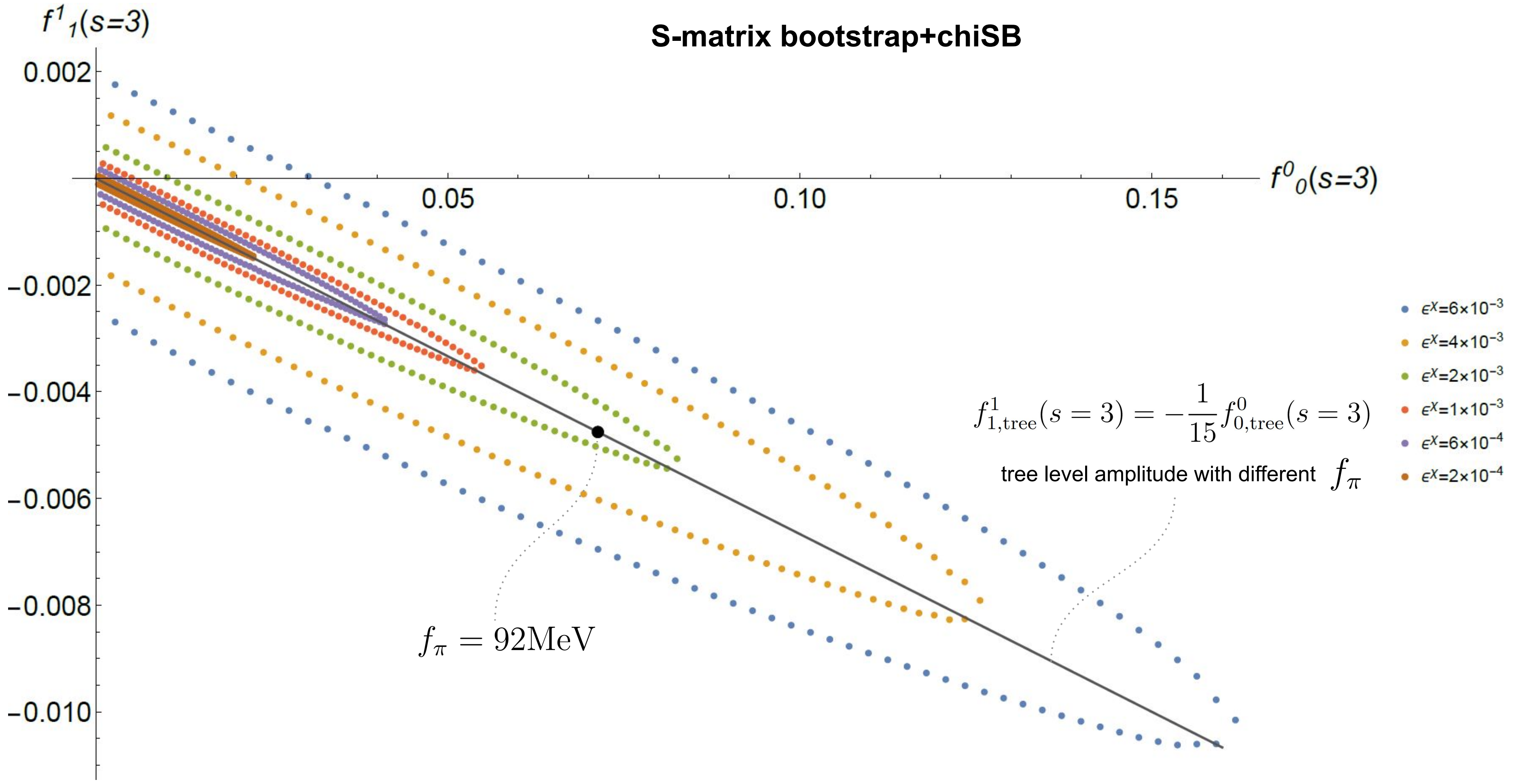
ϵ^χ {
too loose: large deviation from chiSB prediction
too tight: exclude the desired theory

numerics with a series of tolerance

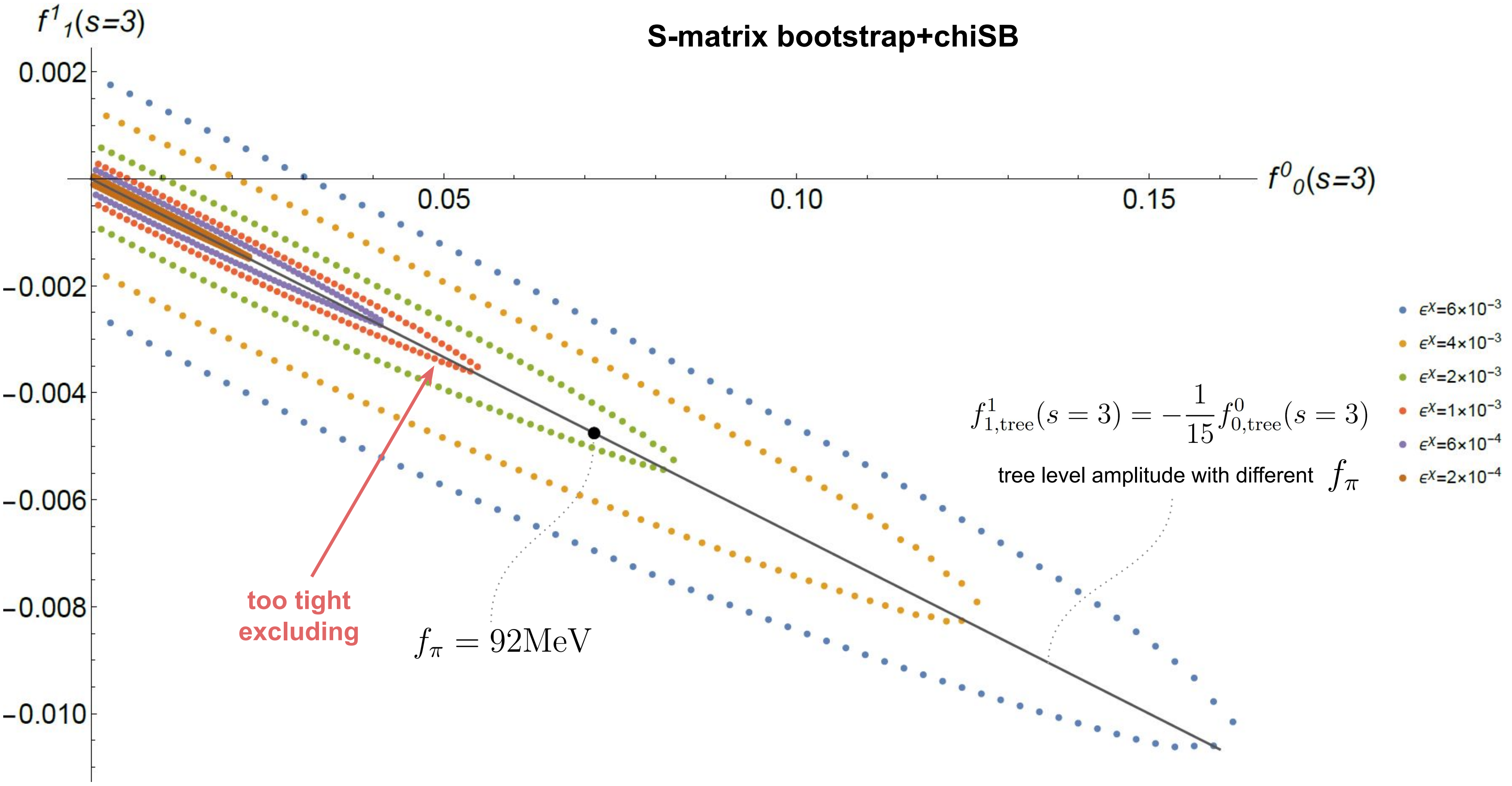
use $f_\pi = 92\text{MeV}$ **to select appropriate tolerance**



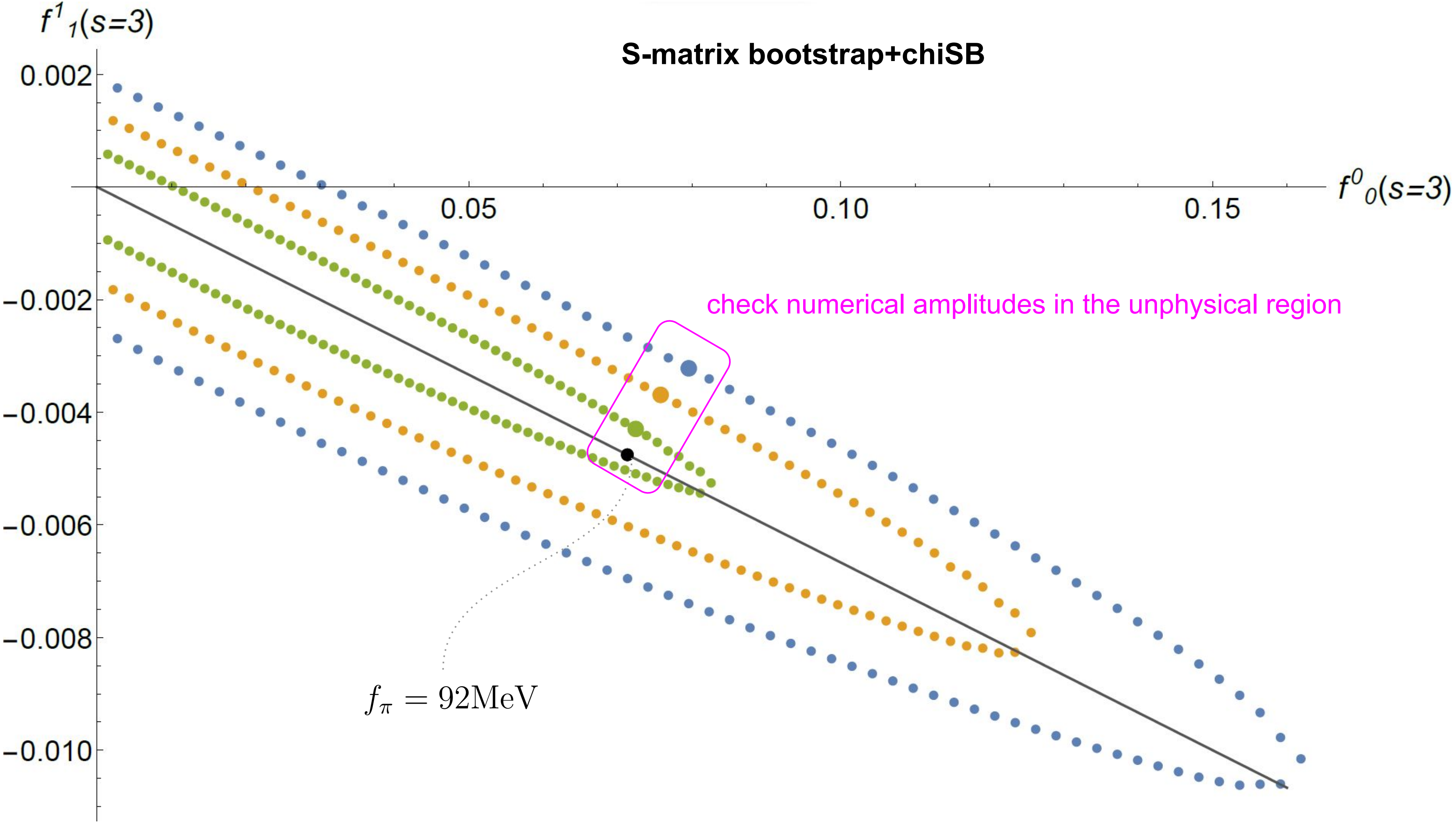
$f^1_1(s=3)$
$f^0_0(s=3)$
−3
−2
−1
1
2
−0.2
−0.4
−0.6
tiny region including weakly coupled
Goldstone boson scattering

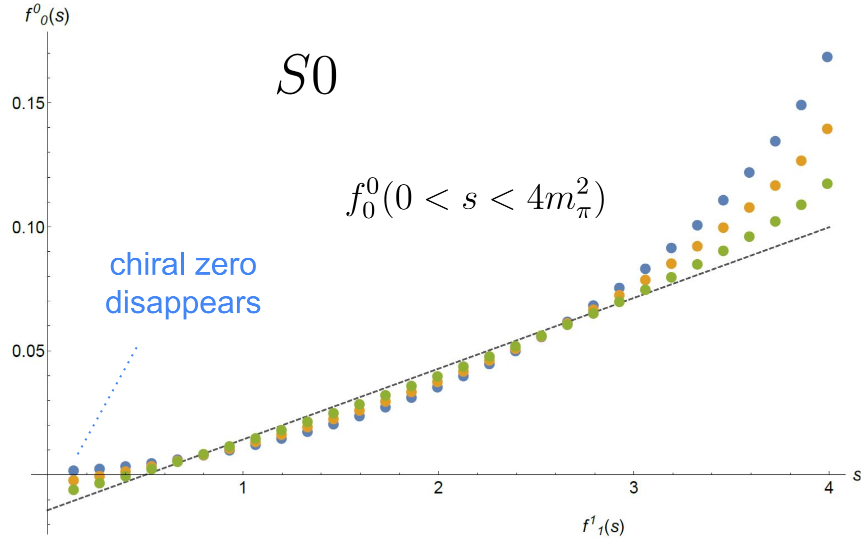


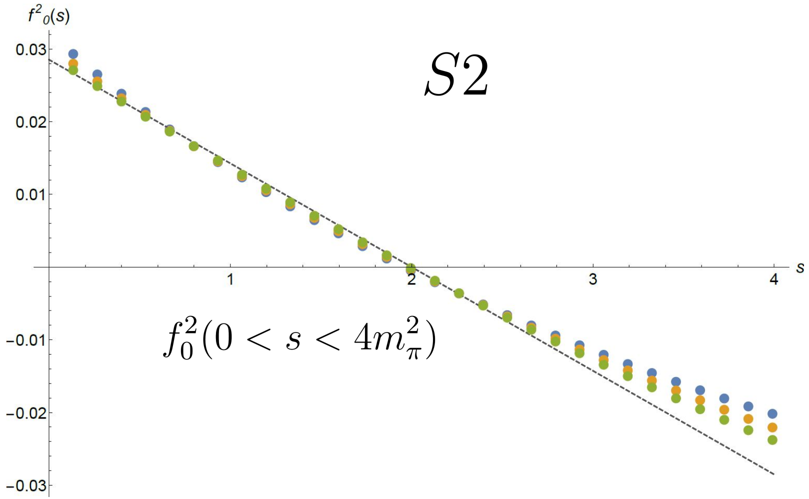
S-matrix bootstrap+chiSB
$f^1{}_1(s=3)$
$f^0{}_0(s=3)$
0.002
−0.002
−0.004
−0.006
−0.008
−0.010
0.05
0.10
0.15
$\epsilon^X=6\times10^{-3}$
$\epsilon^X=4\times10^{-3}$
$\epsilon^X=2\times10^{-3}$
$\epsilon^X=1\times10^{-3}$
$\epsilon^X=6\times10^{-4}$
$\epsilon^X=2\times10^{-4}$
$f^1_{1,\text{tree}}(s=3) = -\frac{1}{15} f^0_{0,\text{tree}}(s=3)$
tree level amplitude with different f_π
$f_\pi = 92\text{MeV}$



S-matrix bootstrap+chiSB
$f^1{}_1(s=3)$
$f^0{}_0(s=3)$
0.002
−0.002
−0.004
−0.006
−0.008
−0.010
0.05
0.10
0.15
$\epsilon^X=6\times10^{-3}$
$\epsilon^X=4\times10^{-3}$
$\epsilon^X=2\times10^{-3}$
$\epsilon^X=1\times10^{-3}$
$\epsilon^X=6\times10^{-4}$
$\epsilon^X=2\times10^{-4}$
$f^1_{1,\mathrm{tree}}(s=3) = -\frac{1}{15} f^0_{0,\mathrm{tree}}(s=3)$
tree level amplitude with different f_π
too tight excluding
$f_\pi = 92\mathrm{MeV}$


S-matrix bootstrap+chiSB
$f^1{}_1(s=3)$
$f^0{}_0(s=3)$
0.002
−0.002
−0.004
−0.006
−0.008
−0.010
0.05
0.10
0.15
check numerical amplitudes in the unphysical region
$f_\pi = 92\text{MeV}$

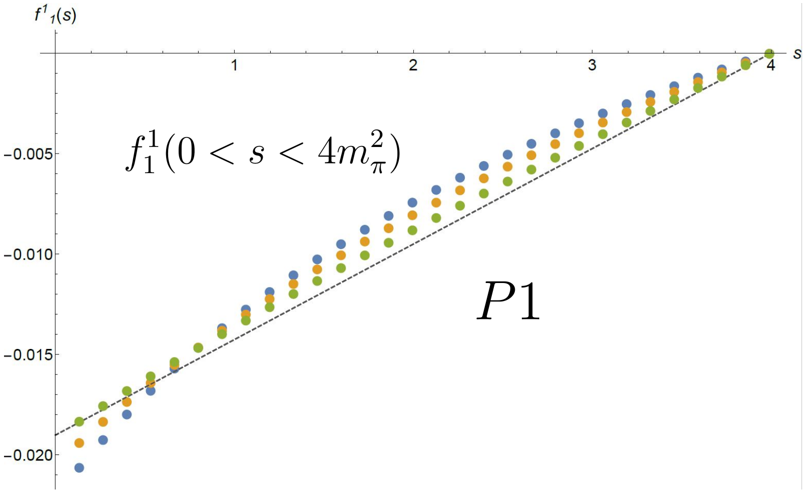




dashed lines: tree-level p.w.

