Automated calculation of N-jet soft functions

Bahman Dehnadi

In collaboration with
Guido Bell,Tobias Mohrmann and Rudi Rahn

UNIVERSITÄT
SIEGEN

Seminar talk, June 12th 2018, Vienna, Austria

Introduction

- Computation of jet cross sections beyond LO complicated by IR-divergences
- Subtraction techniques: subtract IR soft and collinear behaviors from the real emission, then add them back to virtual contributions to cancel the IR poles
- q_T subtraction: e.g. top quark production at hadron colliders

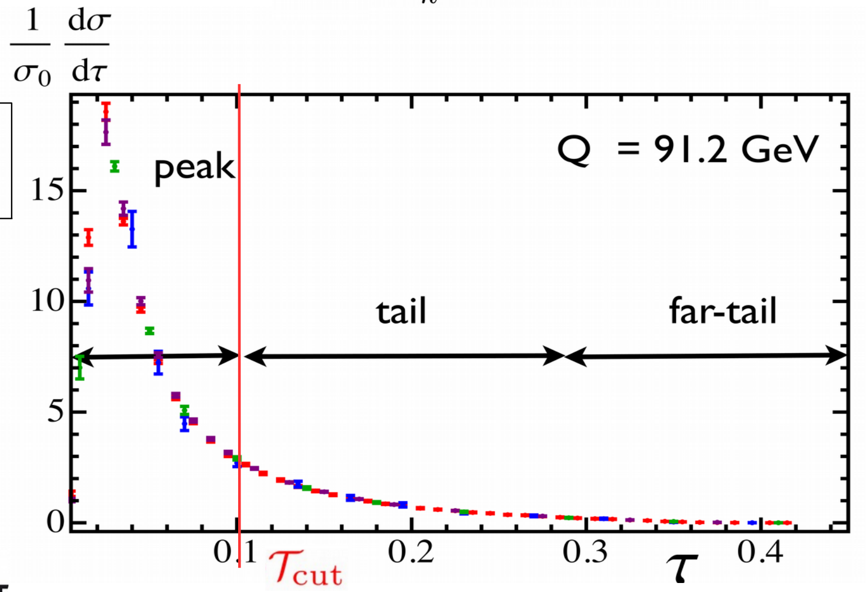
Catani, Grazzini (2007)
Bonciani,Catani, Grazzini, Sargsyan, Torre (2015)

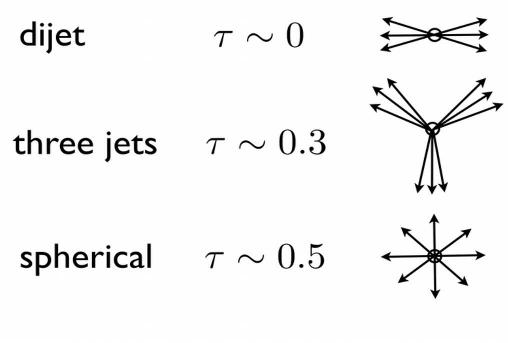
- N-jettiness slicing: e.g. H+jet, W+jet and Z+jet

Boughezal, Focke, Liu, Petriello(2015)
Gaunt, Stahlhofen, Tackmann, Walsh (2015)
Boughezal et. al. (2015)

N-jettiness variable: $$\mathcal{T}_N = \sum_k \min_i \{n_i \cdot p_k\}$$

e⁺e⁻ 2-jettiness (thrust)





Introduction

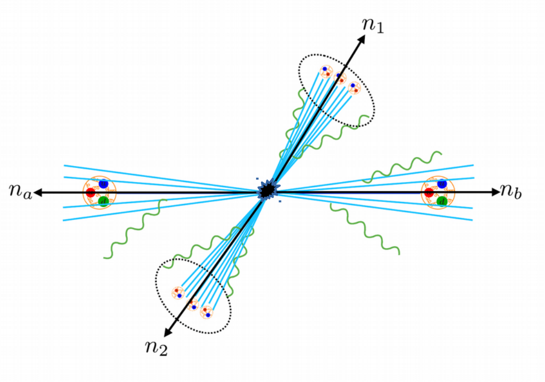
N-jet cross section:

Boughezal, Focke, Liu, Petriello(2015)
Gaunt, Stahlhofen, Tackmann, Walsh (2015)

$$\boxed{\sigma(X) = \int_0 \mathrm{d}\mathcal{T}_N \frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N} = \int_0^{\mathcal{T}_{\mathrm{cut}}} \mathrm{d}\mathcal{T}_N \frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N} + \int_{\mathcal{T}_{\mathrm{cut}}} \mathrm{d}\mathcal{T}_N \frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N}}$$

singular
N-jet final state

non-singualr
N+1-jet / multi-jets



Introduction

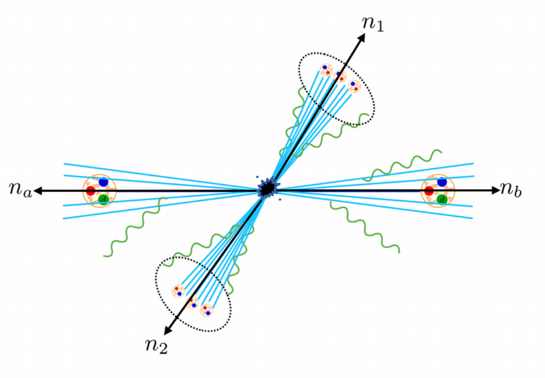
N-jet cross section:

Boughezal, Focke, Liu, Petriello(2015)
Gaunt, Stahlhofen, Tackmann, Walsh (2015)

$$\boxed{\sigma(X) = \int_0 \mathrm{d}\mathcal{T}_N \frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N} = \int_0^{\mathcal{T}_{\mathrm{cut}}} \mathrm{d}\mathcal{T}_N \frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N} + \int_{\mathcal{T}_{\mathrm{cut}}} \mathrm{d}\mathcal{T}_N \frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N}}$$

singular
N-jet final state

non-singualr
N+1-jet / multi-jets



Compute using factorization theorems in soft/collinear limits

$$\frac{\mathrm{d}\sigma(X)}{\mathrm{d}\mathcal{T}_N} = \underset{\text{Hard function}}{\mathbf{H}} \times \underbrace{B_a \otimes B_b}_{\text{Beam functions}} \otimes \underbrace{J_1 \otimes \ldots \otimes J_N}_{\text{Jet functions}} \otimes \underset{\text{N-jet soft function}}{S_N} + \mathcal{O}(\mathcal{T}_N)$$

Hard function: for many processes known to NNLO (e.g. W+jet) Gehrmann et al, (2012)

Beam function: known to NNLO Gaunt et al, (2014)

Jet function: known to NNLO [quark jet function NNNLO] Becher et al (2006), (2010) [Brüser et al. (2018)]

Soft function: for **N=0,1** known to NNLO Gaunt, et al (2015), Boughezal et al (2015), Compbell et al (2017)

This Talk

Calculate the N-jet soft function for N > 1 to NNLO

Automate soft function calculations

Idea: Automation

- Find generic strategy to evaluate soft functions (to NNLO)
- Set up a numerical method based on universal structure of divergences
  - ✓ Isolate singularities with universal phase-space parametrization
  - ✓ Compute observable dependent integrations numerically
  - ✓ SoftSERVE **Bell, Rahn, Talbert (to appear)**
    - Dijet soft functions (two light-like directions)
    - Explicit NNLO results for O(15) observables (e.g. jet grooming, jet vetoes, threshold and transverese momentum ressumation, e^+e^- event shapes)

Aim: extend framework for calculating N-jet soft functions at NNLO

Outline

Automating generic N-jet soft function calculation

(a) NLO: Real emission

Boost invariant parametrization

(b) NNLO: Virtual-Real interference

Double-Real emissions

N-jettiness soft function

(a) Constraints from RGE

(b) 1-jettiness ***Preliminary Results***

(c) 2-jettiness ***Preliminary Results***

Summary and outlook

N-jet soft functions

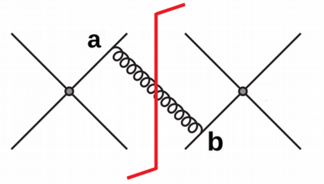
N-jet soft functions at NLO

- N-jet Soft functions

$$S_N(\tau,\mu) = \sum_X \mathcal{M}(\tau,\{k_i\})\,\langle 0|\left(S_{n_1} S_{n_2} S_{n_3}\dots\right)^\dagger |X\rangle\langle X|\left(S_{n_1} S_{n_2} S_{n_3}\dots\right)|0\rangle$$

multiple soft Wilson lines $\boldsymbol{S}_i(x) = \mathcal{P}\exp\left(ig_s\int_{-\infty}^{0} ds\, n_i\cdot A^a(x+sn_i)\,\boldsymbol{T}_i^a\right)$