$f^0_{0,\mathrm{tree}}(s) \quad f^1_{1,\mathrm{tree}}(s) \quad f^2_{0,\mathrm{tree}}(s)$

$f_\pi = 92\mathrm{MeV}$

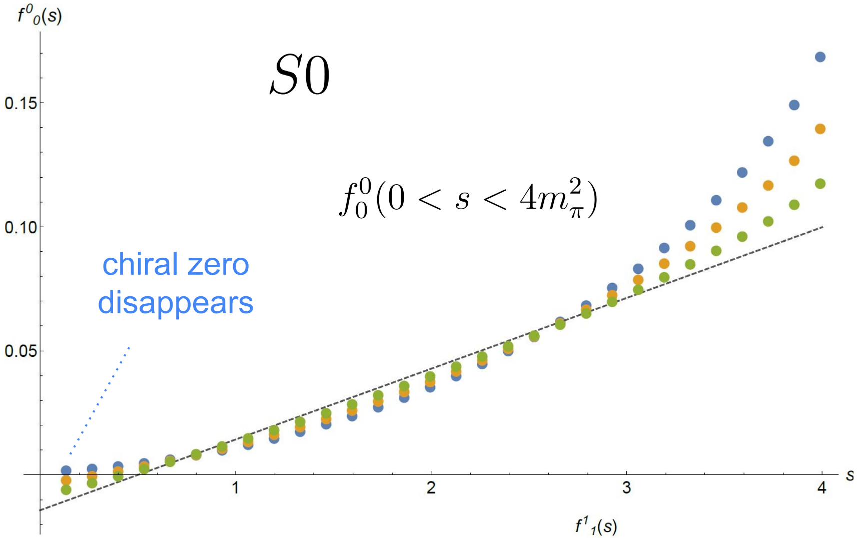


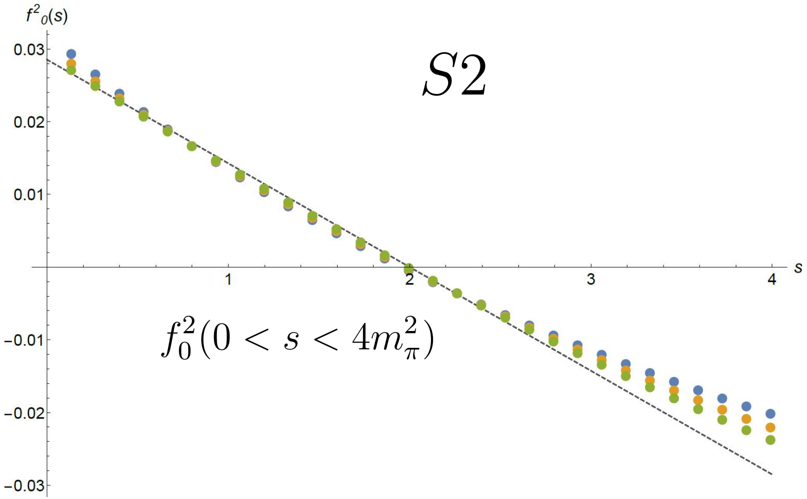
$\epsilon^\chi = 6\times10^{-3}$

$\epsilon^\chi = 4\times10^{-3}$

$\epsilon^\chi = 2\times10^{-3}$

$0 < s < 4m_\pi^2$

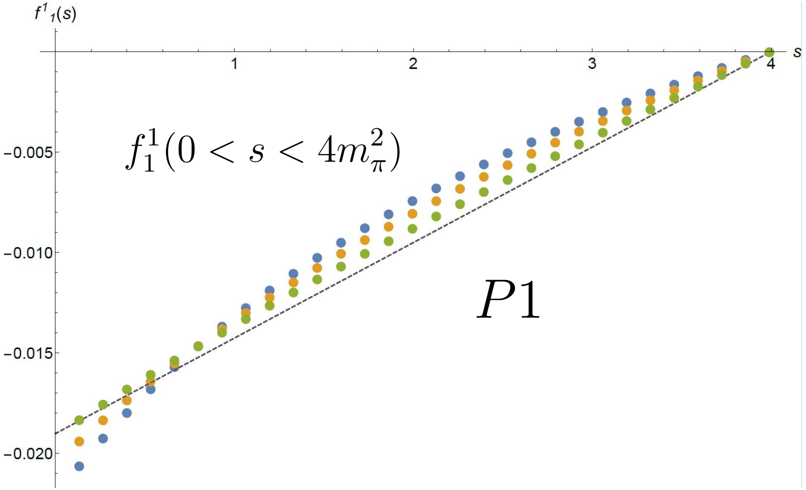




dashed lines: tree-level p.w.

$f^0_{0,\mathrm{tree}}(s)$ $f^1_{1,\mathrm{tree}}(s)$ $f^2_{0,\mathrm{tree}}(s)$

$f_\pi = 92\mathrm{MeV}$

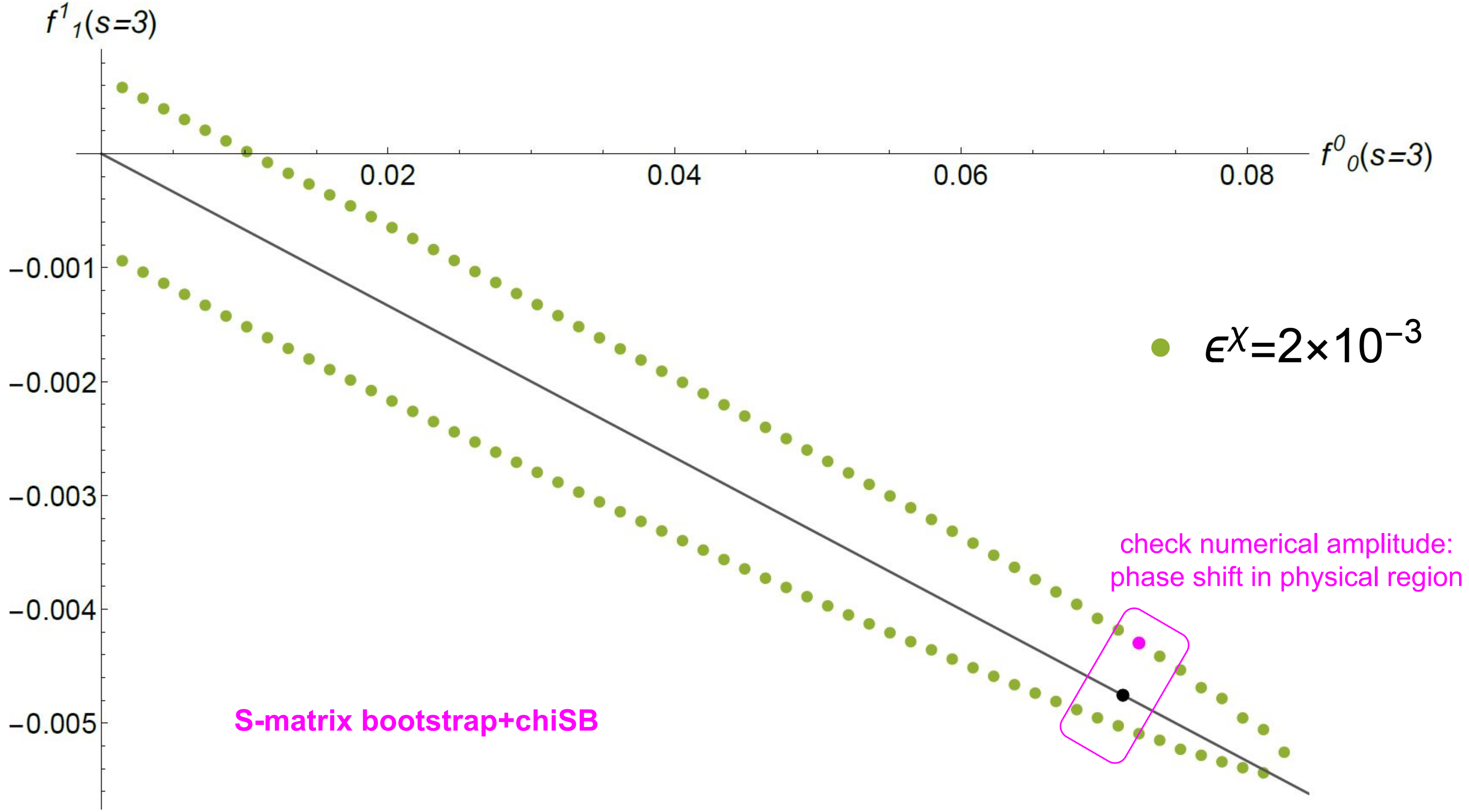


$\epsilon^\chi = 6\times10^{-3}$

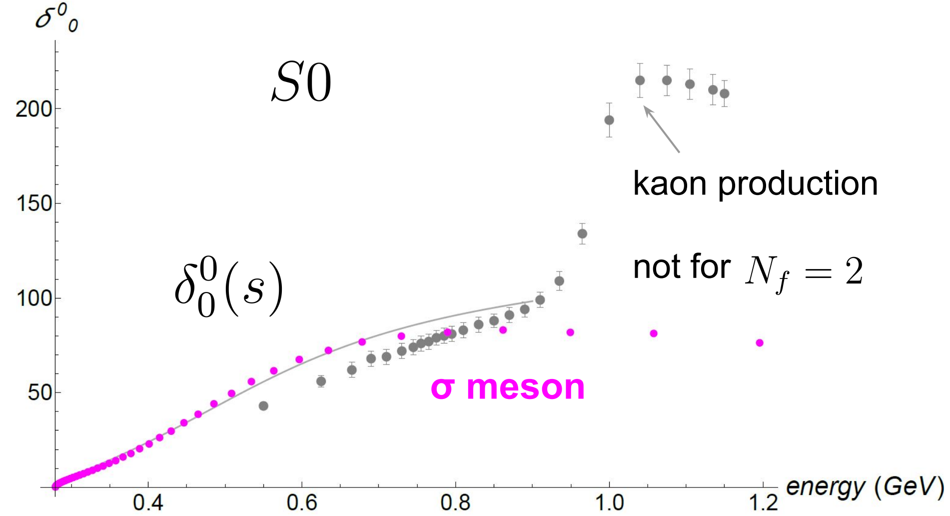
$\epsilon^\chi = 4\times10^{-3}$

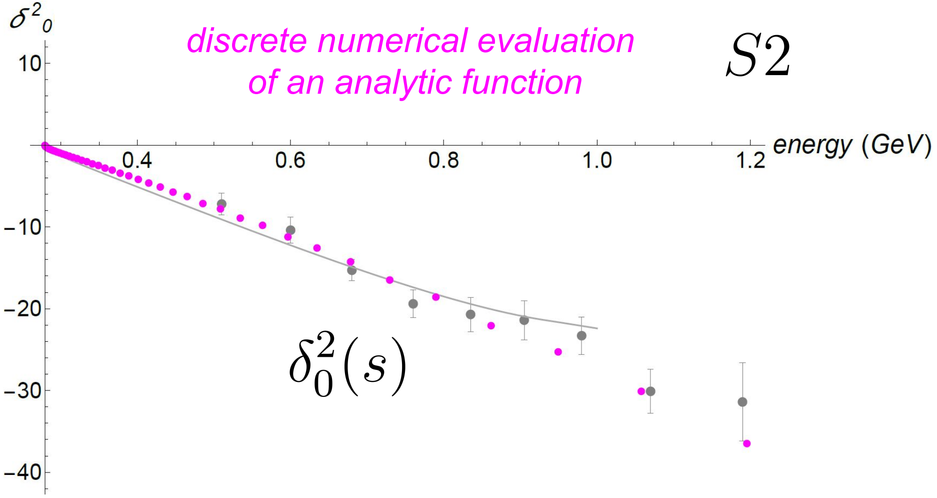
$\epsilon^\chi = 2\times10^{-3}$

$0 < s < 4m_\pi^2$

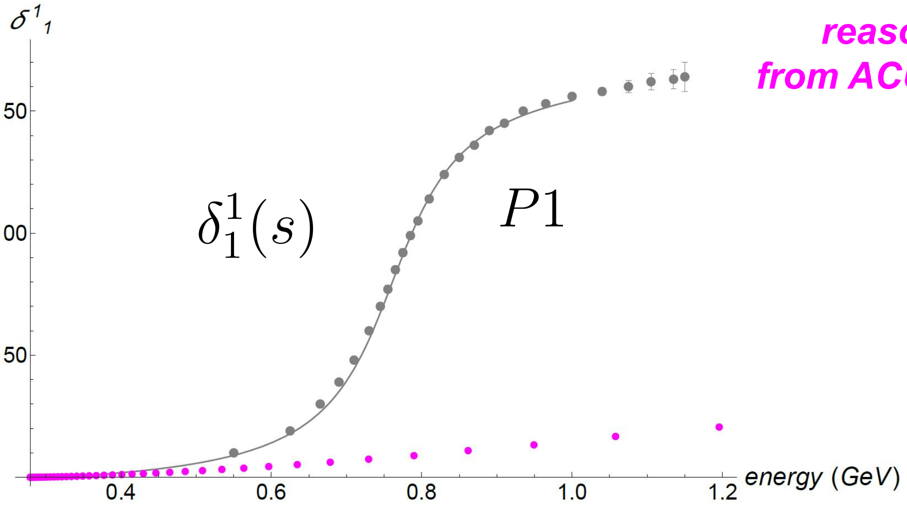


$f^1_1(s=3)$
$f^0_0(s=3)$
0.02
0.04
0.06
0.08
−0.001
−0.002
−0.003
−0.004
−0.005
$\epsilon^\chi=2\times10^{-3}$
check numerical amplitude:
phase shift in physical region
S-matrix bootstrap+chiSB



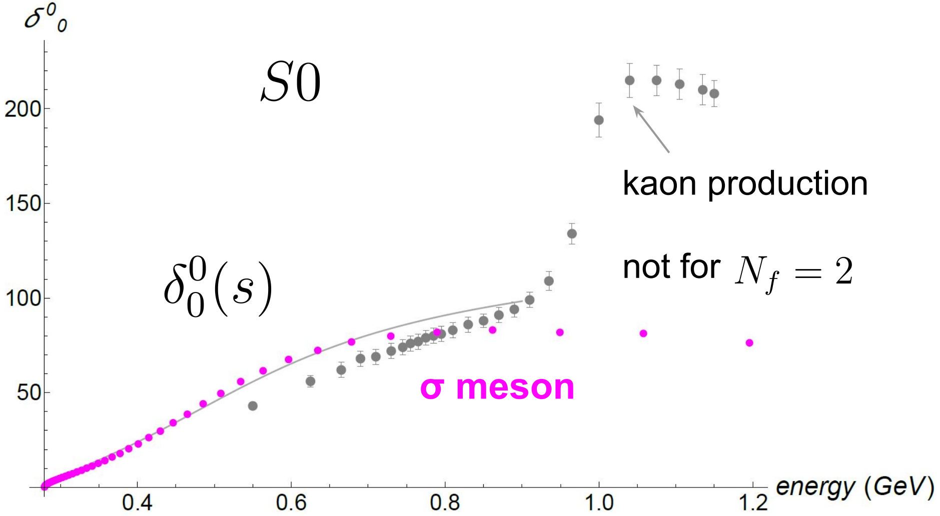


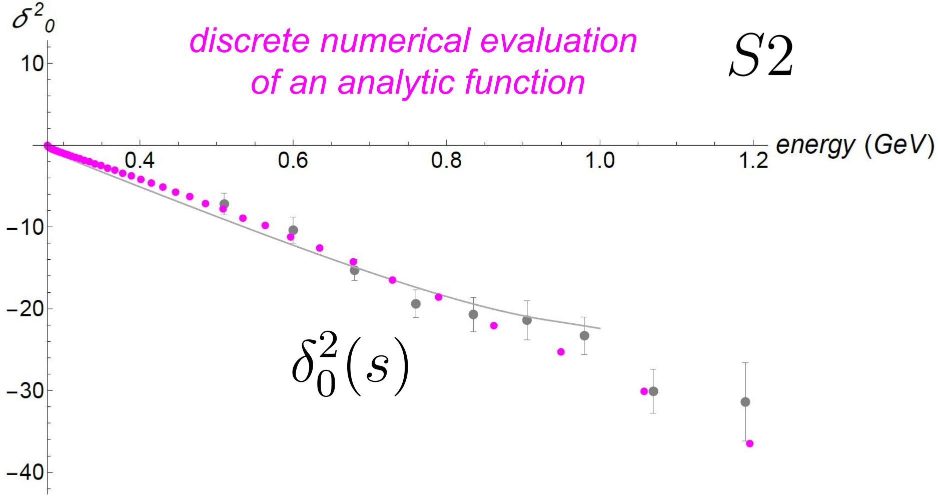
phase shifts with only chiSB (EFT) input



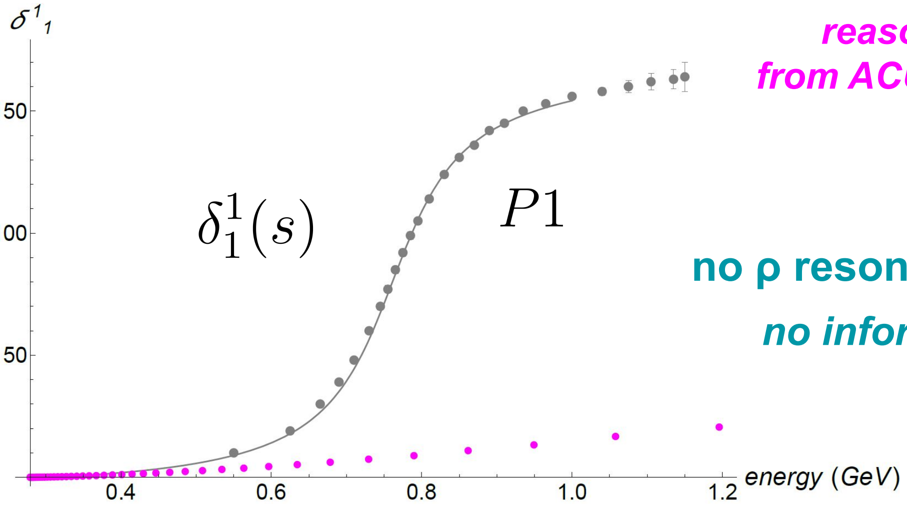
reasonable S0, S2 waves from ACU + weakly coupled EFT

experimental data (gray dots)
[Protopopescu et al, 1973]
[Losty et al, 1974]
pheno fit (gray line)
[Pelaez, Yndurain, 2005]





phase shifts with only chiSB (EFT) input

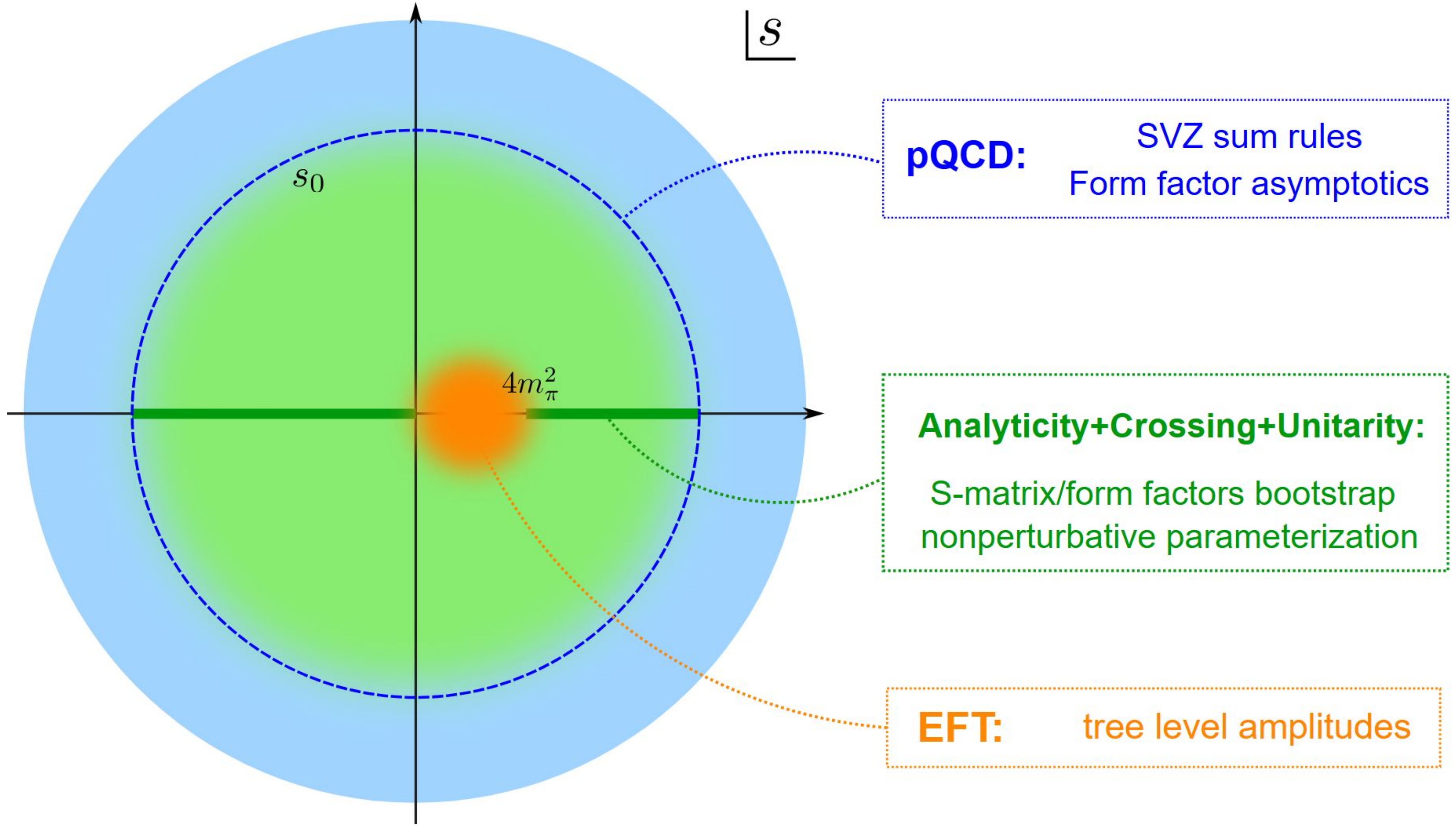


reasonable S0, S2 waves from ACU + weakly coupled EFT

experimental data (gray dots)
[Protopopescu et al, 1973]
[Losty et al, 1974]
pheno fit (gray line)
[Pelaez, Yndurain, 2005]

no ρ resonance without UV info

no information on QCD yet



s
s_0
$4m_\pi^2$
pQCD:
SVZ sum rules
Form factor asymptotics
Analyticity+Crossing+Unitarity:
S-matrix/form factors bootstrap
nonperturbative parameterization
EFT:
tree level amplitudes

S-matrix/form factor bootstrap

[Karateev, Kuhn, Penedones, 2019]

an important development: $|\psi_1\rangle = |p_1, p_2\rangle_{in}\,,\qquad |\psi_2\rangle = |p_1, p_2\rangle_{out}\,,\qquad |\psi_3\rangle = \int dx e^{-i(p_1+p_2)\cdot x}\mathcal{O}(x)|0\rangle$

state created by UV local operator

positive semidefinite matrix $\langle\psi_a|\psi_b\rangle = \begin{pmatrix} 1 & S & \mathcal{F} \\ S^* & 1 & \mathcal{F}^* \\ \mathcal{F}^* & \mathcal{F} & \rho \end{pmatrix} \succeq 0$

S-matrix/form factor bootstrap

[Karateev, Kuhn, Penedones, 2019]

an important development: $|\psi_1\rangle = |p_1, p_2\rangle_{in}\,, \qquad |\psi_2\rangle = |p_1, p_2\rangle_{out}\,, \qquad |\psi_3\rangle = \int dx e^{-i(p_1+p_2)\cdot x}\mathcal{O}(x)|0\rangle$

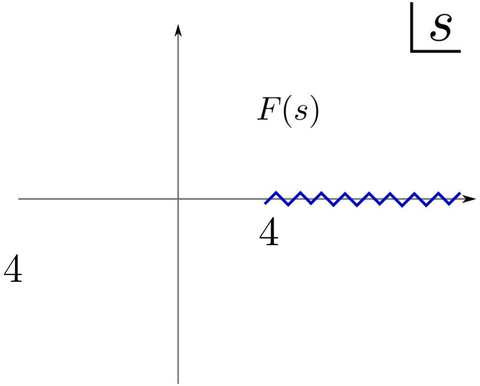
state created by UV local operator

positive semidefinite matrix $\langle\psi_a|\psi_b\rangle = \begin{pmatrix} 1 & S & \mathcal{F} \\ S^* & 1 & \mathcal{F}^* \\ \mathcal{F}^* & \mathcal{F} & \rho \end{pmatrix} \succeq 0$

2-particle form factor: ${}_{\text{out}}\langle p_1, p_2|\mathcal{O}(0)|0\rangle = F(s)$ *analytic function of s*

$$F(s) = \frac{1}{\pi}\int_4^\infty dx \frac{\text{Im}F(x)}{x-s}$$ +subtractions

spectral density: $\int \frac{d^4x}{(2\pi)^4} e^{iPx}\langle 0|\mathcal{O}^\dagger(x)\mathcal{O}(0)|0\rangle = \rho(s)$ *supported at $s > 4$*



S-matrix/form factor bootstrap

[Karateev, Kuhn, Penedones, 2019]

an important development: $|\psi_1\rangle = |p_1, p_2\rangle_{in}\,,\qquad |\psi_2\rangle = |p_1, p_2\rangle_{out}\,,\qquad |\psi_3\rangle = \int dx e^{-i(p_1+p_2)\cdot x}\mathcal{O}(x)|0\rangle$

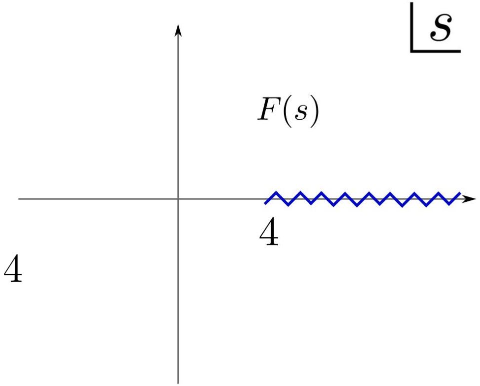
state created by UV local operator

positive semidefinite matrix $\langle\psi_a|\psi_b\rangle = \begin{pmatrix} 1 & S & \mathcal{F} \\ S^* & 1 & \mathcal{F}^* \\ \mathcal{F}^* & \mathcal{F} & \rho \end{pmatrix} \succeq 0$

2-particle form factor: ${}_{\text{out}}\langle p_1, p_2|\mathcal{O}(0)|0\rangle = F(s)$ *analytic function of s*

$$F(s) = \frac{1}{\pi}\int_4^\infty dx \frac{\text{Im}F(x)}{x-s}$$ +subtractions

spectral density: $\int \frac{d^4x}{(2\pi)^4} e^{iPx}\langle 0|\mathcal{O}^\dagger(x)\mathcal{O}(0)|0\rangle = \rho(s)$ *supported at* $s > 4$



bootstrap variables: $\{\rho_{1,2}(x,y), \ldots, \text{Im}F(x), \rho(x)\}$

allow connection with UV theory

Current correlators from the UV gauge theory

to connect with UV gauge theory

$$
\begin{array}{c} \\ \langle \text{in}|_{P',I,\ell} \\ \langle \text{out}|_{P',I,\ell} \\ \langle 0|\mathcal{O}^\dagger_{P',I,\ell} \end{array}
\begin{array}{c} \begin{array}{ccc} |\text{in}\rangle_{P,I,\ell} & |\text{out}\rangle_{P,I,\ell} & \mathcal{O}_{P,I,\ell}|0\rangle \end{array} \\ \begin{pmatrix} 1 & S_\ell^I(s) & \mathcal{F}_\ell^I \\ S_\ell^{I*}(s) & 1 & \mathcal{F}_\ell^{I*} \\ \mathcal{F}_\ell^{I*} & \mathcal{F}_\ell^I & \rho_\ell^I(s) \end{pmatrix} \end{array} \succeq 0 \qquad s > 4 \quad \forall \ell, I
$$

construct operators from gauge theory with desired quantum numbers

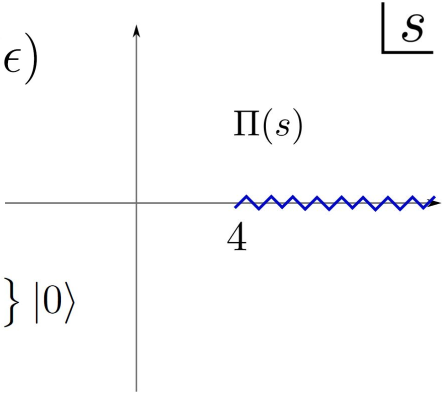
Current correlators from the UV gauge theory

to connect with UV gauge theory

$$
\begin{array}{c} \langle \text{in}|_{P',I,\ell} \\ \langle \text{out}|_{P',I,\ell} \\ \langle 0|\mathcal{O}^{\dagger}_{P',I,\ell} \end{array}
\overset{\begin{array}{ccc} |\text{in}\rangle_{P,I,\ell} & |\text{out}\rangle_{P,I,\ell} & \mathcal{O}_{P,I,\ell}|0\rangle \end{array}}{\begin{pmatrix} 1 & S^I_\ell(s) & \mathcal{F}^I_\ell \\ S^{I*}_\ell(s) & 1 & \mathcal{F}^{I*}_\ell \\ \mathcal{F}^{I*}_\ell & \mathcal{F}^I_\ell & \rho^I_\ell(s) \end{pmatrix}} \succeq 0 \qquad s > 4 \quad \forall \ell, I
$$

construct operators from gauge theory with desired quantum numbers

$$\rho^I_\ell(s) = 2\,\mathrm{Im}\Pi^I_\ell(x + i\epsilon)$$



e.g. isospin 1, spin 1
vector (electromagnetic) current

$$P1 \quad : \quad j^\mu_V(x) = \frac{1}{2}(\bar{u}\gamma^\mu u - \bar{d}\gamma^\mu d) \qquad \Pi^1_1(s) = i\int \frac{d^4x}{(2\pi)^4} e^{iPx} \langle 0|\hat{T}\left\{ j^\dagger_\sigma(x) j_\sigma(0)\right\}|0\rangle$$