$\mathcal{M}(\tau,\{k_i\})$ **generic measurement function**



✓ One-loop: Virtual corrections scaleless, real emissions diagrams contribute

✓ N-jet soft function at NLO: $S_N = \sum_{a\neq b} T_a\cdot T_b\, S_{ab}$

Catani, Grazzini (2000)
Catani, Seymour (1996)

$$S_{ab} \sim \int \mathrm{d}^d k\,\delta(k^2)\,\theta(k^0)\,\mathcal{M}(\tau,\{k_i\})\,|\mathcal{A}_{ab}(k)|^2$$

dipole matrix element

$$|\mathcal{A}_{ab}(k)|^2 \sim \frac{n_a\cdot n_b}{2\,n_a\cdot k\,n_b\cdot k}$$

Dijet matrix element

$$|\mathcal{A}(k)|^2 \sim \frac{n_+\cdot n_-}{2\,k_-\,k_+}$$

UNIVERSITÄT SIEGEN

Setup of calculation at NLO

Strategy

1. Boost invariant parametrization: use the **transverse momentum** and **rapidity** measure in the frame where each pair of **dipoles are back to back**

$$k_T = \sqrt{\frac{2\, k_a\, k_b}{n_{ab}}} \qquad y = \frac{k_a}{k_b} \qquad \begin{aligned} n_{ab} &\equiv n_a \cdot n_b \\ k_X &\equiv n_X \cdot k \end{aligned}$$

- **Parameterizing the solid angle:** Sudakov decomposition is a **Lorenz covariant relation**

$$k^\mu = k_b \frac{n_a^\mu}{n_{ab}} + k_a \frac{n_b^\mu}{n_{ab}} \underbrace{- k_{x_3} n^\mu_{x_3} - k_{x_4} n^\mu_{x_4} + \dots}_{k_\perp^\mu}$$

$$\begin{aligned} k_{x_3} &= -k_T \, \cos(\theta_1) \\ k_{x_4} &= -k_T \, \cos(\theta_2) \, \sin(\theta_1) \\ k_{x_d} &= -k_T \, \cos(\theta_{d-2}) \, \sin(\theta_{d-3}) \dots \sin(\theta_1) \end{aligned}$$

UNIVERSITÄT SIEGEN

Setup of calculation at NLO

2. Generic measurement function (inspired by **Laplace space**)

$$\boxed{\mathcal{M}(\tau;k) = \exp\Big(-\tau\, k_T\, y^{n/2}\,\sqrt{n_{ab}/2}\, f(y,\theta_1,\theta_2)\Big)}$$

- k_T dependence fixed on dimensional grounds
- $f(y,\theta_1,\theta_2)$ finite and non-zero in collinear limit y → 0
- Factorized part of kinematic dependences on n_{ab} : improves numerical convergence
- External kinematics are limited to 4-dim $\longrightarrow$ **2 angles** for N-jet processes

UNIVERSITÄT SIEGEN

Setup of calculation at NLO

3. Integrate k_T analytically

4. Derive a master formula

$$S_{ab}(\tau,\mu) \sim \overbrace{\frac{\Gamma(-2\epsilon)}{\Gamma(-\epsilon)}}^{\textbf{Soft divergence}} \left(\sqrt{n_{ab}/2}\,\tau e^{\gamma_E}\mu\right)^{2\epsilon}$$

$$\times \int_0^1 \mathrm{d}y \int_{-1}^1 \mathrm{d}\cos\theta_1\,\mathrm{d}\cos\theta_2\ \sin^{-1-2\,\epsilon}\theta_1\ \sin^{-2-2\,\epsilon}\theta_2\ \overbrace{y^{-1+n\epsilon}}^{\textbf{Collinear divergences}}\underbrace{\Big[f(y,\theta_1,\theta_2)\Big]^{2\epsilon}}_{\textbf{Measurement function}}$$

✓ Singularities from k_T → 0 and y → 0 are factorized

N-jettiness soft function

5. Isolate singularities with standard subtraction techniques:

$$\int_0^1 dx\ x^{-1+n\varepsilon}\ f(x) = \int_0^1 dx\ x^{-1+n\varepsilon}\ [\ \underbrace{f(x) - f(0)}_{\text{finite}} + \underbrace{f(0)}_{1/\varepsilon}\ ]$$

Two approaches

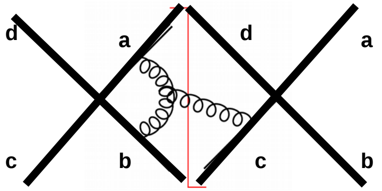
- pySecDec (results shown in this talk) Borowka, Heinrich, Jahn, Jones, Kerner, Schlenk,Zirke (2017)

  general implementation of sector decomposition algorithm
  Cuba library for numerical integrations

- SoftSERVE (in progress) Bell, Rahn, Talbert (to appear)

  C++ implementation for N-jet soft function
  Cuba library for numerical integrations

Setup of calculation at NNLO



- ✓ Two-Loop: Virtual corrections scaleless
- ✓ Real-Virtual contribution Catani, Grazzini (2000)

$$S_{\rm RV} = \sum_{a\neq b} T_a \cdot T_b \, S^R_{ab} + \underbrace{\sum_{a\neq b\neq c} (\lambda_{ab} - \lambda_{ak} - \lambda_{bk}) \, f_{ABC} \, T^A_a \, T^B_b \, T^C_c \, S^{Im}_{abc}}$$

$$\lambda_{XY} = \begin{cases} +1 & \text{if X and Y are both incoming/outgoing} \\ 0 & \text{otherwise} \end{cases}$$

Three-parton correlation (process dependent)

- ➢ **dipole contribution :** follow the same strategy of NLO

$$|\mathcal{A}^R_{ab}(k)|^2 \sim \left(\frac{n_{ab}}{2\,k_a\,k_b}\right)^{1+\epsilon}$$

Dijet matrix element

$$|\mathcal{A}(k)|^2 \sim \left(\frac{n_+ \cdot n_-}{2\,k_+\,k_-}\right)^{1+\epsilon}$$

- ➢ **tripole contribution:**
  - ✓ only present in processes with four or more hard partons
  - ✓ choose dipole n_a- n_c and follow the same strategy of NLO

$$|\mathcal{A}^{Im}_{abc}(k)|^2 \sim \left(\frac{n_{ac}}{2\,k_a\,k_c}\right)\left(\frac{n_{ab}}{2\,k_a\,k_b}\right)^{\epsilon}$$

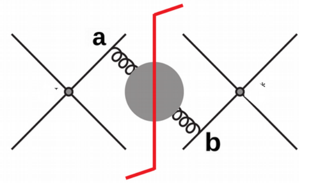
Setup of calculation at NNLO

- ✓ Double real corrections:

Catani, Grazzini (2000)

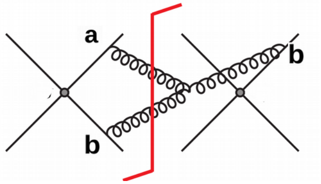
I) radiation of soft q$\bar{q}$ pair

$$S_N^{q\bar{q}} = T_F\, n_f \sum_{a\neq b} T_a \cdot T_b\, S_{ab}^{T_F n_f}$$



II) radiation of double-real gluons

$$S_N^{gg} = C_A \sum_{a\neq b} T_a \cdot T_b\, S_{ab}^{C_A}$$



III) tripole and quadrupole contributions are accounted for by non-abelian exponentiation

assume non-abelian exponentiation

- ➢ $T_F\, n_f$ structure

$$S_{ab}^{T_F n f} \sim \int \mathrm{d}^d k\, \delta(k^2)\, \theta(k^0) \int \mathrm{d}^d l\, \delta(l^2)\, \theta(l^0)\, \mathcal{M}(\tau; k, l) \left|\mathcal{A}_{ab}(k,l)\right|^2_{T_F n f}$$

matrix element

$$\left|\mathcal{A}_{ab}(k,l)\right|^2_{T_F n f} \sim \frac{2\, k\cdot l(k_i + l_i)(k_j + l_j) - (k_i\, l_j - l_i\, k_j)^2}{(k_i + l_i)^2\, (k_j + l_j)^2 \boxed{(2\, k \cdot l)^2}}$$

overlapping divergence

$$\left|\mathcal{A}(k,l)\right|^2_{C_F T_F n f} \sim \frac{2\, k\cdot l(k_- + l_-)(k_+ + l_+) - (k_-\, l_+ - l_-\, k_+)^2}{(k_- + l_-)^2\, (k_+ + l_+)^2\, (2\, k \cdot l)^2}$$