$\vdots$

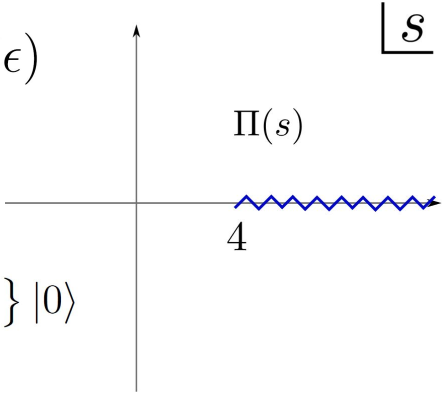
Current correlators from the UV gauge theory

to connect with UV gauge theory

$$
\begin{array}{c} \\ \langle \text{in}|_{P',I,\ell} \\ \langle \text{out}|_{P',I,\ell} \\ \langle 0|\mathcal{O}^\dagger_{P',I,\ell} \end{array}
\begin{array}{c} \begin{array}{ccc} |\text{in}\rangle_{P,I,\ell} & |\text{out}\rangle_{P,I,\ell} & \mathcal{O}_{P,I,\ell}|0\rangle \end{array} \\ \begin{pmatrix} 1 & S^I_\ell(s) & \mathcal{F}^I_\ell \\ S^{I*}_\ell(s) & 1 & \mathcal{F}^{I*}_\ell \\ \mathcal{F}^{I*}_\ell & \mathcal{F}^I_\ell & \rho^I_\ell(s) \end{pmatrix} \end{array} \succeq 0 \qquad s > 4 \quad \forall \ell, I
$$

construct operators from gauge theory with desired quantum numbers

$$\rho^I_\ell(s) = 2\,\text{Im}\Pi^I_\ell(x + i\epsilon)$$



e.g. isospin 1, spin 1
vector (electromagnetic) current

$$P1 \quad : \quad j^\mu_V(x) = \frac{1}{2}(\bar{u}\gamma^\mu u - \bar{d}\gamma^\mu d)$$

$$\vdots$$

$$\Pi^1_1(s) \quad = \quad i\int \frac{d^4x}{(2\pi)^4} e^{iPx} \langle 0|\hat{T}\left\{ j^\dagger_\sigma(x) j_\sigma(0)\right\}|0\rangle$$

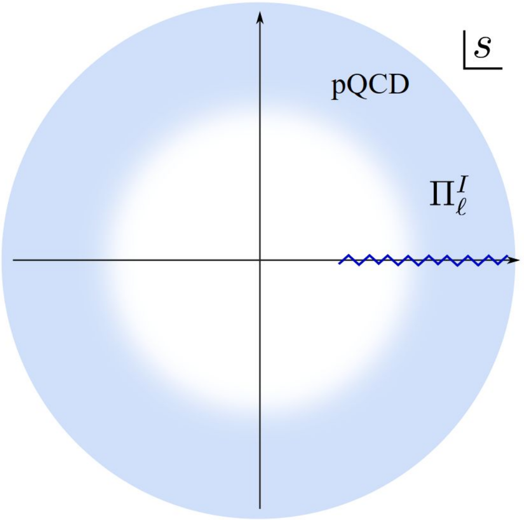
large spacelike momenta — asymptotic free region with pQCD computation

SVZ expansion

[Shifman, Vainshtein, Zakharov, 1979]

OPE: $$T\{j(x)j(0)\} = C_{\mathbb{1}}(x)\ \mathbb{1} + \sum_{\mathcal{O}} C_{\mathcal{O}}(x)\ \mathcal{O}(0)$$

$$\langle 0|T\{j(x)j(0)\}|0\rangle = C_{\mathbb{1}}(x) + C_{\bar{q}q}(x)\ \langle 0|m_q\bar{q}q|0\rangle + C_{G^2}(x)\ \langle 0|\frac{\alpha_s}{\pi}G^a_{\mu\nu}G^{a\,\mu\nu}|0\rangle + \ldots$$



SVZ expansion

[Shifman, Vainshtein, Zakharov, 1979]

OPE: $$T\{j(x)j(0)\} = C_{\mathbb{1}}(x)\ \mathbb{1} + \sum_{\mathcal{O}} C_{\mathcal{O}}(x)\ \mathcal{O}(0)$$

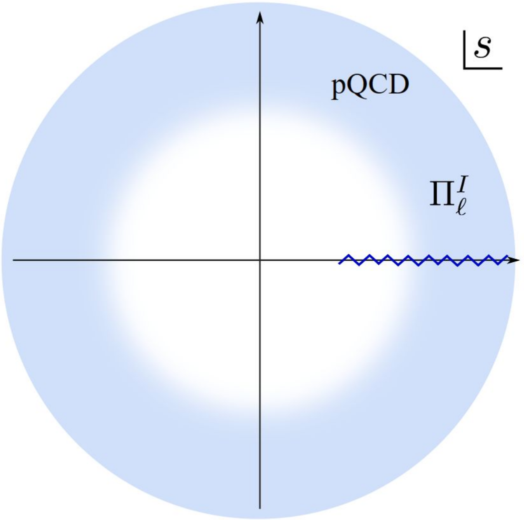
$$\langle 0|T\{j(x)j(0)\}|0\rangle = C_{\mathbb{1}}(x) + C_{\bar{q}q}(x)\ \langle 0|m_q\bar{q}q|0\rangle + C_{G^2}(x)\ \langle 0|\frac{\alpha_s}{\pi}G^a_{\mu\nu}G^{a\,\mu\nu}|0\rangle + \ldots$$

SB vacuum

quark condensate

gluon condensate

pQCD computation



SVZ expansion

[Shifman, Vainshtein, Zakharov, 1979]

OPE: $$T\{j(x)j(0)\} = C_{\mathbb{1}}(x)\ \mathbb{1} + \sum_{\mathcal{O}} C_{\mathcal{O}}(x)\ \mathcal{O}(0)$$

$$\langle 0|T\{j(x)j(0)\}|0\rangle = C_{\mathbb{1}}(x) + C_{\bar{q}q}(x)\ \langle 0|m_q\bar{q}q|0\rangle + C_{G^2}(x)\ \langle 0|\frac{\alpha_s}{\pi}G^a_{\mu\nu}G^{a\,\mu\nu}|0\rangle + \ldots$$

SB vacuum

quark condensate

gluon condensate

Fourier transform

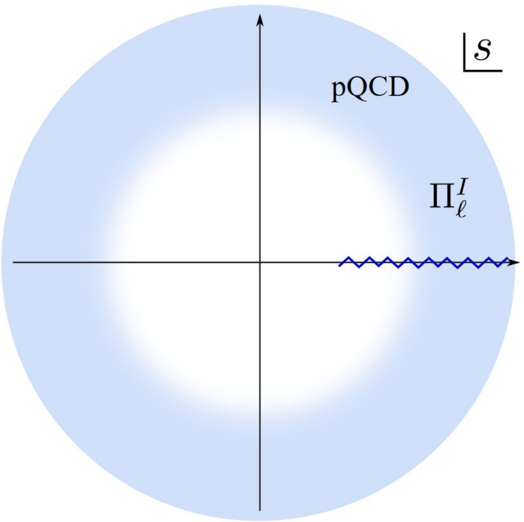
pQCD computation

large s expansion of vacuum polarization: e.g. vector current

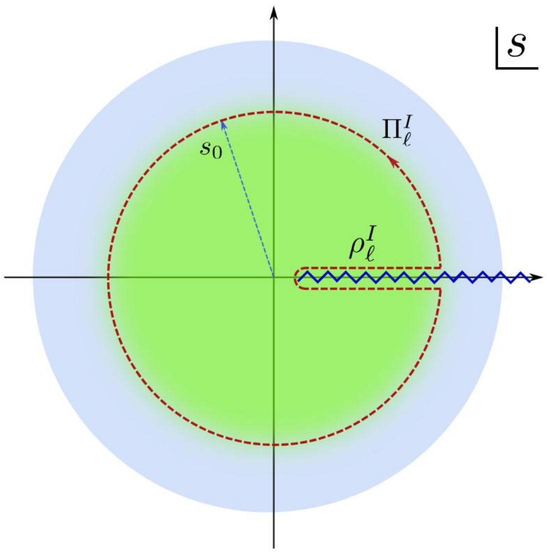
$$\Pi_1^1(s) = \frac{1}{2}\frac{1}{(2\pi)^4}\left\{-\frac{1}{4\pi^2}\left(1+\frac{\alpha_s}{\pi}\right) s\ln(-\frac{s}{\mu^2}) + \frac{1}{12s}\langle\frac{\alpha_s}{\pi}G^2\rangle + \frac{1}{s}\langle m_q\bar{q}q\rangle + \ldots\right\}$$

⋮

$N_c = 3$



Finite energy sum rule



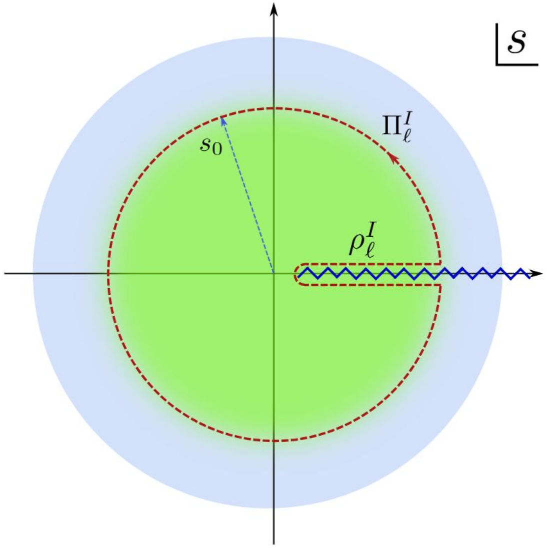
connect pQCD with bootstrap at s_0

contour integral $s^n \Pi(s)$ vanishes

SVZ

$$\int_4^{s_0} \rho(x) x^n dx = -s_0^{n+1} \int_0^{2\pi} e^{i(n+1)\varphi} \Pi(s_0 e^{i\varphi}) d\varphi$$

Finite energy sum rule



connect pQCD with bootstrap at s_0

contour integral $s^n\Pi(s)$ vanishes

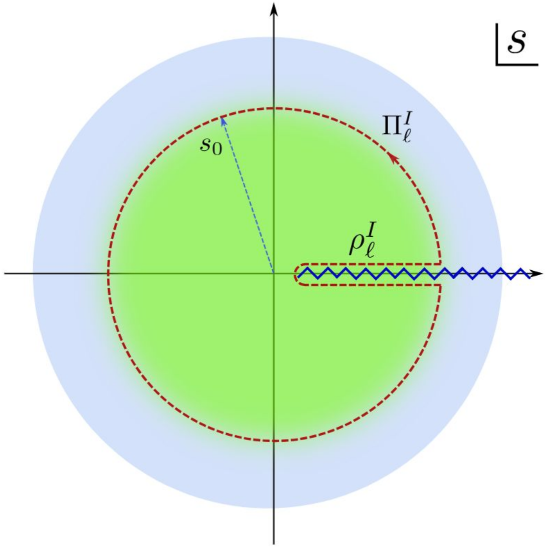
SVZ

$$\int_4^{s_0} \rho(x)x^n dx = -s_0^{n+1}\int_0^{2\pi} e^{i(n+1)\varphi}\Pi(s_0 e^{i\varphi})d\varphi$$

bootstrap variables ← **gauge theory information**

linear constraints

Finite energy sum rule



connect pQCD with bootstrap at s_0

contour integral $s^n\Pi(s)$ vanishes

SVZ

$$\int_4^{s_0} \rho(x) x^n dx = -s_0^{n+1} \int_0^{2\pi} e^{i(n+1)\varphi} \Pi(s_0 e^{i\varphi}) d\varphi$$

bootstrap variables ← **gauge theory information**

linear constraints

$$P1 \quad : \quad \frac{1}{s_0^{n+2}} \int_4^{s_0} \rho_1^1(x) x^n dx = \frac{1}{2(2\pi)^4} \left\{ \frac{1}{2\pi(n+2)} \left(1 + \frac{\alpha_s}{\pi}\right) - \frac{\delta_n \pi}{6 s_0^2} \langle \frac{\alpha_s}{\pi} G^2 \rangle - \frac{\delta_n 2\pi}{s_0^2} \langle m_q \bar{q} q \rangle + \dots \right\}, \quad n \geq -1$$

to extract from bootstrap in the future

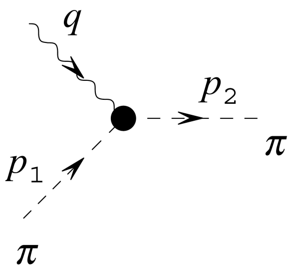
condensates suppressed at large s_0, not used as input

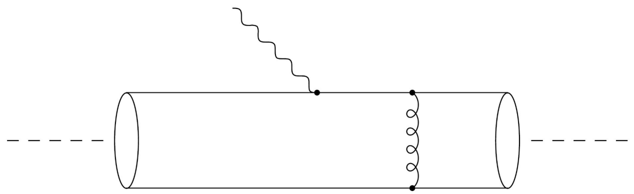
Asymptotic behavior of form factor from pQCD

perturbative QCD also controls asymptotic behavior of form factors

e.g. electromagnetic FF $\langle \pi(p_2) | J^\mu_{\rm em}(0) | \pi(p_1) \rangle = (p_1^\mu + p_2^\mu) F_\pi(q^2)$

$$q = (p_2 - p_1)$$





$$F_\pi(s) \simeq -\frac{16\pi\alpha_s(s) f_\pi^2}{s}$$

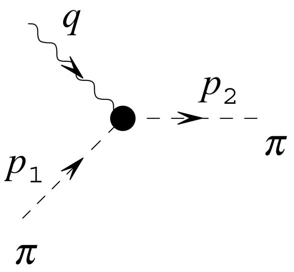
[Lepage, Brodsky, 1979]

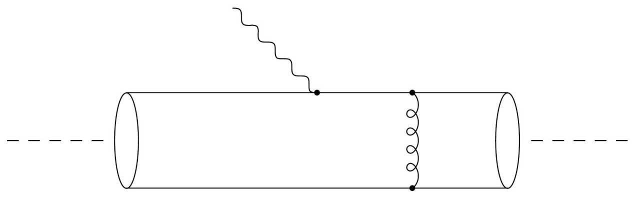
Asymptotic behavior of form factor from pQCD

perturbative QCD also controls asymptotic behavior of form factors

e.g. electromagnetic FF $\langle\pi(p_2)|J^{\mu}_{\text{em}}(0)|\pi(p_1)\rangle = (p_1^{\mu} + p_2^{\mu})F_{\pi}(q^2)$

$$q = (p_2 - p_1)$$





$$F_{\pi}(s) \simeq -\frac{16\pi\alpha_s(s)f_{\pi}^2}{s}$$

[Lepage, Brodsky, 1979]

evaluate to estimate

in practical numerical implementation
suffices to require smallness above $s = s_0$

$$||\mathcal{F}(s > s_0)|| \lesssim \epsilon$$

Gauge theory parameters: numerical input

$N_f = 2 \qquad N_c = 3$ *for comparison with experiments*

$$s_0 = (1.2\,\mathrm{GeV})^2, \quad \alpha_s \simeq 0.41, \quad m_u \simeq 4\,\mathrm{MeV} \quad m_d \simeq 7.3\,\mathrm{MeV}$$

$$s_0 = (2\,\mathrm{GeV})^2, \quad \alpha_s \simeq 0.31, \quad m_u \simeq 3.6\,\mathrm{MeV} \quad m_d \simeq 6.5\,\mathrm{MeV}$$

Gauge theory parameters: numerical input

$N_f = 2 \qquad N_c = 3$ *for comparison with experiments*

$$s_0 = (1.2\,\mathrm{GeV})^2, \quad \alpha_s \simeq 0.41, \quad m_u \simeq 4\,\mathrm{MeV} \quad m_d \simeq 7.3\,\mathrm{MeV}$$

$$s_0 = (2\,\mathrm{GeV})^2, \quad \alpha_s \simeq 0.31, \quad m_u \simeq 3.6\,\mathrm{MeV} \quad m_d \simeq 6.5\,\mathrm{MeV}$$

FESR

$$\frac{1}{s_0^{n+2}} \int_4^{s_0} \rho_0^0(x) x^n dx \simeq \frac{6.23 \times 10^{-7}}{n+2}$$

$$\frac{1}{s_0^{n+2}} \int_4^{s_0} \rho_1^1(x) x^n dx \simeq \frac{5.62 \times 10^{-5}}{n+2}$$

$$\frac{1}{s_0^{n+3}} \int_4^{s_0} \rho_2^0(x) x^n dx \simeq \frac{5.13 \times 10^{-5}}{n+3}$$

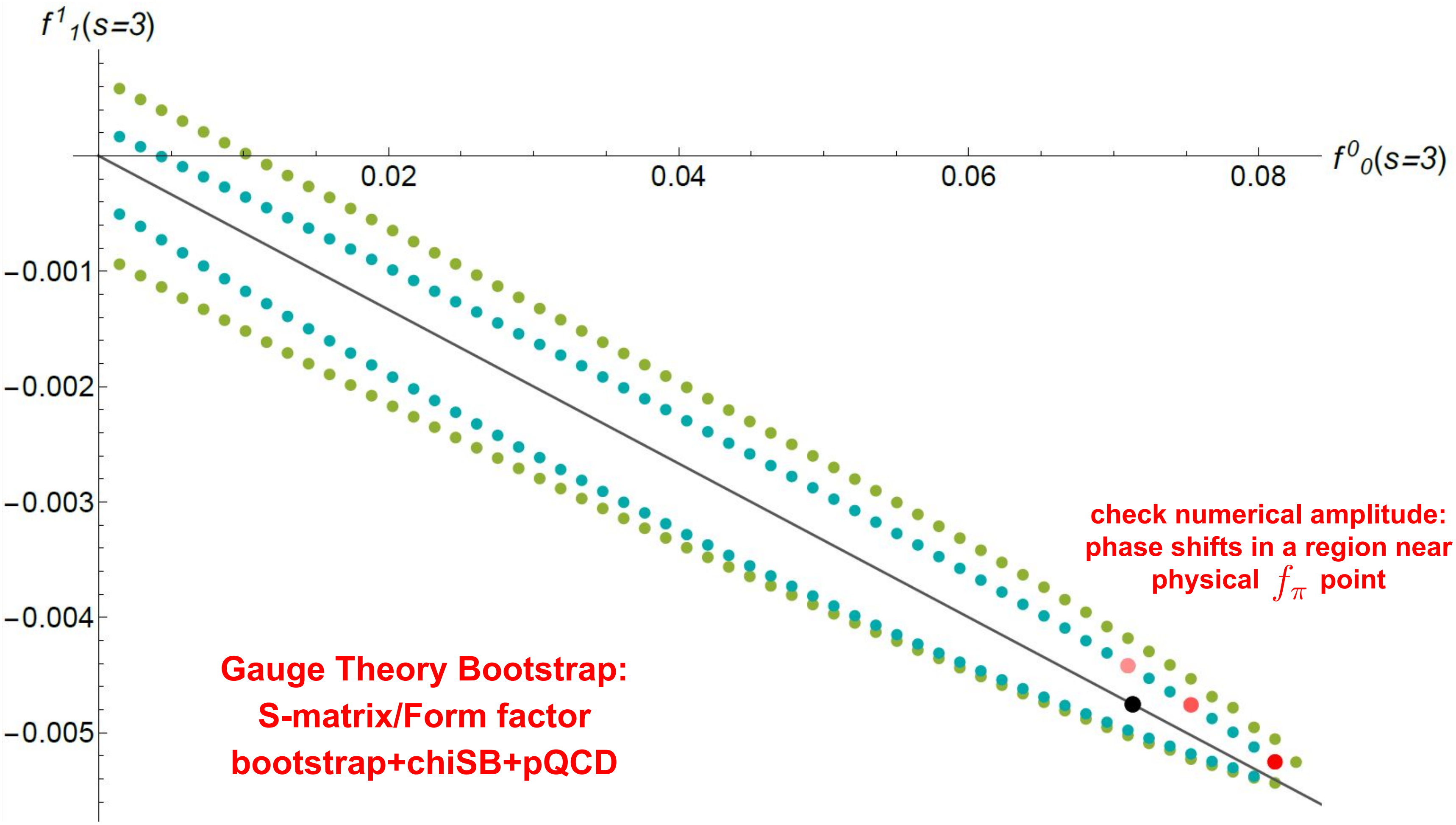
FF asymptotics

$$|\mathcal{F}_0^0(s > s_0)|^2 \lesssim 3 \times 10^{-8}$$

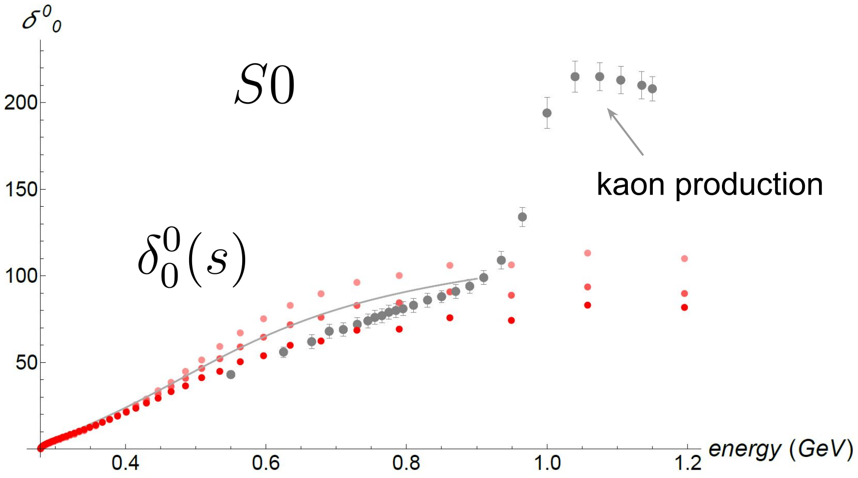
$$|\mathcal{F}_1^1(s > s_0)|^2 \lesssim 2 \times 10^{-6}$$

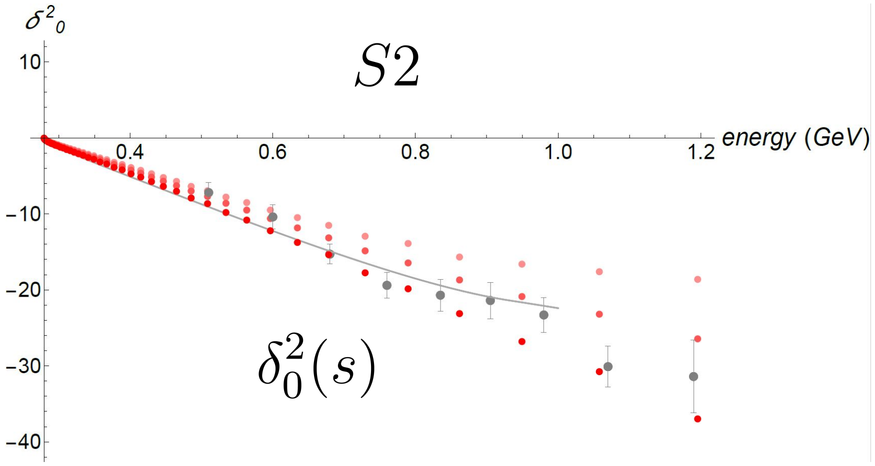
$$|\mathcal{F}_2^0(s > s_0)|^2 \lesssim 4 \times 10^{-2}$$

...



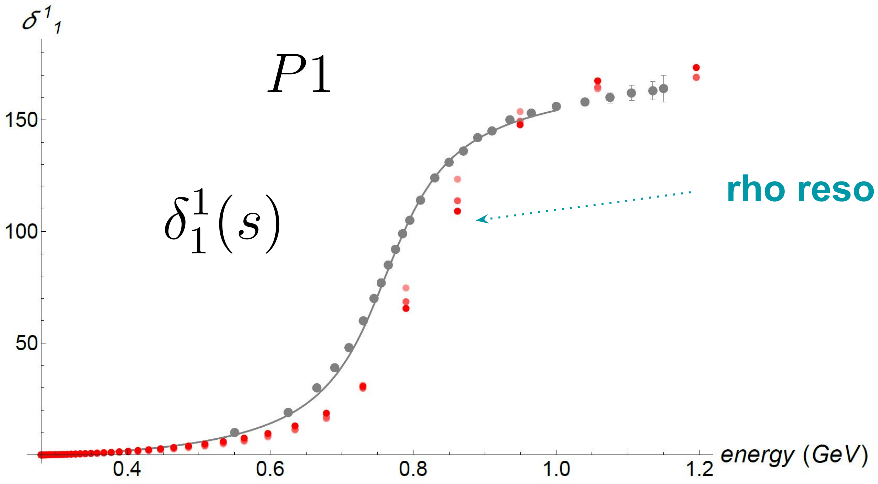
$f^1{}_1(s=3)$
$f^0{}_0(s=3)$
0.02
0.04
0.06
0.08
−0.001
−0.002
−0.003
−0.004
−0.005
check numerical amplitude:
phase shifts in a region near
physical f_π point
Gauge Theory Bootstrap:
S-matrix/Form factor
bootstrap+chiSB+pQCD





phase shifts up to 1.2GeV

Gauge Theory Bootstrap

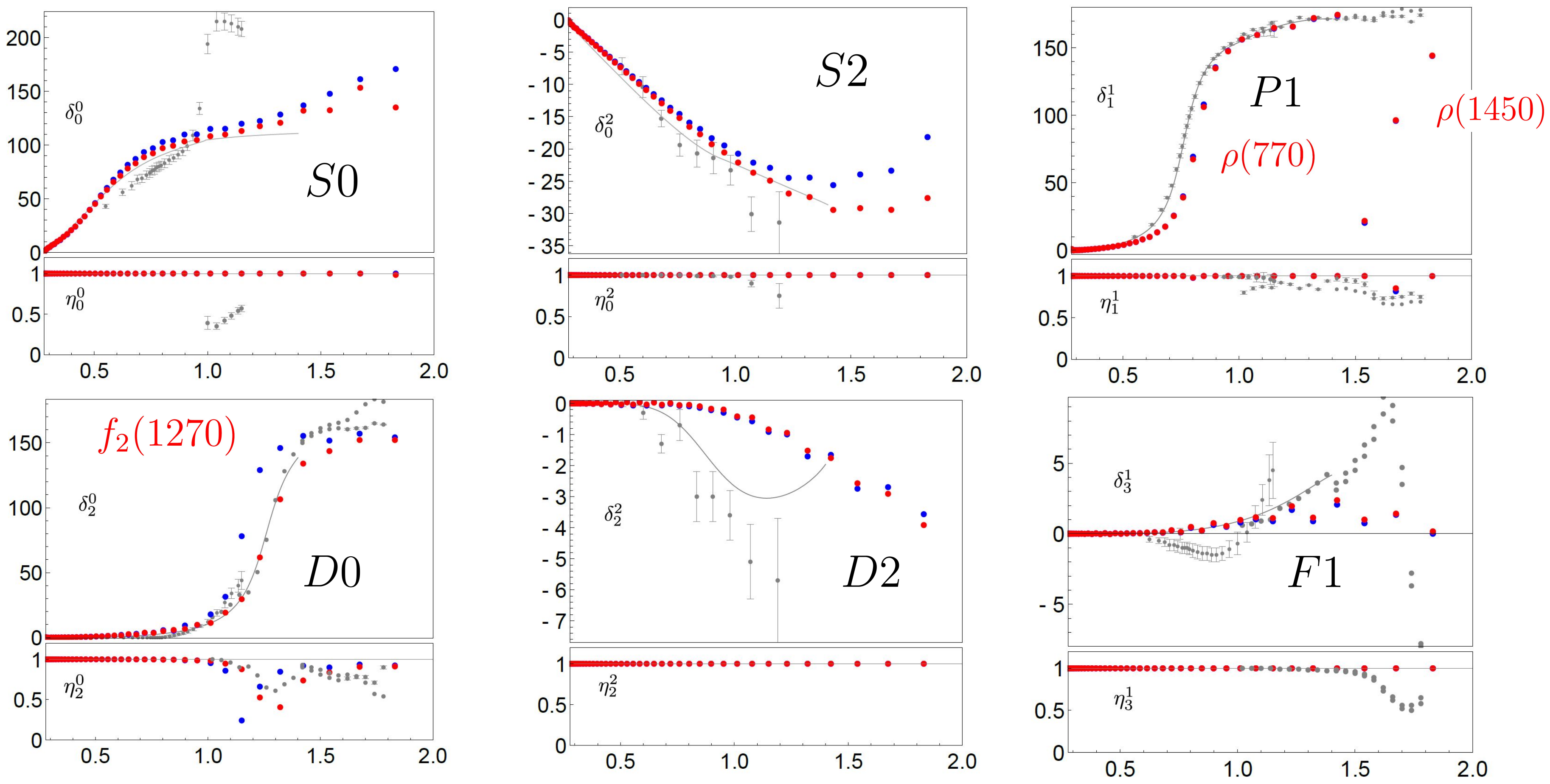


rho resonance, info from UV

experimental data (gray dots)
[Protopopescu et al, 1973]
[Losty et al, 1974]
pheno fit (gray line)
[Pelaez, Yndurain, 2005]

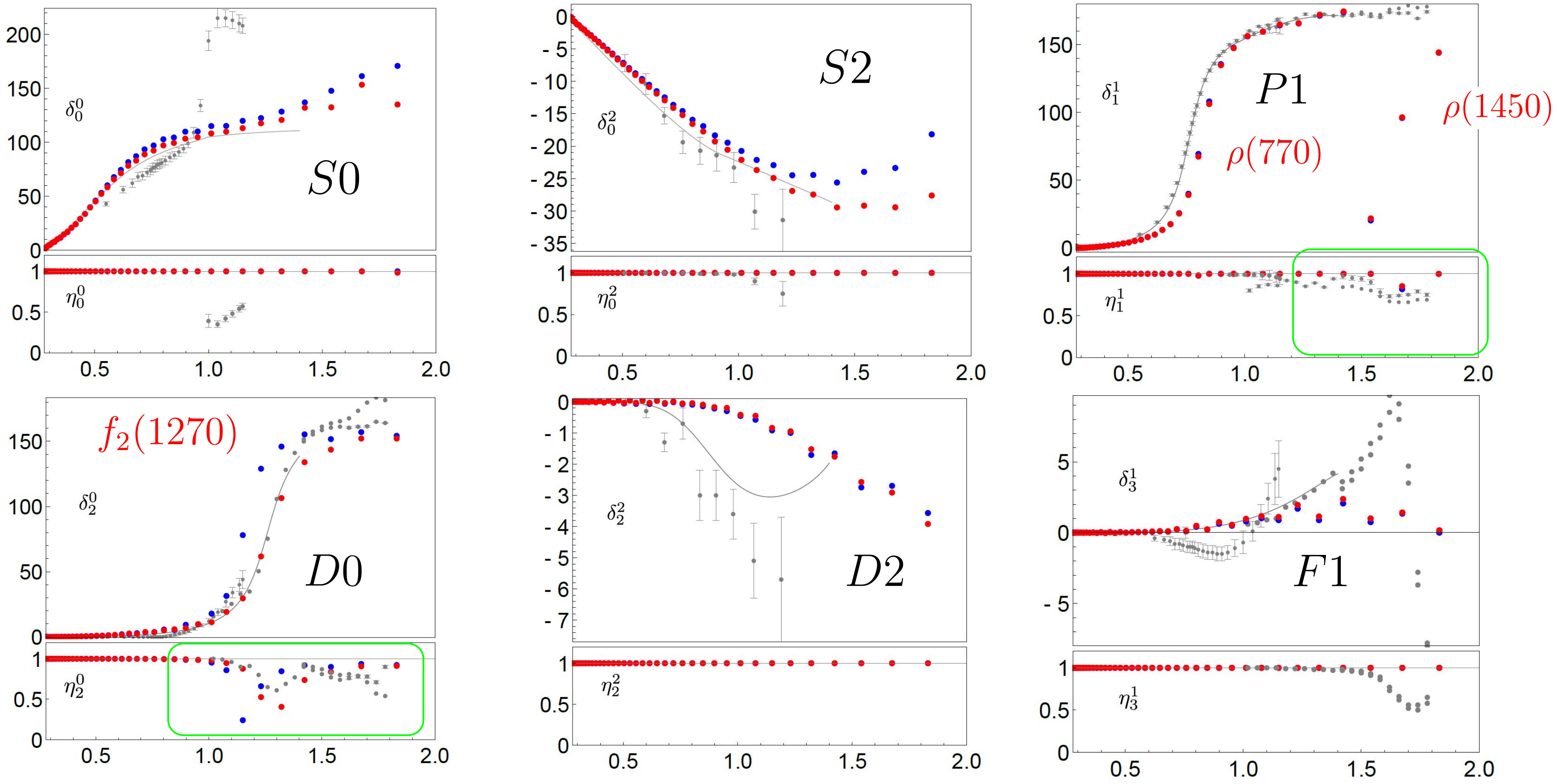
Gauge Theory Bootstrap

phase shifts up to 2GeV



Gauge Theory Bootstrap

phase shifts up to 2GeV



Low energy parameters: threshold expansion

scattering lengths and effective range parameters

$$\mathrm{Re} f_\ell^I(s) \overset{k\to 0}{\simeq} \frac{2m_\pi}{\pi} k^{2\ell}\left(a_\ell^I + b_\ell^I k^2 + \dots\right) \qquad k = \frac{\sqrt{s-4m_\pi^2}}{2}$$