UNIVERSITÄT SIEGEN

Setup of calculation at NNLO

Strategy

1. Parametrization: collective and relative variables related to a two body system

$$p_T = \sqrt{\frac{2}{n_{ab}}(k_a+l_a)(k_b+l_b)} \qquad a = \frac{k_b\,l_a}{k_a\,l_b} = \sqrt{\frac{y_l}{y_k}}$$

$$y = \frac{k_a+l_a}{k_b+l_b} \qquad b = \sqrt{\frac{k_a\,k_b}{l_a\,l_b}} = \frac{k_T}{l_T}$$

2. Generic form of the measurement function: five angles in transverse plane

$$\boxed{\mathcal{M}(\tau;k,l) = \exp\Big(-\tau\,p_T\,y^{n/2}\,\sqrt{n_{ab}/2}\;F(a,b,y,\theta_{kl},\theta_{nk_1},\theta_{nk_2},\theta_{nl_1},\theta_{nl_2})\Big)}$$

- p_T dependence fixed on dimensional grounds
- $F(a,b,y,\theta_{kl},\theta_{nk_1},\theta_{nk_2},\theta_{nl_1},\theta_{nl_2})$ finite and non-zero for y → 0
- External kinematics are limited to 4-dim ⟶ **5 angles** for N-jet processes

Setup of calculation at NNLO

3&4. Integrate k_T analytically and obtain the master formula

$$S_{C_F T_F n_f}(\tau,\mu) \sim \frac{\overbrace{\Gamma(-4\epsilon)}^{\textbf{Soft divergence}}}{\Gamma(-\epsilon)\Gamma(1/2-\epsilon)} \left(\tau e^{\gamma_E}\mu\right)^{4\epsilon} \int_0^1 \mathrm{d}y\,\mathrm{d}a\,\mathrm{d}b$$

$$\times \int_{-1}^{1} \mathrm{d}\cos\theta_{kl}\ \mathrm{d}\cos\theta_{nk_1}\ \mathrm{d}\cos\theta_{nk_2}\ \sin^{-1-2\,\epsilon}\theta_{kl}\ \sin^{-2-2\,\epsilon}\theta_{nk_1}\ \sin^{-3-2\,\epsilon}\theta_{nk_2}$$

$$\times \int_{-1}^{1} \mathrm{d}\cos\theta_{nl_1}\ \mathrm{d}\cos\theta_{nl_2}\ \sin^{-1-2\,\epsilon}\theta_{nl_1}\ \sin^{-2-2\,\epsilon}\theta_{nl_2}$$

$$\times \underbrace{\frac{y^{-1+2n\epsilon}}{\left(1+a^2-2\,a\ \cos\theta_{kl}\right)^2}}_{\textbf{Collinear divergences}} \underbrace{\left[F(a,b,y,\theta_{kl},\theta_{nk_1},\theta_{nk_2},\theta_{nl_1},\theta_{nl_2})\right]^{4\epsilon}}_{\textbf{Measurement function}} \underbrace{\mathcal{J}(a,b,y,\epsilon)}_{\textbf{Matrix element Jacobian}}$$

Overlapping singularity:
I) Sector decomposition
II) Factorize the singularity with a simple change of variable

UNIVERSITÄT SIEGEN

Applications:

N-jettiness soft function

Solving RGE

- RGE for the renormalized soft function and the counterterm

$$\mu \frac{\mathrm{d}\, S(\tau,\mu)}{\mathrm{d}\mu} = \frac{1}{2}\gamma_s\, S(\tau,\mu) + \frac{1}{2} S(\tau,\mu)\,\gamma_s^\dagger$$

$$\mu \frac{\mathrm{d}\, Z_S(\tau,\mu)}{\mathrm{d}\mu} = -\frac{1}{2}\gamma_s S(\tau,\mu)$$

$$\boxed{i\pi\alpha_s^2 \left[\sum_{a\neq b} T_a\cdot T_b \ln(\sqrt{2\,n_{ab}})\,,\sum_{c\neq d} T_c\cdot T_d\,\Delta_{cd}\right]}$$

- Soft anomalous dimension given by consistency relation(RG invariance)

$$\gamma_s = \Gamma_{\text{cusp}}\left[-2\sum_{a\neq b} T_a\cdot T_b \ln\left(\sqrt{2\,n_{ab}}\,\mu\,\bar{\tau}\right) + i\pi\sum_{a\neq b} T_a\cdot T_b\Delta_{ab}\right] + \gamma_s^{\text{non}-\text{cusp}}$$

related to the anomalous dimension of hard Wilson Coefficient from matching QCD to SCET

$$\Delta_{ab} = \begin{cases} +1 & \text{if a and b are both incoming/outgoing} \\ 0 & \text{otherwise} \end{cases}$$

Solve iteratively for the bare soft function (provides a cross check for the poles)

$$S^{\text{bare}}(\tau) = Z_S(\tau,\mu)\, S(\tau,\mu)\, Z_S^\dagger(\tau,\mu)$$

N-jettiness soft function

The soft function in Laplace space

$$Z_\alpha = 1 - \left(\frac{\alpha_s}{4\pi}\right)\frac{\beta_0}{\epsilon}$$

$$\bar{\tau} = \tau\, e^{\gamma_E}$$

$$S(\tau,\mu) = 1 + \left(\frac{Z_\alpha\,\alpha_s}{4\,\pi}\right)\sum_{a\neq b}\mathbf{T_a}\cdot\mathbf{T_b}\left(\sqrt{2\,n_{ab}}\,\mu\,\bar{\tau}\right)^{2\epsilon} S^{(1)}_{ab}(\epsilon)$$
$$+\left(\frac{Z_\alpha\,\alpha_s}{4\,\pi}\right)^2\Bigg[\sum_{a\neq b}\mathbf{T_a}\cdot\mathbf{T_b}\left(\sqrt{2\,n_{ab}}\,\mu\,\bar{\tau}\right)^{4\epsilon} S^{(2)}_{ab}(\epsilon) + \sum_{a\neq b\neq c}\mathbf{f_{ABC}}\,\mathbf{T^A_a}\,\mathbf{T^B_b}\,\mathbf{T^C_c}\left(\mu\,\bar{\tau}\right)^{4\epsilon}\, S^{(2,Im)}_{ab}(\epsilon)$$
$$+\frac{1}{2}\sum_{a\neq b, c\neq d}\mathbf{T_a}\cdot\mathbf{T_b}\,\mathbf{T_c}\cdot\mathbf{T_d}\left(2\sqrt{n_{ab}\,n_{cd}}\,\mu^2\,\bar{\tau}^2\right)^{2\epsilon} S^{(1)}_{ab}(\epsilon)S^{(1)}_{cd}(\epsilon)\ \Bigg] + \mathcal{O}(\alpha_s^3)$$

Jouttenus, Stewart, Tackmann, Waalewijn (2011)

$$S^{(1)}_{ab}(\epsilon) = \overbrace{\frac{2}{\epsilon^2} + \frac{0}{\epsilon} + \mathbf{I^1_{ab}}}^{\text{known results for any number of jets}} + \epsilon\, K^1_{ab}$$

$$S^{(2)}_{ab}(\epsilon) = \Bigg(T_F n_f\left[-\frac{2}{3\,\epsilon^3} - \frac{10}{9\,\epsilon^2} + \frac{1}{\epsilon}\left(-\frac{56}{7} + \frac{\pi^2}{9} - \frac{4}{3}\mathbf{I^1_{ab}}\right) + I^{T_F n_f}_{ab}\right]$$
$$+ C_A\left[\frac{0}{\epsilon^4} + \frac{11}{6\,\epsilon^3} + \frac{1}{\epsilon^2}\left(\frac{67}{18} - \frac{\pi^2}{6}\right) + \frac{1}{\epsilon}\left(\frac{202}{27} - \frac{11\,\pi^2}{36} - 7\zeta_3 + \frac{11}{3}\mathbf{I^1_{ab,c}}\right) + I^{C_A}_{ab}\right]\Bigg)$$