	W	**GTB**	CGL	PY
$a_0^{(0)}$	0.16	0.178, 0.182	0.220 ± 0.005	0.230 ± 0.010
$a_0^{(2)}$	-0.046	-0.0369, -0.0378	-0.0444 ± 0.0010	-0.0422 ± 0.0022
$b_0^{(0)}$	0.18	0.287, 0.290	0.280 ± 0.001	0.268 ± 0.010
$b_0^{(2)}$	-0.092	-0.064, -0.066	-0.080 ± 0.001	-0.071 ± 0.004
$a_1^{(1)}$	31	28.0, 28.4	37.0 ± 0.13	$38.1 \pm 1.4\ (\times 10^{-3})$
$b_1^{(1)}$	0	2.86, 3.37	5.67 ± 0.13	$4.75 \pm 0.16\ (\times 10^{-3})$
$a_2^{(0)}$	0	12.6, 12.3	17.5 ± 0.3	$18.0 \pm 0.2\ (\times 10^{-4})$
$a_2^{(2)}$	0	2.87, 2.81	1.70 ± 0.13	$2.2 \pm 0.2\ (\times 10^{-4})$

Low energy parameters: pion charge radii

threshold expansion of the form factors:

scalar form factor: $$F_0^0(s) = F_0^0(0)\left[1 + \frac{1}{6}s\langle r^2\rangle_S^\pi + \ldots\right]$$

vector form factor: $$F_1^1(s) = 1 + \frac{1}{6}s\langle r^2\rangle_V^\pi + \ldots$$

	GTB	Exp. fits
$\langle r^2\rangle_S^\pi$	0.64, 0.61	$0.61 \pm 0.04\,\mathrm{fm}^2$
$\langle r^2\rangle_V^\pi$	0.388, 0.381	$0.439 \pm 0.008\,\mathrm{fm}^2$

Low energy parameters: chiral Lagrangian coefficients

calculate **the chiral Lagrangian coefficients**

$$\bar{\ell}_{1,2,4,6}$$

$$a_{D0} = \frac{1}{1440\pi^3 f_\pi^4}\left\{\bar{l}_1 + 4\bar{l}_2 - \frac{53}{8}\right\} + \ldots$$

$$a_{D2} = \frac{1}{1440\pi^3 f_\pi^4}\left\{\bar{l}_1 + \bar{l}_2 - \frac{103}{40}\right\} + \ldots$$

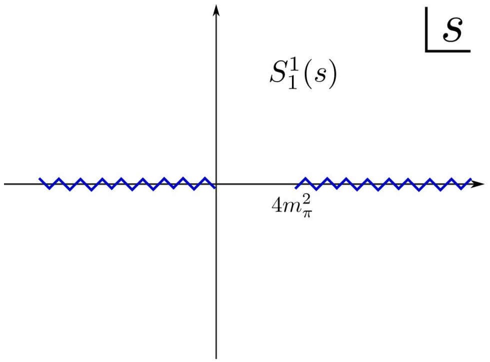
$$F_0(s) = 1 + \frac{s}{16\pi^2 f_\pi^2}\left(\bar{l}_4 - \frac{13}{12}\right) + \ldots$$

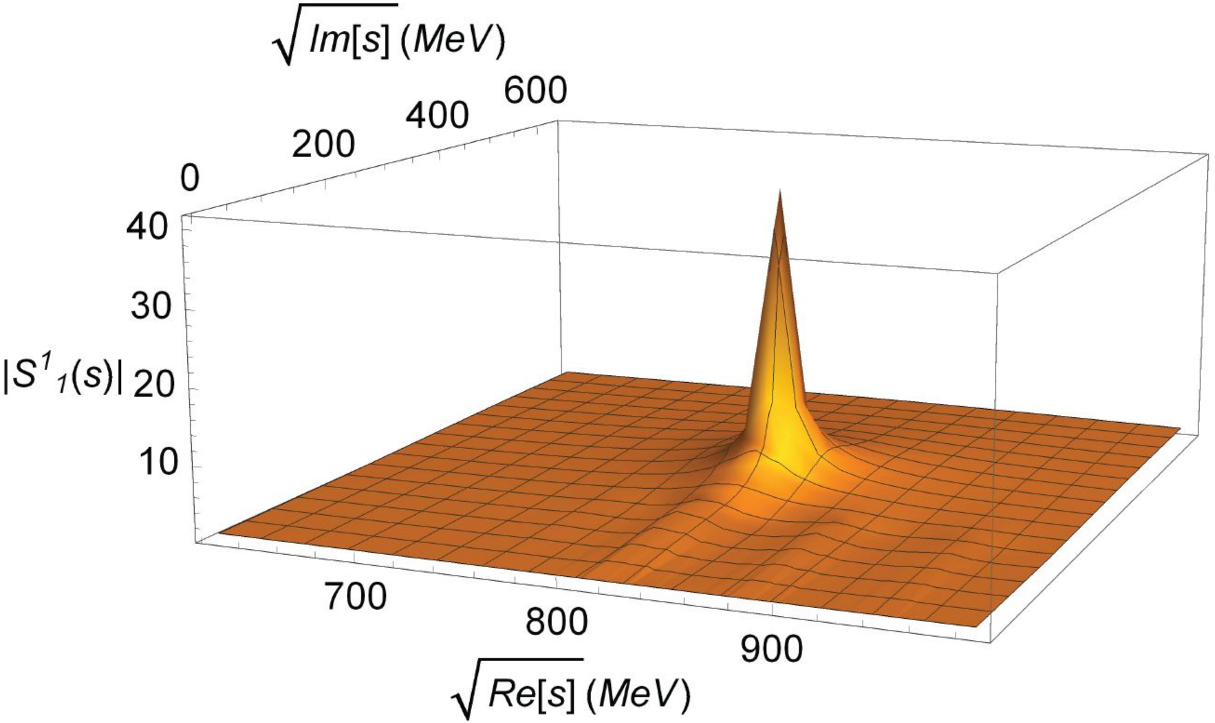
$$F_1(s) = 1 + \frac{s}{96\pi^2 f_\pi^2}(\bar{l}_6 - 1) + \ldots$$

[Gasser, Leutwyler, 1984]

	GTB	GL	Bij	CGL
$\bar{l}_1$	0.92, 0.93	-2.3 ± 3.7	-1.7 ± 1.0	-0.4 ± 0.6
$\bar{l}_2$	4.1, 4.0	6.0 ± 1.3	6.1 ± 0.5	4.3 ± 0.1
$\bar{l}_4$	4.7, 4.6	4.3 ± 0.9	4.4 ± 0.3	4.4 ± 0.2
$\bar{l}_6$	14.3, 14.1	16.5 ± 1.1	$16.0 \pm 0.5 \pm 0.7$	

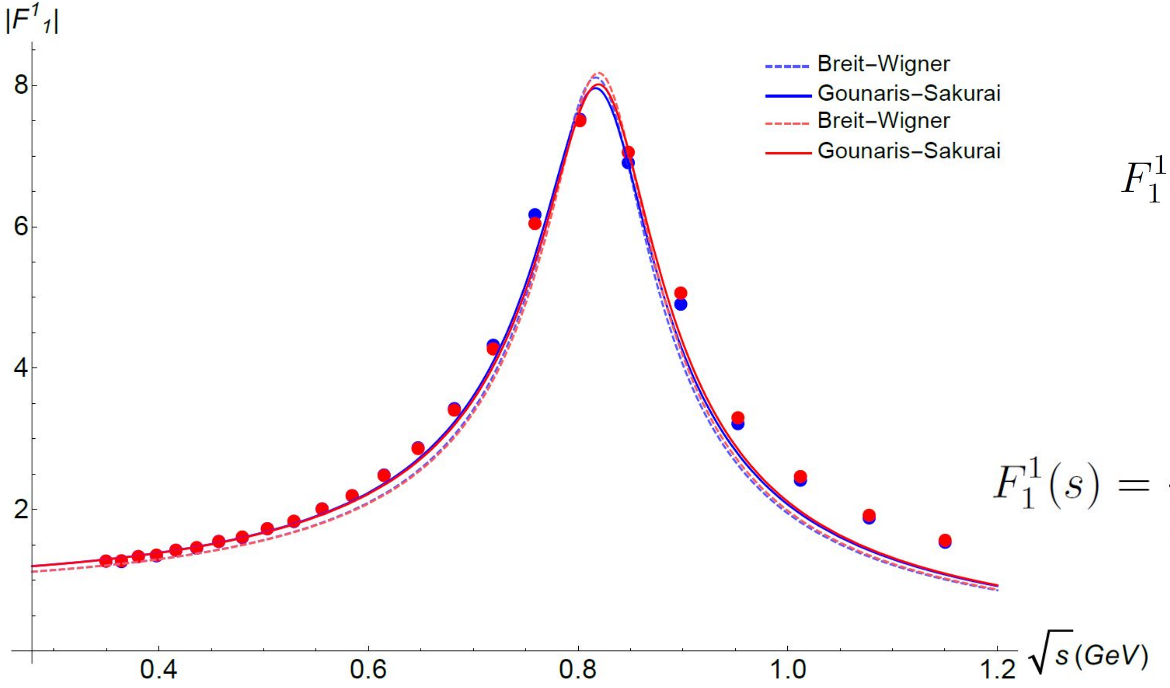
rho meson as pole on the second sheet of $S_1^1(s)$





	GTB	PDG
$\mathrm{Re}(\sqrt{s_\rho})$	829, 832	$761 - 765 \pm 0.23$ MeV
$\mathrm{Im}(\sqrt{s_\rho})$	63, 64	$71 - 74 \pm 0.8$ MeV

Vector (electromagnetic) form factor and rho meson



Breit-Wigner form

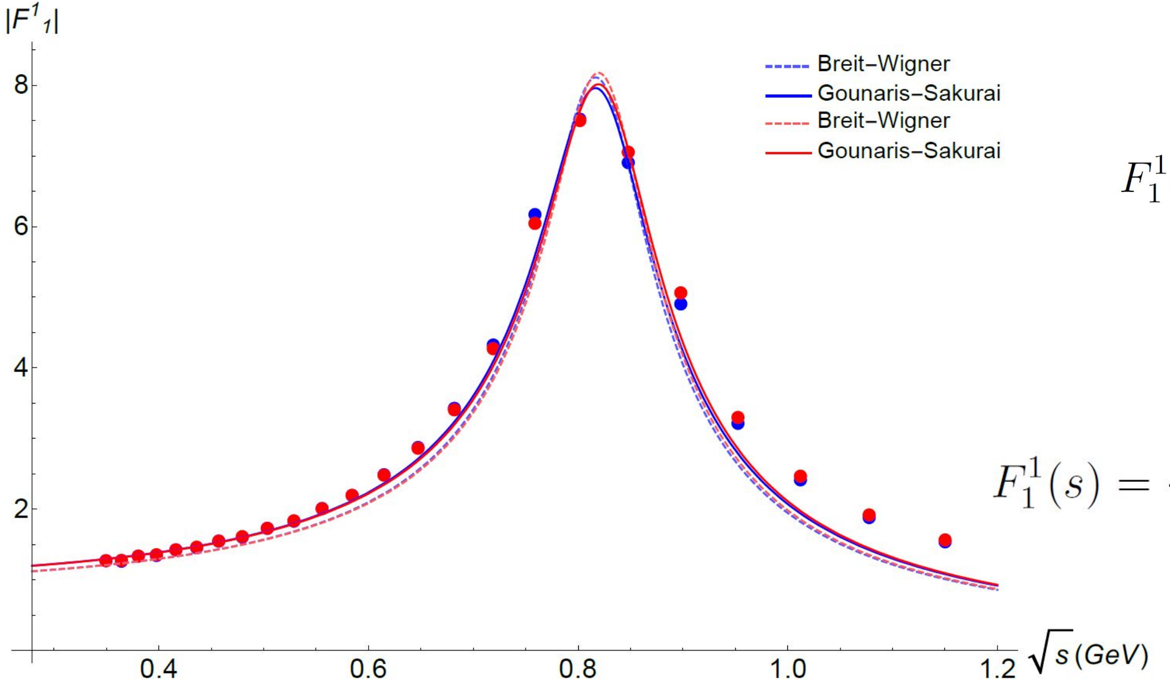
$$F_1^1(s) = -\frac{m_\rho^2}{s - m_\rho^2 + i m_\rho \Gamma_\rho \theta(s - 4m_\pi^2)}$$

Gounaris-Sakurai form

$$F_1^1(s) = \frac{m_\rho^2[1 + d\Gamma_\rho/m_\rho]}{(m_\rho^2 - s) - i m_\rho \Gamma_\rho (q/q_\rho)^3 (m_\rho/\sqrt{s})}$$

	GTB	PDG
m_ρ	836, 839	775 ± 0.23 MeV
Γ_ρ	111, 111	149.1 ± 0.8 MeV

Vector (electromagnetic) form factor and rho meson



Breit-Wigner form

$$F_1^1(s) = -\frac{m_\rho^2}{s - m_\rho^2 + i m_\rho \Gamma_\rho \theta(s - 4m_\pi^2)}$$

Gounaris-Sakurai form

$$F_1^1(s) = \frac{m_\rho^2[1 + d\Gamma_\rho/m_\rho]}{(m_\rho^2 - s) - i m_\rho \Gamma_\rho (q/q_\rho)^3 (m_\rho/\sqrt{s})}$$