This work: Preliminary Results

poles are known from RGE

UNIVERSITÄT SIEGEN

One-jettiness in pp collision

$$S^{(1)}_{ab}(\epsilon) = \frac{\mathbf{C^1_{-2}}}{\epsilon^2} + \frac{\mathbf{C^1_{-1}}}{\epsilon} + \mathbf{I^1_{ab}} + \epsilon\,\mathbf{K^1_{ab}}$$

$$S^{(2)}_{ab}(\epsilon) = \left(\mathbf{T_F n_f}\left[\frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{T_F n_f}_{ab}}\right] + \mathbf{C_A}\left[\frac{\mathbf{C^2_{-4}}}{\epsilon^4} + \frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{C_A}_{ab}}\right]\right)$$

$n_{12} = 2$

$n_{13} = 1 - \cos(\theta)$

$n_{23} = 2 - n_{13}$

1-Loop

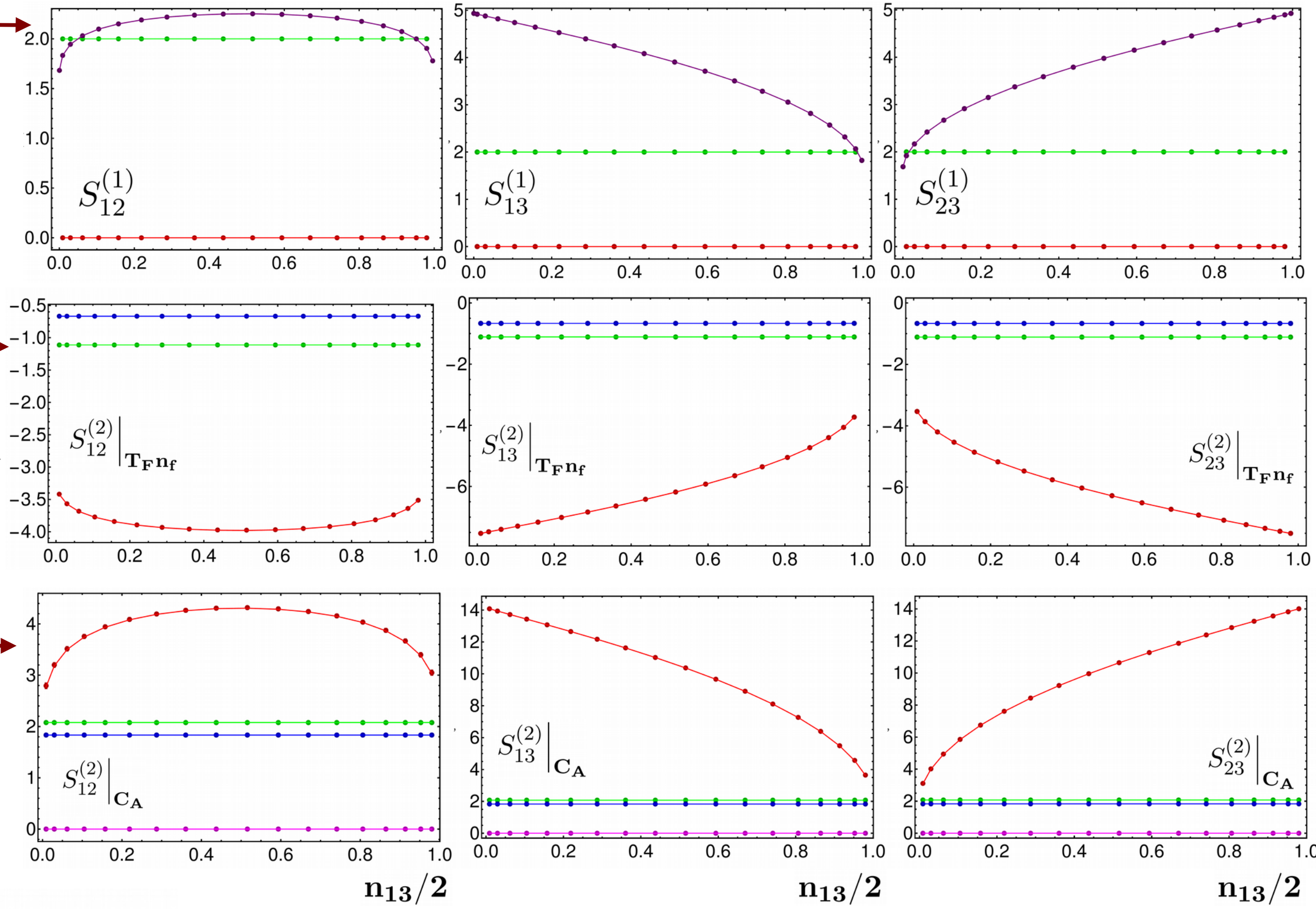
Relative Uncertainty < 0.01%

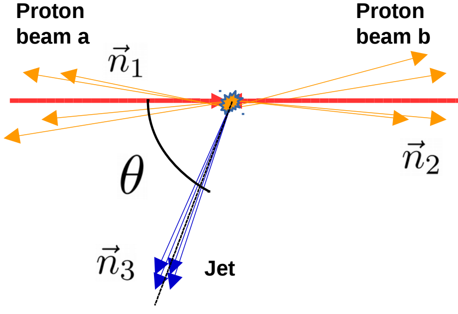
2-Loop $T_F\, n_f$

Relative Uncertainty < 0.1%

2-Loop C_A

Relative Uncertainty < 0.4%





*Our numerical results using VEGAS **(dots)** agree within the uncertainty with the known **results at NLO** and the **divergent terms at NNLO (lines)**.*

Preliminary Results

UNIVERSITÄT SIEGEN

Our predictions

One-jettiness in pp collision

$$S^{(1)}_{ab}(\epsilon) = \frac{\mathbf{C^1_{-2}}}{\epsilon^2} + \frac{\mathbf{C^1_{-1}}}{\epsilon} + \mathbf{I^1_{ab}} + \epsilon\,\mathbf{K^1_{ab}}$$

$$S^{(2)}_{ab}(\epsilon) = \left(\mathbf{T_F n_f}\left[\frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{T_F n_f}_{ab}}\right] + \mathbf{C_A}\left[\frac{\mathbf{C^2_{-4}}}{\epsilon^4} + \frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{C_A}_{ab}}\right]\right)$$

$n_{12} = 2$

$n_{13} = 1 - \cos(\theta)$

$n_{23} = 2 - n_{13}$

1-Loop

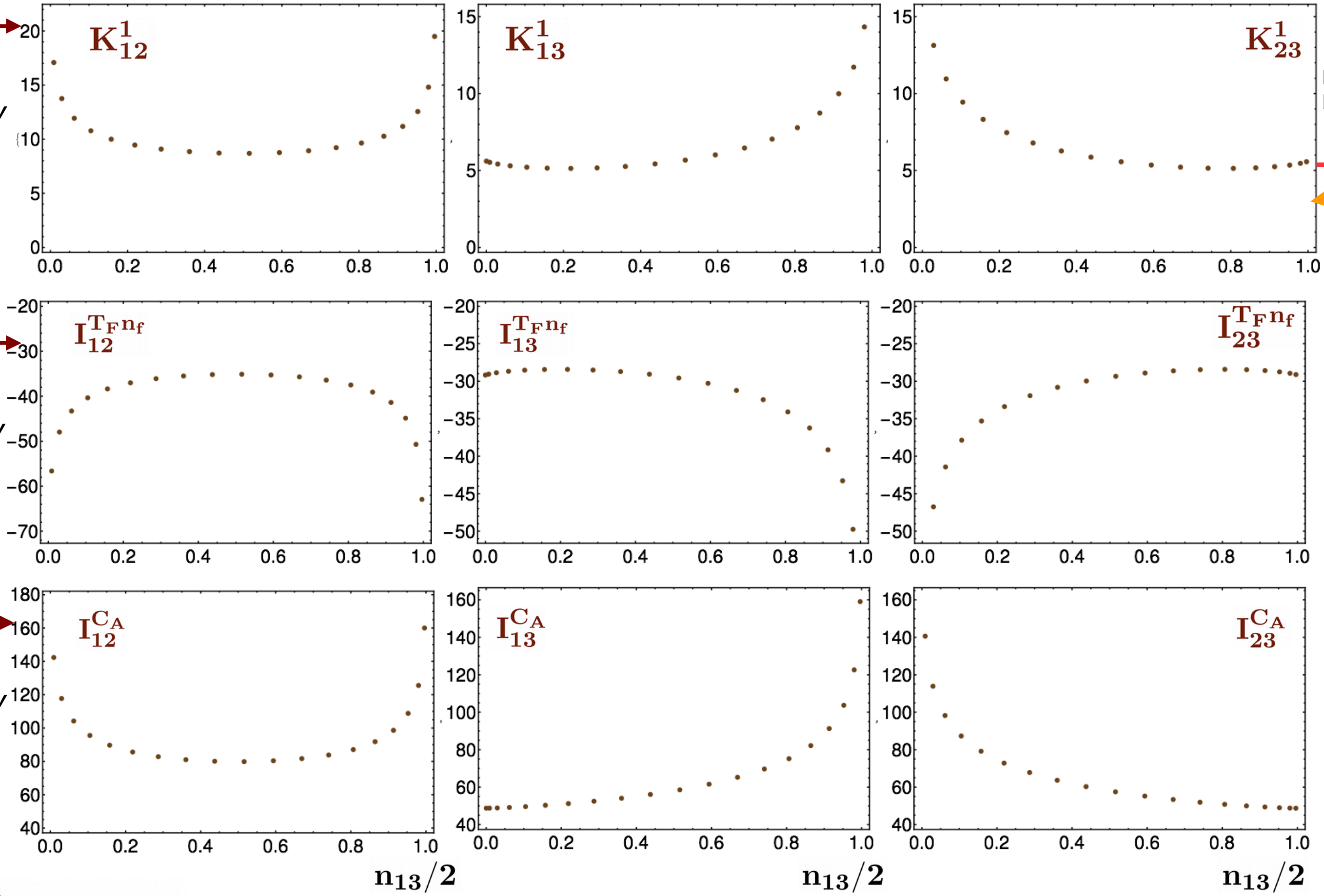
Relative Uncertainty < 0.01%

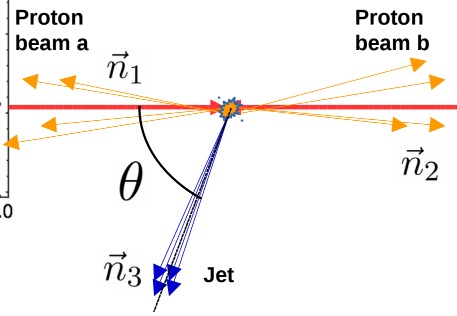
2-Loop $T_F n_f$

Relative Uncertainty < 0.1%

2-Loop C_A

Relative Uncertainty < 0.3%





Preliminary Results

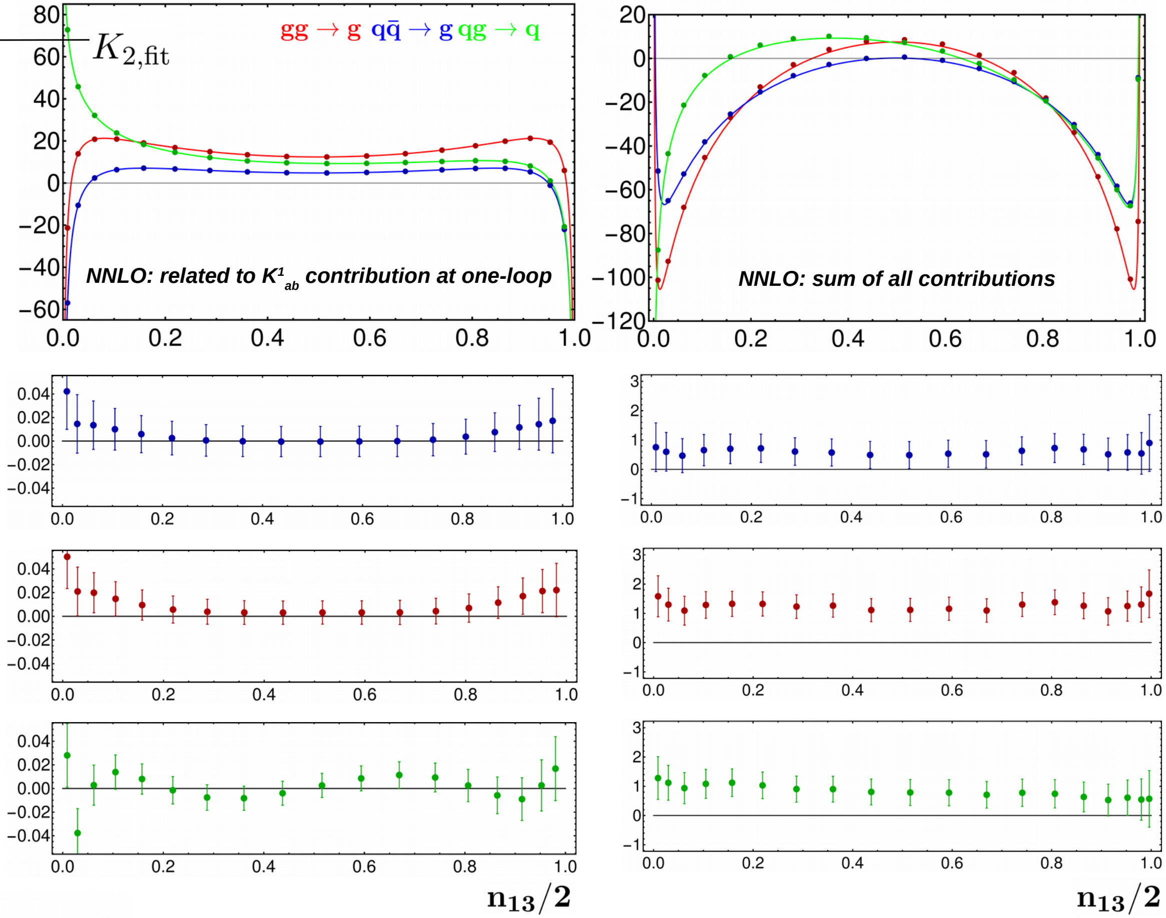
Numerical checks

One-jettiness in pp collision

Sum of the dipole contributions and **color factors** at NNLO for different partonic channels $gg \to g$, $q\bar{q} \to g$, $qg \to q$ in the distribution space (coefficients of $\delta(\mathcal{T}_1)$). **Our results (dots)** vs. fit result in **Ref.[2] (lines)** .

$$n_{12} = 2$$
$$n_{13} = 1 - \cos(\theta)$$
$$n_{23} = 2 - n_{13}$$

See Ref [2]



Difference between our results and the fit result in Ref [2]

Ref. [1]: Boughezal, Liu, Petriello (2015)

Ref. [2]: Campbell, Ellis, Mondini, Williams (2017)

Ref. [1]: provides one plot for qg → q

Ref. [2] provides useful fits to their numerical results. However we could not reconstruct their uncertainties!