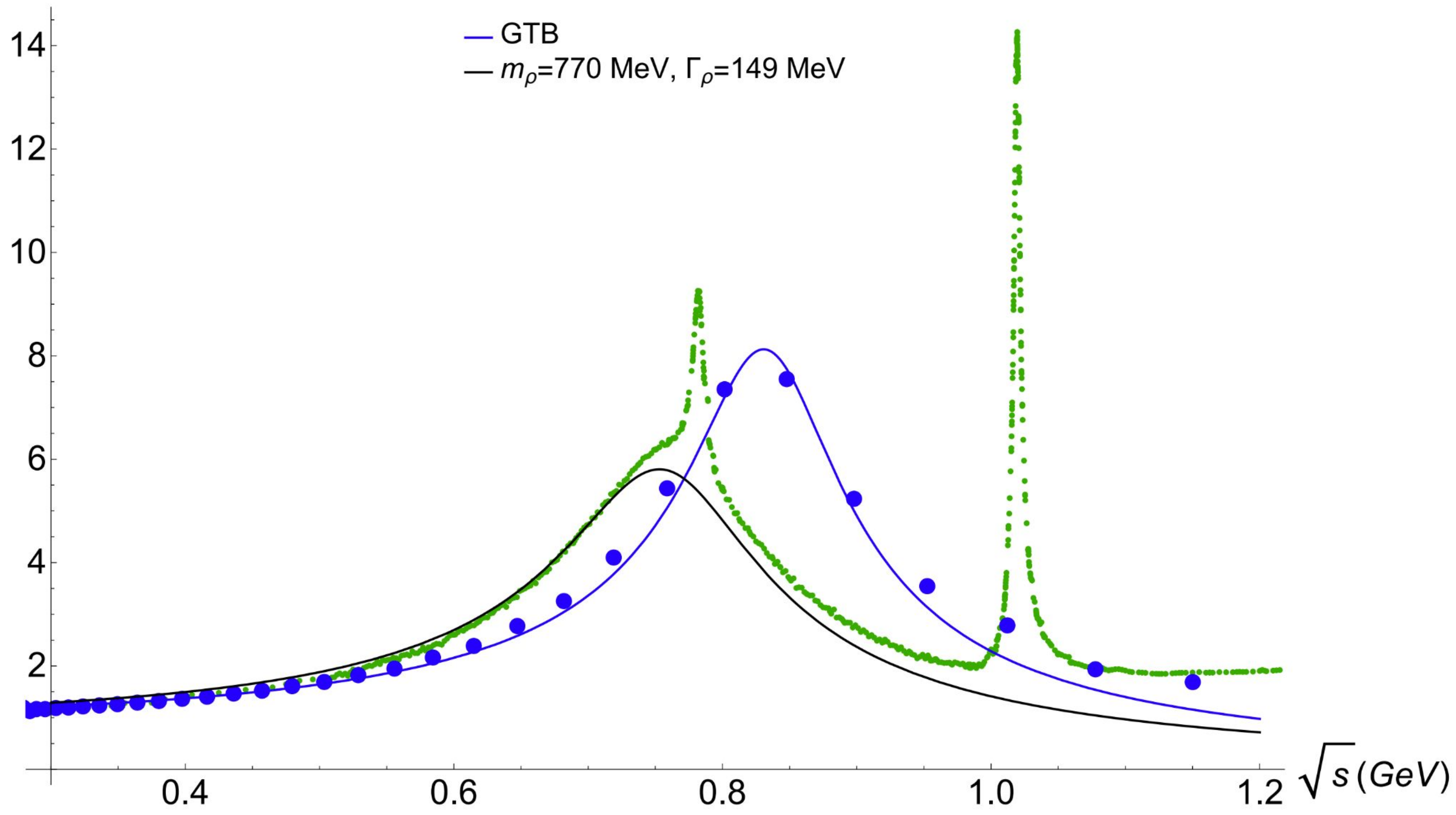
couplings $\Gamma_\rho = g_{\rho\pi\pi}^2 \frac{m_\rho}{48\pi}\left[1 - \frac{4m_\pi^2}{m_\rho^2}\right]^{\frac{3}{2}}$

$g_{\rho\pi\pi} = 4.9, \quad 4.9$

$g_{\rho\pi\pi} = 6$

	GTB	PDG
m_ρ	836, 839	775 ± 0.23 MeV
Γ_ρ	111, 111	149.1 ± 0.8 MeV

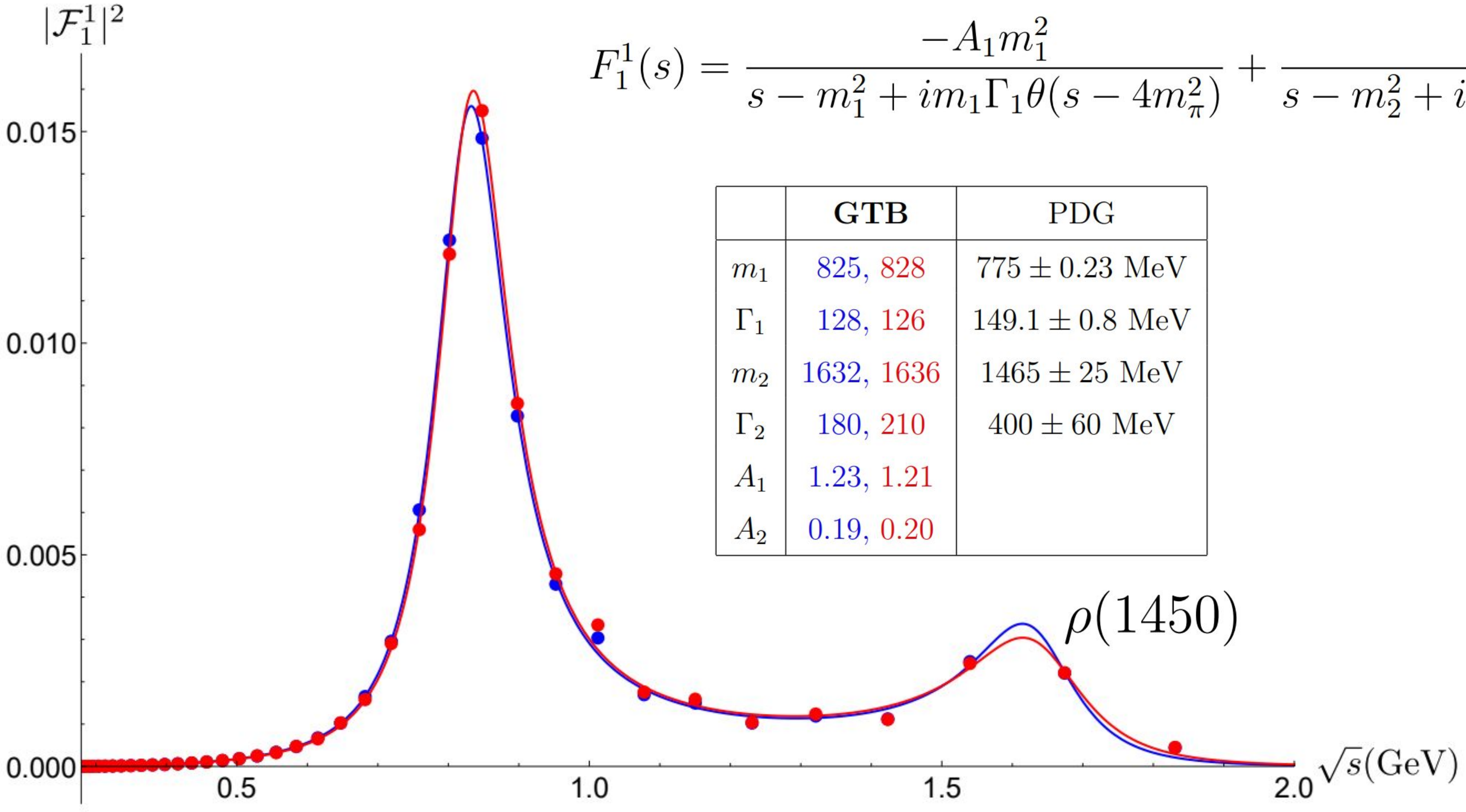
$e^+e^- \longrightarrow$ hadrons



Vector (electromagnetic) form factor and rho meson

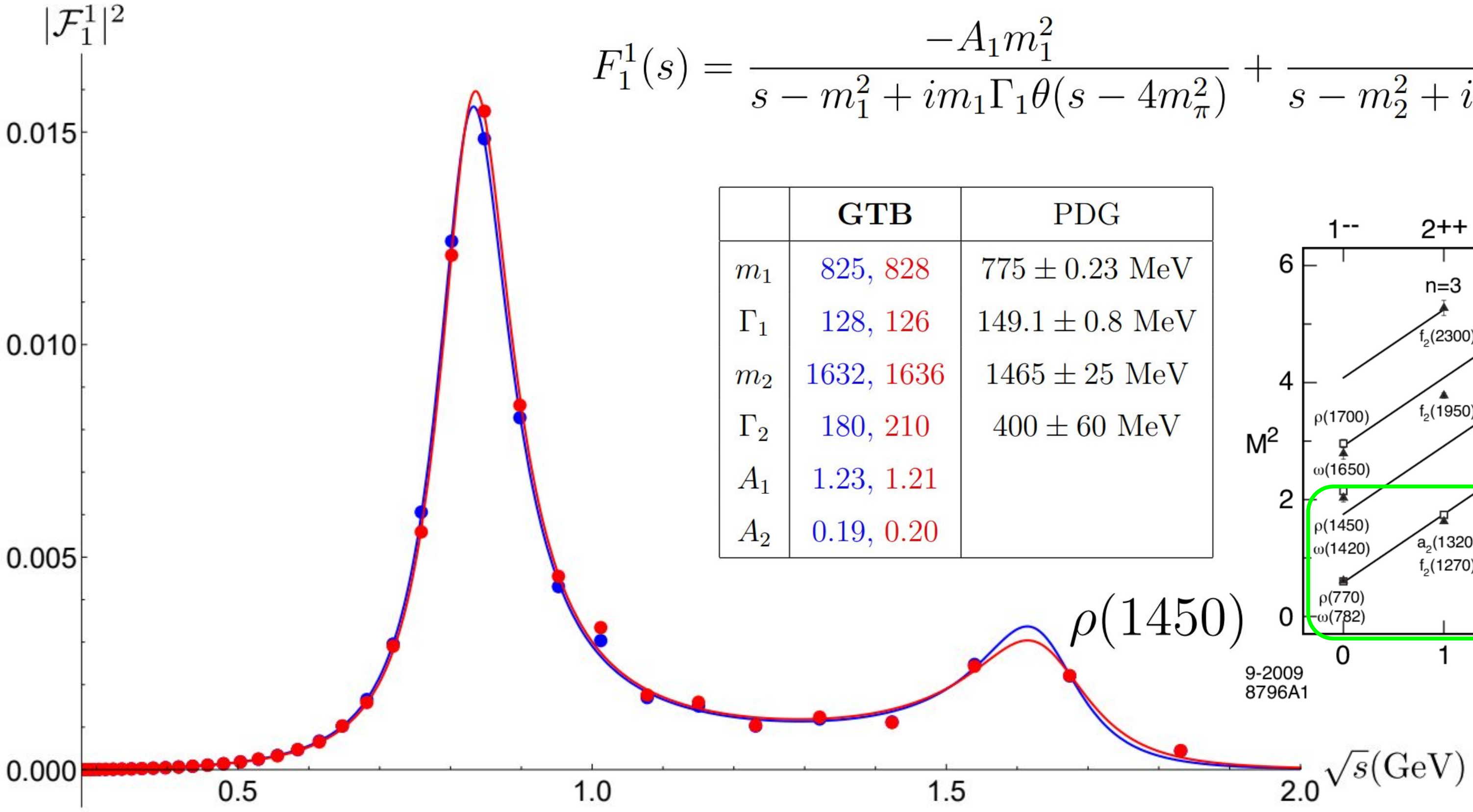
$$F_1^1(s) = \frac{-A_1 m_1^2}{s - m_1^2 + i m_1 \Gamma_1 \theta(s - 4m_\pi^2)} + \frac{A_2 m_2^2}{s - m_2^2 + i m_2 \Gamma_2 \theta(s - 4m_\pi^2)}$$

	GTB	PDG
m_1	825, 828	775 ± 0.23 MeV
Γ_1	128, 126	149.1 ± 0.8 MeV
m_2	1632, 1636	1465 ± 25 MeV
Γ_2	180, 210	400 ± 60 MeV
A_1	1.23, 1.21	
A_2	0.19, 0.20	

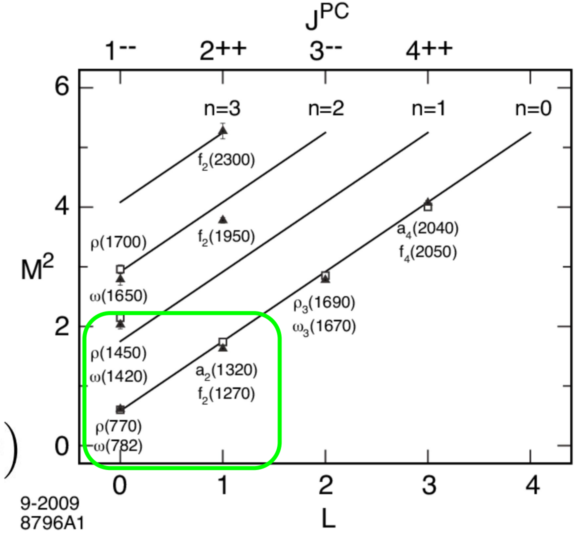


Vector (electromagnetic) form factor and rho meson

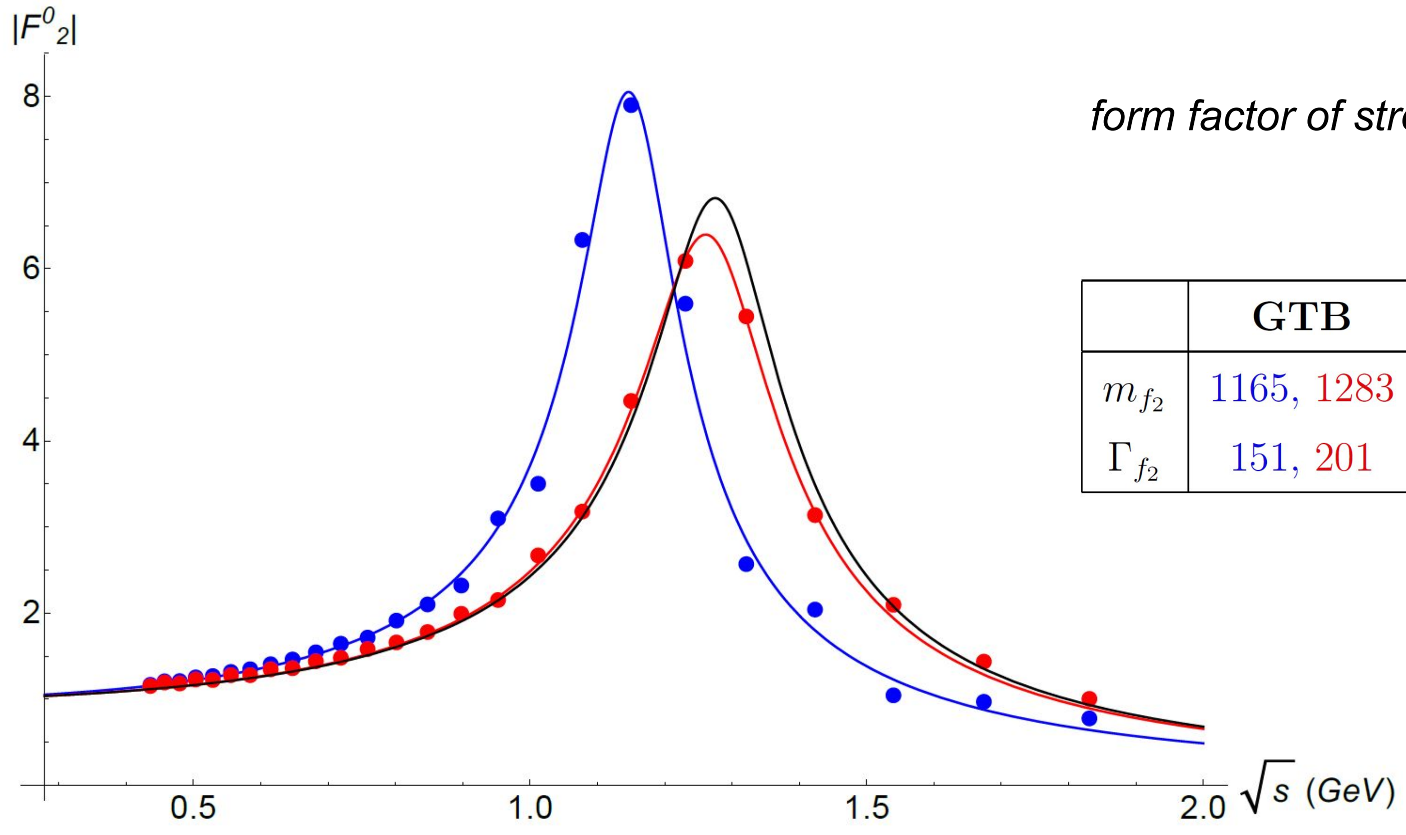
$$F_1^1(s) = \frac{-A_1 m_1^2}{s - m_1^2 + i m_1 \Gamma_1 \theta(s - 4m_\pi^2)} + \frac{A_2 m_2^2}{s - m_2^2 + i m_2 \Gamma_2 \theta(s - 4m_\pi^2)}$$