Our numerical error estimates (w.i.p)

Preliminary Results

UNIVERSITÄT SIEGEN

Numerical checks

Two-jettiness in pp collision

$$S^{(1)}_{ab}(\epsilon) = \frac{\mathbf{C^{1}_{-2}}}{\epsilon^2} + \frac{\mathbf{C^{1}_{-1}}}{\epsilon} + \mathbf{I^{1}_{ab}} + \epsilon\,\mathbf{K^{1}_{ab}}$$

$$S^{(2)}_{ab}(\epsilon) = \left(\mathbf{T_F n_f}\left[\frac{\mathbf{C^{2}_{-3}}}{\epsilon^3} + \frac{\mathbf{C^{2}_{-2}}}{\epsilon^2} + \frac{\mathbf{C^{2}_{-1}}}{\epsilon} + \mathbf{I^{T_F n_f}_{ab}}\right] + \mathbf{C_A}\left[\frac{\mathbf{C^{2}_{-4}}}{\epsilon^4} + \frac{\mathbf{C^{2}_{-3}}}{\epsilon^3} + \frac{\mathbf{C^{2}_{-2}}}{\epsilon^2} + \frac{\mathbf{C^{2}_{-1}}}{\epsilon} + \mathbf{I^{C_A}_{ab}}\right]\right)$$

$n_{12} = 2$

$n_{13} = 1 - \cos(\theta)$

$n_{23} = 2 - n_{13}$

1-Loop

Relative Uncertainty < 0.01%

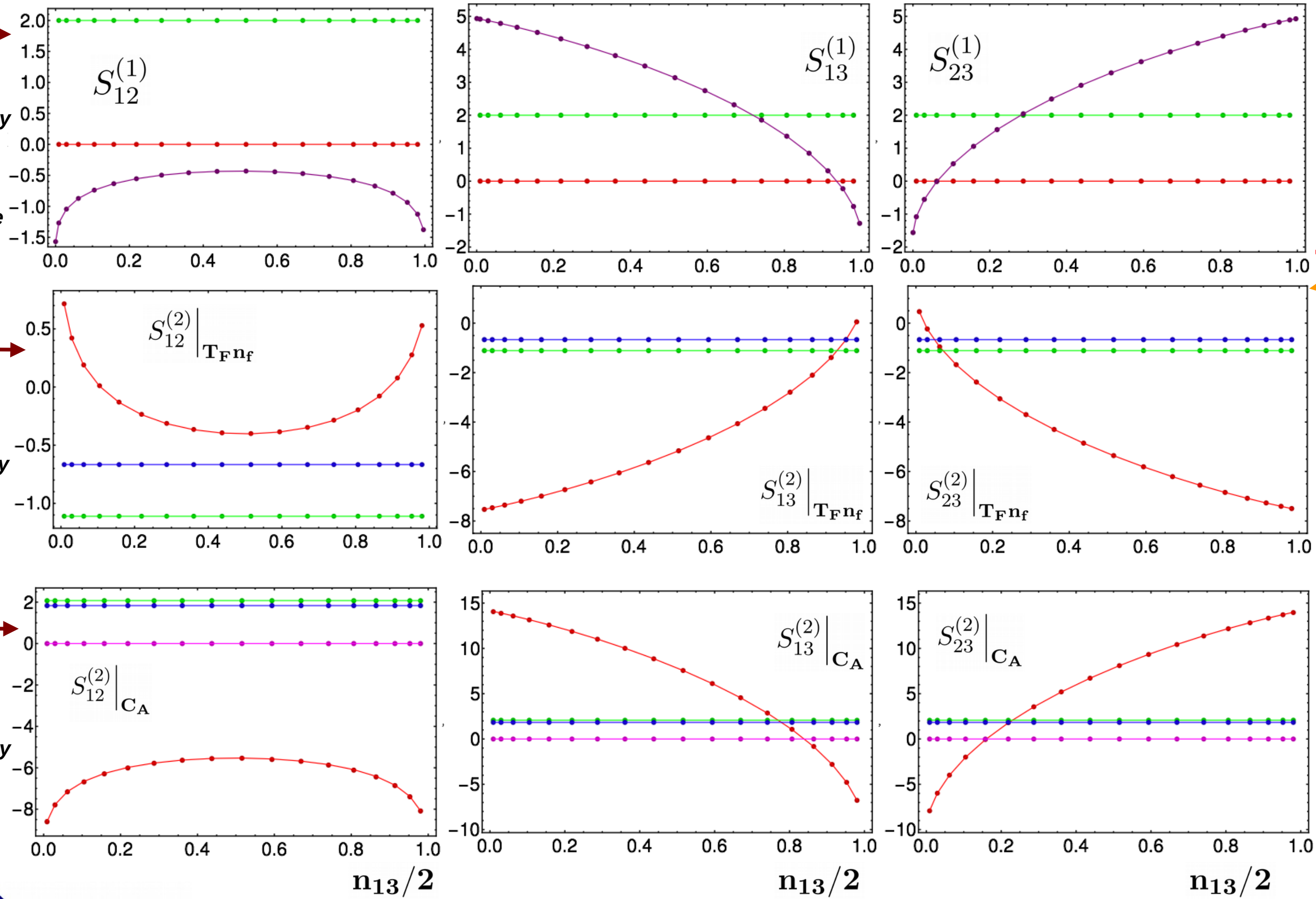
The first Purple line < 1%

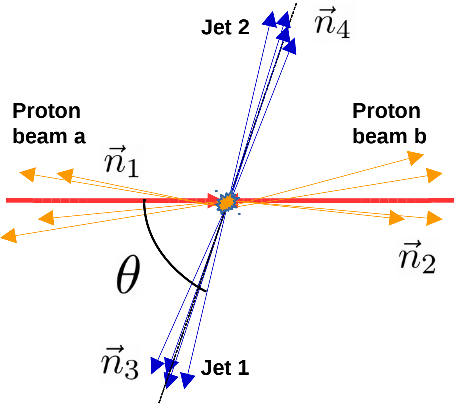
2-Loop $T_F n_f$

Relative Uncertainty < 0.1%

2-Loop C_A

Relative Uncertainty < 0.5%





*Our numerical results using VEGAS **(dots)** agree within the uncertainty with the known **results at NLO** and the **divergent terms at NNLO (lines)**.*

Preliminary Results

UNIVERSITÄT SIEGEN

New predictions

Two-jettiness in pp collision

$$S^{(1)}_{ab}(\epsilon) = \frac{\mathbf{C^1_{-2}}}{\epsilon^2} + \frac{\mathbf{C^1_{-1}}}{\epsilon} + \mathbf{I^1_{ab}} + \epsilon\,\mathbf{K^1_{ab}}$$

$$S^{(2)}_{ab}(\epsilon) = \left(\mathbf{T_F n_f}\left[\frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{T_F n_f}_{ab}}\right] + \mathbf{C_A}\left[\frac{\mathbf{C^2_{-4}}}{\epsilon^4} + \frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{C_A}_{ab}}\right]\right)$$

$n_{12} = 2$

$n_{13} = 1 - \cos(\theta)$

$n_{23} = 2 - n_{13}$

1-Loop

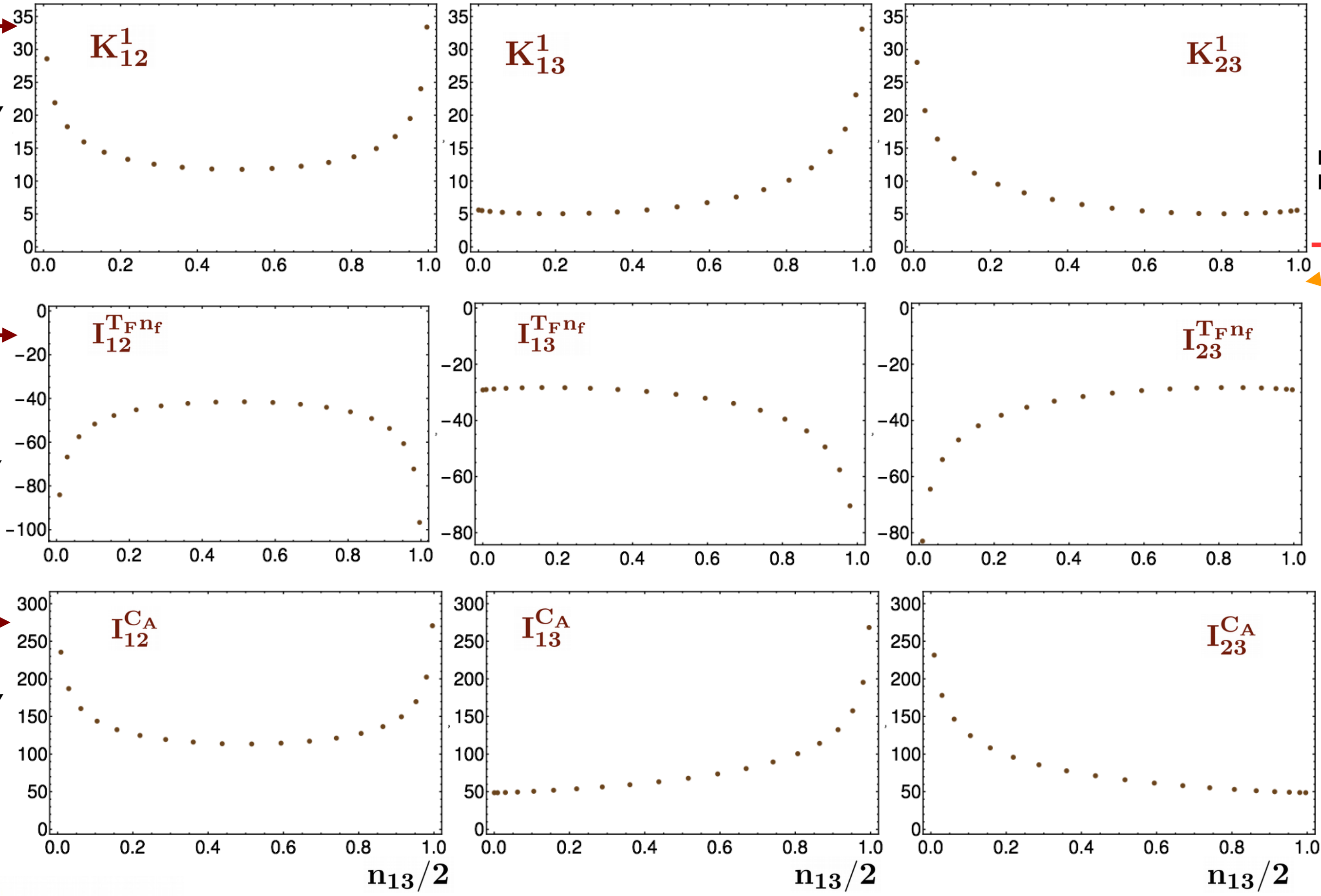
Relative Uncertainty < 0.2%

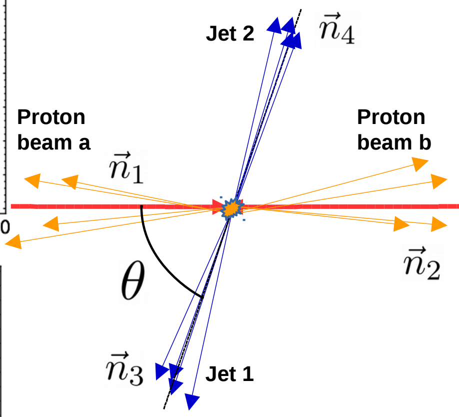
2-Loop $T_F n_f$

Relative Uncertainty < 0.1%

2-Loop C_A

Relative Uncertainty < 0.3%





Preliminary Results

UNIVERSITÄT SIEGEN

Two-jettiness in pp collision

$$\sum_{a\neq b\neq c} f_{ABC}\, T_a^A\, T_b^B\, T_c^C\, S_{ab}^{(2,Im)}(\epsilon) = f_{ABC}\, T_1^A\, T_2^B\, T_3^C\, I^{(2,Im)}(\epsilon)$$

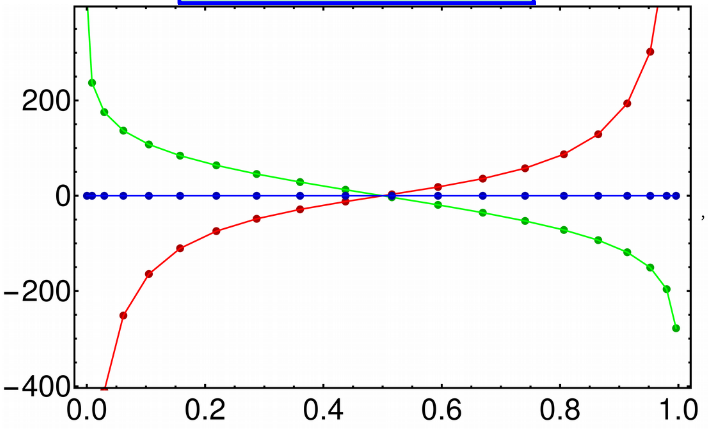
$$I^{(2,Im)}(\epsilon) = \frac{C_{-3}^{(2,Im)}}{\epsilon^3} + \frac{C_{-2}^{(2,Im)}}{\epsilon^2} + \frac{C_{-1}^{(2,Im)}}{\epsilon} + I_{ab}^{(2,Im)}$$

$n_{12} = 2$

$n_{13} = 1 - \cos(\theta)$

$n_{23} = 2 - n_{13}$

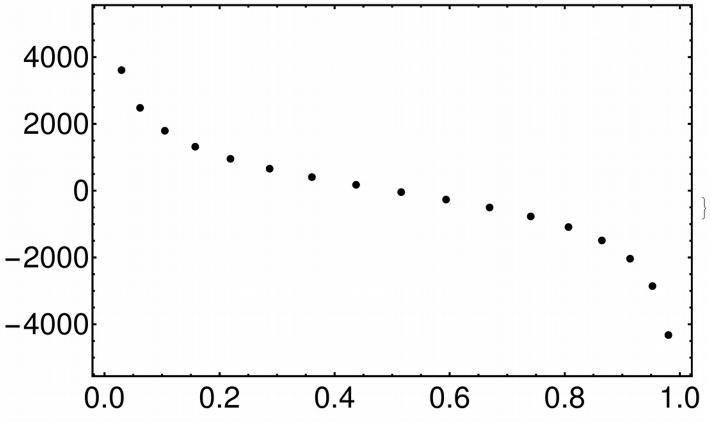
Numerical checks



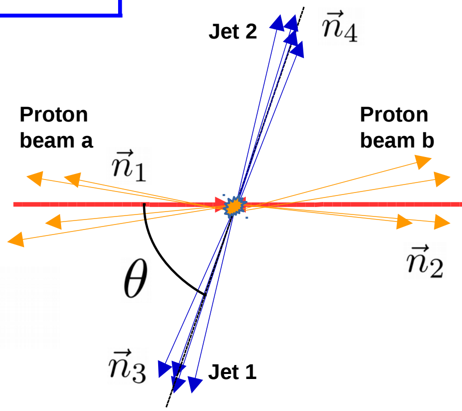
$\mathbf{n_{13}/2}$

Relative Uncertainty < 0.3%

New predictions



$\mathbf{n_{13}/2}$



- ✓ *Our numerical results using VEGAS **(dots)** agree within the uncertainty with the known **results at NLO** and the **divergent terms at NNLO (lines)**.*

Preliminary Results

Conclusions and outlook

Conclusions

✓ Systematic extension of our framework for automated calculations of N-jet soft functions

- First step assumes non-abelian exponentiation and SCET-1 type observable

✓ NNLO results

- Numerical results for 1-jettiness soft function
- **First numerical results for 2-jettiness soft function**
- **Our calculation allows to extend the N-jettiness slicing technique to processes with 2 jets**
- A reliable error estimate needs further studies (w.i.p)

Outlook

➢ Other observables on the horizon (angularities, boosted-tops, hadronic event shapes, etc) (w.i.p)

- may trigger new ideas for subtraction techniques

➢ N-jet implementation in SoftSERVE (w.i.p)