	GTB	PDG
m_1	825, 828	775 ± 0.23 MeV
Γ_1	128, 126	149.1 ± 0.8 MeV
m_2	1632, 1636	1465 ± 25 MeV
Γ_2	180, 210	400 ± 60 MeV
A_1	1.23, 1.21	
A_2	0.19, 0.20	



Gravitational form factor and f_2 meson



form factor of stress energy tensor

	GTB	PDG
m_{f_2}	1165, 1283	1275.4 ± 0.6 MeV
Γ_{f_2}	151, 201	186.6 ± 2.3 MeV

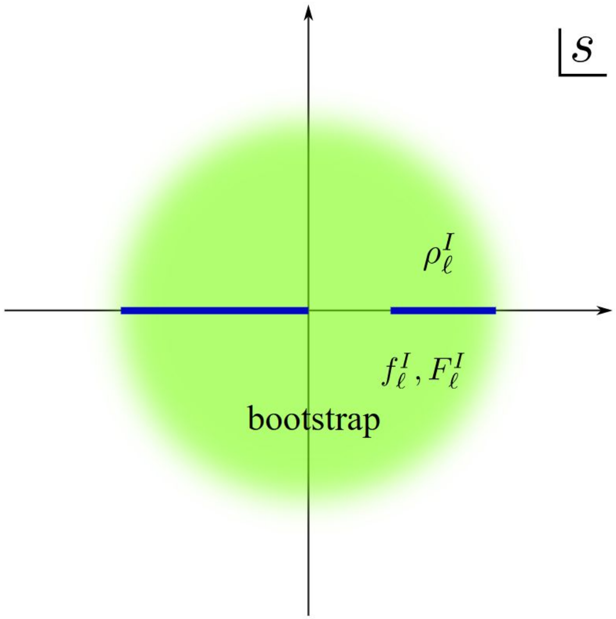
Saturation of positive semidefinite matrix

positive semidefinite
$$\begin{pmatrix} 1 & S & \mathcal{F} \\ S^* & 1 & \mathcal{F}^* \\ \mathcal{F}^* & \mathcal{F} & \rho \end{pmatrix} \succeq 0 \qquad \forall I,\ \ell,\ s$$

iff all its principal minors are non-negative

$$\rho + S^*\mathcal{F}^2 + S(\mathcal{F}^*)^2 - 2|\mathcal{F}|^2 - \rho|S|^2 \geq 0$$

$$\rho \geq 0 \qquad |\mathcal{F}|^2 \leq \rho \qquad |S|^2 \leq 1$$



Saturation of positive semidefinite matrix

positive semidefinite
$$\begin{pmatrix} 1 & S & \mathcal{F} \\ S^* & 1 & \mathcal{F}^* \\ \mathcal{F}^* & \mathcal{F} & \rho \end{pmatrix} \succeq 0 \qquad \forall I,\ \ell,\ s$$

iff all its principal minors are non-negative

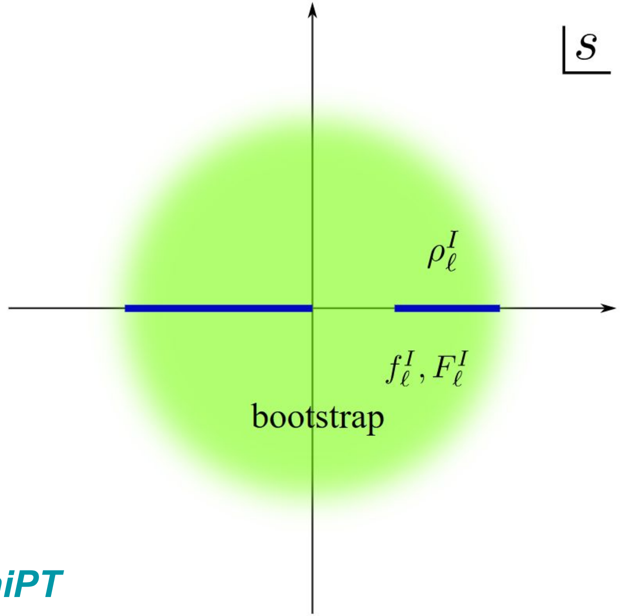
$$\rho + S^*\mathcal{F}^2 + S(\mathcal{F}^*)^2 - 2|\mathcal{F}|^2 - \rho|S|^2 \geq 0$$

$$\rho \geq 0 \qquad |\mathcal{F}|^2 \leq \rho \qquad |S|^2 \leq 1$$

saturation: $\rho = |\mathcal{F}|^2$

$$|S| = 1 \quad S = \frac{\mathcal{F}}{\mathcal{F}^*}$$

Watson theorem



saturation connects quantities controlled by pQCD and chiPT

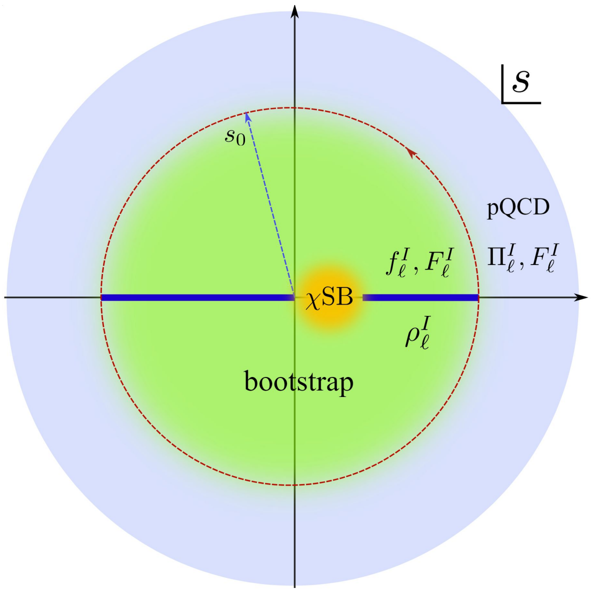
How the Gauge Theory Bootstrap works

energy

gauge theory info

$$\int_4^{s_0} \rho(x) x^n dx = -s_0^{n+1} \int_0^{2\pi} e^{i(n+1)\varphi} \Pi(s_0 e^{i\varphi}) d\varphi \quad F_\pi(s) \simeq -\frac{16\pi\alpha_s(s) f_\pi^2}{s} \quad , \dots$$

pQCD



very low energy behavior $A(s,t,u) \simeq \frac{4}{\pi} \frac{s - m_\pi^2}{32\pi f_\pi^2}$ **chiSB**

How the Gauge Theory Bootstrap works

energy

gauge theory info

$$\int_4^{s_0} \rho(x) x^n dx = -s_0^{n+1} \int_0^{2\pi} e^{i(n+1)\varphi} \Pi(s_0 e^{i\varphi}) d\varphi \quad F_\pi(s) \simeq -\frac{16\pi \alpha_s(s) f_\pi^2}{s} \ , \ \dots$$

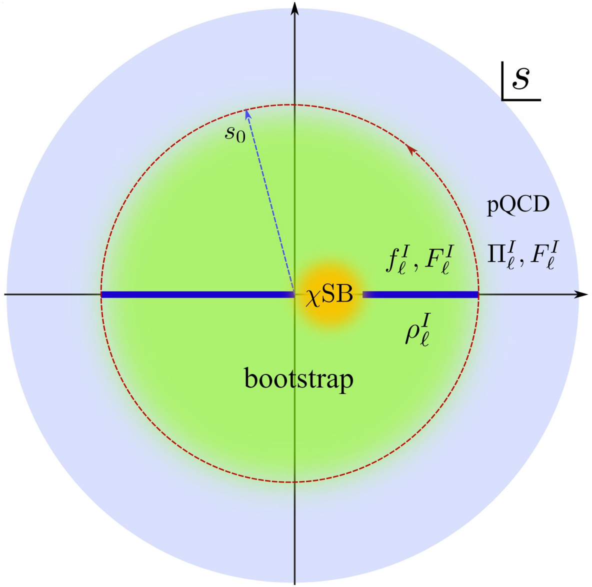
pQCD

$$\rho(s) = |\mathcal{F}(s)|^2$$

dispersion relation

$$\mathcal{F}(s) = \sqrt{\rho(s)} e^{i\alpha(s)} \quad \ln \mathcal{F}(s) = \frac{1}{2} \ln \rho(s) + i\alpha(s)$$

very low energy behavior $A(s,t,u) \simeq \frac{4}{\pi} \frac{s - m_\pi^2}{32\pi f_\pi^2}$ **chiSB**



How the Gauge Theory Bootstrap works

energy

gauge theory info

$$\int_4^{s_0} \rho(x) x^n dx = -s_0^{n+1} \int_0^{2\pi} e^{i(n+1)\varphi} \Pi(s_0 e^{i\varphi}) d\varphi \quad F_\pi(s) \simeq -\frac{16\pi\alpha_s(s) f_\pi^2}{s}, \ \dots$$

pQCD

Gauge Theory Bootstrap

$$\rho(s) = |\mathcal{F}(s)|^2$$

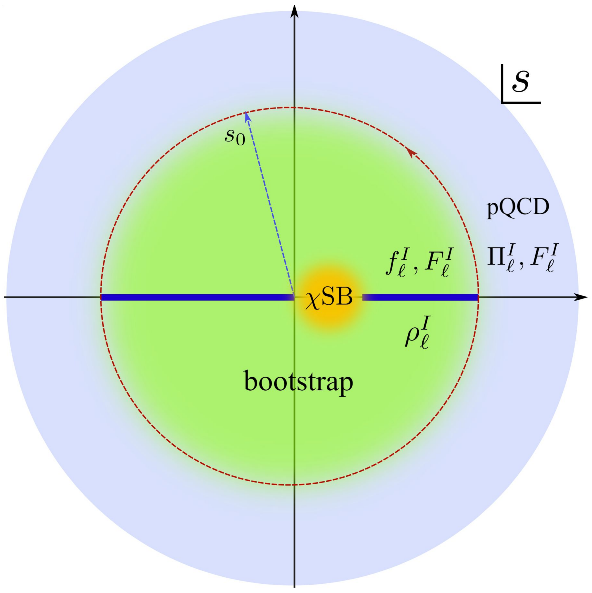
dispersion relation

$$\mathcal{F}(s) = \sqrt{\rho(s)} e^{i\alpha(s)} \quad \ln \mathcal{F}(s) = \frac{1}{2} \ln \rho(s) + i\alpha(s)$$

$$|S(s)| = 1 \qquad e^{2i\alpha(s)} = \frac{\mathcal{F}(s)}{\mathcal{F}^*(s)} = S(s) = e^{2i\delta(s)}$$

very low energy behavior $\quad A(s,t,u) \simeq \frac{4}{\pi} \frac{s - m_\pi^2}{32\pi f_\pi^2}$

chiSB



Conclusions

- Gauge Theory Bootstrap:

using only $N_c \; N_f \; m_q \; \Lambda_{\text{QCD}}$ $f_\pi \; m_\pi$ to remove in the future development

gauge theory parameters ($N_c \; N_f \; m_q \; \Lambda_{\text{QCD}}$) *universal low energy parameters* ($f_\pi \; m_\pi$)

strongly coupled low energy physics of asymptotically free gauge theories

Conclusions

- Gauge Theory Bootstrap:

  using only $N_c \; N_f \; m_q \; \Lambda_{\mathrm{QCD}}$ (*gauge theory parameters*) $f_\pi \; m_\pi$ (*universal low energy parameters*) to remove in the future development

  strongly coupled low energy physics of asymptotically free gauge theories

- Numerical test with $N_f = 2 \; N_c = 3$ find good agreement with experiments

  Results suggest: we are on the right track for ***solving QCD*** (gauge theories)

Prospects

- many future explorations in the framework:

  tuning gauge theory parameters → low energy dynamics

  (hadron spectrum, couplings)

  interplay between gauge theory vs. chiral dynamics (e.g. S0 vs. P1)

Prospects

- many future explorations in the framework:

  tuning gauge theory parameters → low energy dynamics

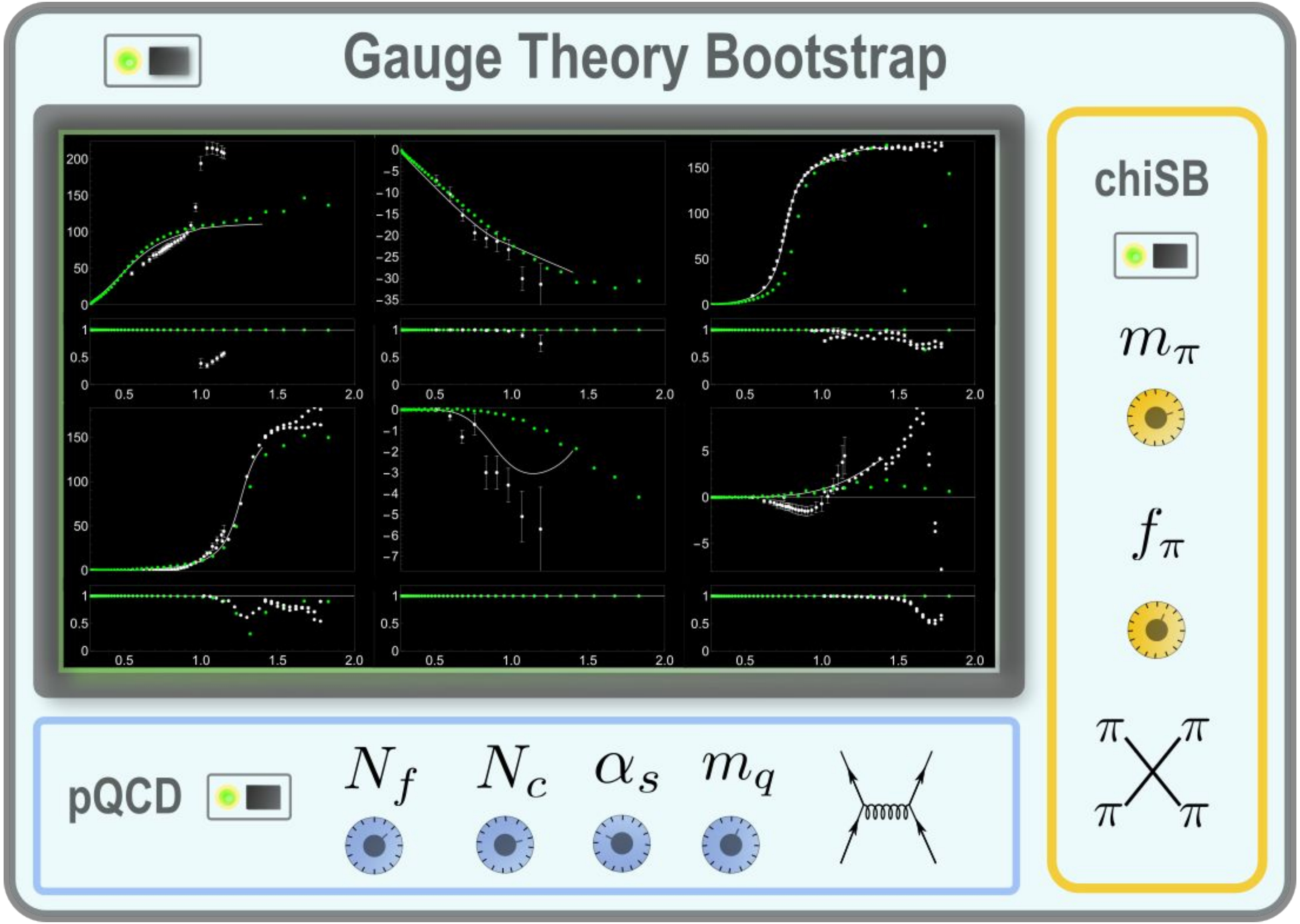
  (hadron spectrum, couplings)

  interplay between gauge theory vs. chiral dynamics (e.g. S0 vs. P1)

- Fast machine precision numerics (~20min on average laptop),

  Need a lot of improvement to be more robust

Ancillary files (details):
- GTB_numerics.m
- GTB_numerics.nb



Gauge Theory Bootstrap
chiSB
m_π
f_π
pQCD
N_f
N_c
α_s
m_q

Thank you!