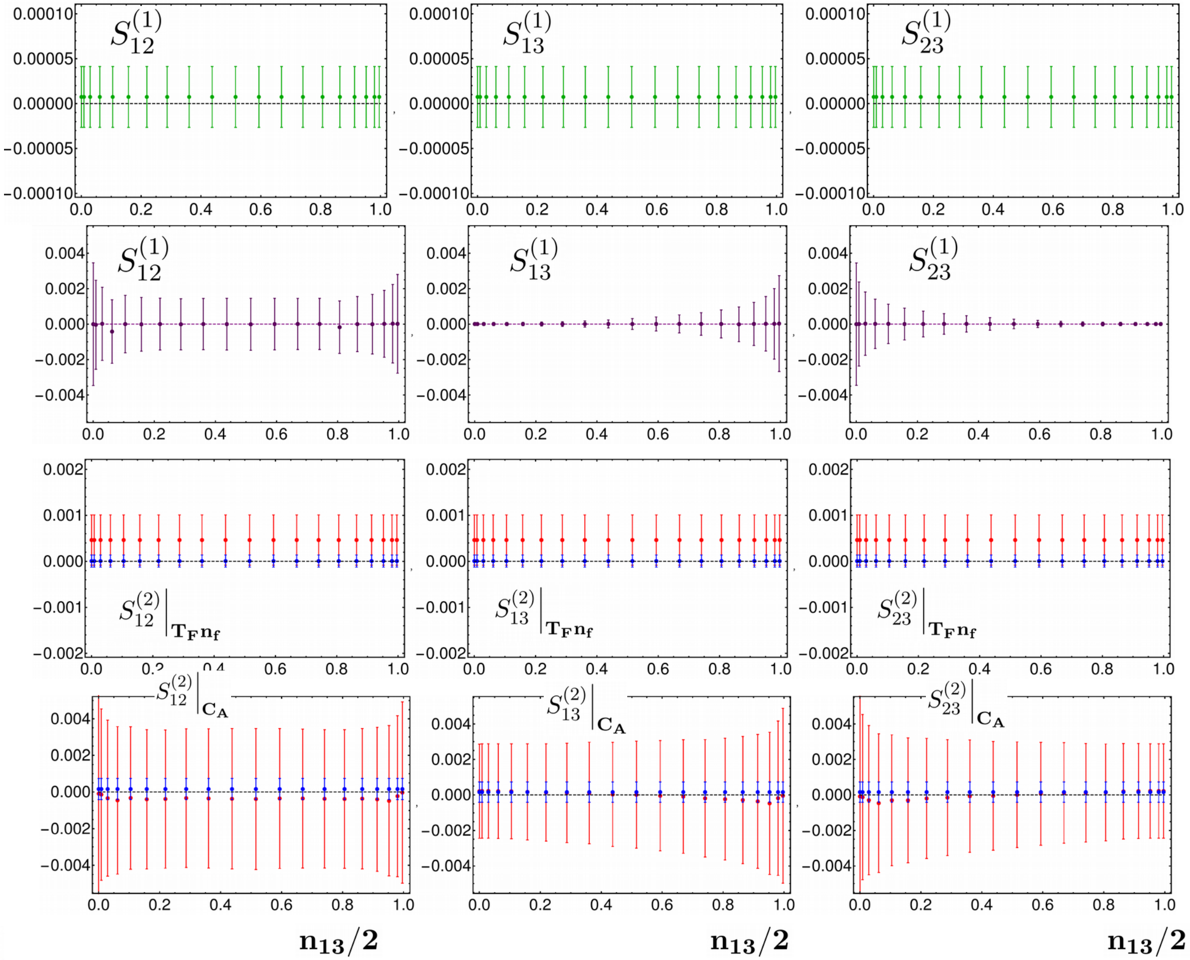
Thank you for your attention!

Back up slides

One-jettiness (RGE vs Numerics)

$$S_{ab}^{(1)}(\epsilon) = \frac{\mathbf{C^1_{-2}}}{\epsilon^2} + \frac{\mathbf{C^1_{-1}}}{\epsilon} + \mathbf{I^1_{ab}} + \epsilon\,\mathbf{K^1_{ab}}$$

$$S_{ab}^{(2)}(\epsilon) = \left(\mathbf{T_F n_f}\left[\frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{T_F n_f}_{ab}}\right] + \mathbf{C_A}\left[\frac{\mathbf{C^2_{-4}}}{\epsilon^4} + \frac{\mathbf{C^2_{-3}}}{\epsilon^3} + \frac{\mathbf{C^2_{-2}}}{\epsilon^2} + \frac{\mathbf{C^2_{-1}}}{\epsilon} + \mathbf{I^{C_A}_{ab}}\right]\right)$$



Preliminary Results

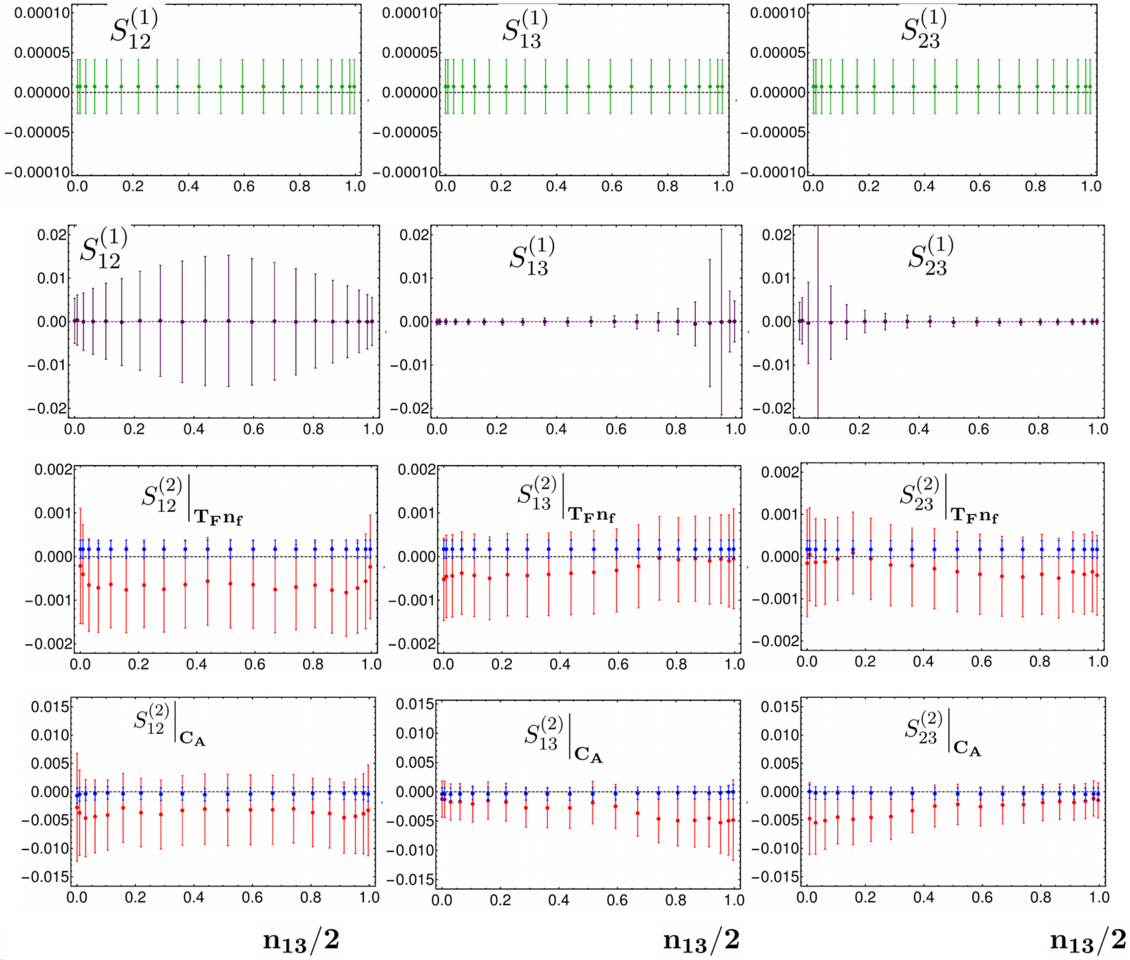
UNIVERSITÄT SIEGEN

Vienna 2018

Two-jettiness (RGE vs Numerics)

$$S^{(1)}_{ab}(\epsilon) = \frac{\mathbf{C^{1}_{-2}}}{\epsilon^2} + \frac{\mathbf{C^{1}_{-1}}}{\epsilon} + \mathbf{I^{1}_{ab}} + \epsilon\,\mathbf{K^{1}_{ab}}$$

$$S^{(2)}_{ab}(\epsilon) = \left(\mathbf{T_F n_f}\left[\frac{\mathbf{C^{2}_{-3}}}{\epsilon^3} + \frac{\mathbf{C^{2}_{-2}}}{\epsilon^2} + \frac{\mathbf{C^{2}_{-1}}}{\epsilon} + \mathbf{I^{T_F n_f}_{ab}}\right] + \mathbf{C_A}\left[\frac{\mathbf{C^{2}_{-4}}}{\epsilon^4} + \frac{\mathbf{C^{2}_{-3}}}{\epsilon^3} + \frac{\mathbf{C^{2}_{-2}}}{\epsilon^2} + \frac{\mathbf{C^{2}_{-1}}}{\epsilon} + \mathbf{I^{C_A}_{ab}}\right]\right)$$



Preliminary Results

UNIVERSITÄT SIEGEN