Silver Linings from the Flavor Sector

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Based on different collaborations with S. Descotes-Genon, J. Matias, F. Mescia, T. Hurth, L. Hofer, S. Neshatpour and F. Mahmoudi

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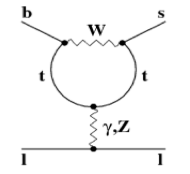
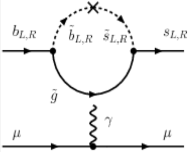
Effective Operators for Flavor Physics

Effective lagrangian at the hadronic scale

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}_{\text{QED}} + \sum_i \left(\frac{c_i^{\text{SM}}}{M_W^2} + \frac{c_i^{\text{NP}}}{\Lambda_1^2} + \cdots \right) \mathcal{O}_i^{(6)} + \cdots$$

- The coefficients c_i^{SM} are suppressed:
  1. Arise at the loop level
  2. Proportional to $(V_{\text{CKM}}^\dagger)_{ts}(V_{\text{CKM}})_{tb}$
  3. GIM $\to$ Zero for degenerate quarks
  4. Helicity suppression: only $V-A$ interactions, etc.
- Test the SM hyphothesis $c_i^{\text{NP}} = 0$
- Measure $c_i^{\text{NP}}/\Lambda_{\text{NP}}$ to learn about the NP.
- Complementarity with high-p_T searches (Λ_{NP}).



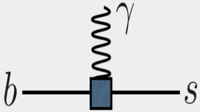


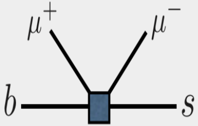
Radiative and Dileptonic $b \to s$ Operators

Significant progress made recently regarding $b \to s\gamma/s\ell\ell$ operators

Radiative and Dileptonic $b \to s$ Operators

$$
\begin{aligned}
\mathcal{O}_{7(\prime)} &= [\bar{s}\sigma^{\mu\nu}P_{R(L)}b]F_{\mu\nu} \\
\mathcal{O}_{9(\prime)} &= [\bar{s}\gamma^{\mu}P_{L(R)}b][\bar{\ell}\gamma_{\mu}\ell] \\
\mathcal{O}_{10(\prime)} &= [\bar{s}\gamma^{\mu}P_{L(R)}b][\bar{\ell}\gamma_{\mu}\gamma_5\ell]
\end{aligned}
$$





Effective Hamiltonian

$$
\mathcal{H}_{\rm eff}^{\rm sl} = -\frac{4G_F}{\sqrt{2}}V_{tb}V_{ts}^{*}\Big[\mathcal{C}_7\mathcal{O}_7 + \mathcal{C}_{7'}\mathcal{O}_{7'} + \mathcal{C}_9\mathcal{O}_9 + \mathcal{C}_{9'}\mathcal{O}_{9'} + \mathcal{C}_{10}\mathcal{O}_{10} + \mathcal{C}_{10'}\mathcal{O}_{10'}\Big]
$$

Note: We write $\mathcal{C}_i = \mathcal{C}_i^{\rm SM} + \mathcal{C}_i^{\rm NP}$:

$$
\mathcal{C}_{7\rm eff}^{\rm SM} = -0.29,\ \mathcal{C}_9^{\rm SM} = 4.07,\ \mathcal{C}_{10}^{\rm SM} = -4.31,\ \mathcal{C}_{i'}^{\rm SM} \simeq 0
$$

Radiative and Dileptonic $b \to s$ Operators

How to Measure Radiative and Dileptonic Operators?

1. Identify decay modes and observables most sensitive to such ops

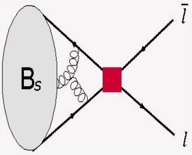
Decay modes for $b \to s\gamma$ and $b \to s\ell\ell$

$B \to X_s\gamma$	$B \to X_s\ell\ell$	$B_s \to \ell^+\ell^-$
$B \to K^*\gamma$	$B \to K^*\ell\ell$	$B \to K\ell\ell$

2. Compute the observables in the effective theory
3. Buy a full set of non-perturbative parameters from the Black Market
4. Fit the data, extract CL intervals for the $\mathcal{C}_i(m_b)$.
5. Interpret the results.

Note: We will fit directly to $\mathcal{C}_i^{\rm NP}(m_b)$

$$\mathcal{L} = \mathcal{L}_{QED+QCD} + \mathcal{C}_9^{(\prime)}\,[\bar{s}\gamma^\mu P_{L(R)} b]\,[\bar{\ell}\gamma_\mu \ell] + \mathcal{C}_{10}^{(\prime)}\,[\bar{s}\gamma^\mu P_{L(R)} b]\,[\bar{\ell}\gamma_\mu\gamma_5 \ell] + \cdots$$



$$\mathcal{A}_9^{(\prime)} = \mathcal{C}_9^{(\prime)}\,\langle \bar{\ell}\ell|\bar{\ell}\gamma_\mu\ell|0\rangle\,\underbrace{\langle 0|\bar{s}\gamma^\mu P_{L(R)} b|B_s\rangle}_{\sim p_B^\mu = p_\ell^\mu + p_{\bar{\ell}}^\mu} = 0 + \mathcal{O}(\alpha)$$

Contributions from $\mathcal{O}_7$ and other 4-quark ops are zero like $\mathcal{A}_9^{(\prime)}$.

$$\to \mathcal{A}_{10}^{(\prime)} = \mathcal{C}_{10}^{(\prime)}\,\langle \bar{\ell}\ell|\bar{\ell}\gamma_\mu\gamma_5\ell|0\rangle\langle 0|\bar{s}\gamma^\mu P_{L(R)} b|B_s\rangle = \mp i f_{B_s}\mathcal{C}_{10}^{(\prime)}\, m_\ell\,[\bar{u}_\ell\gamma_5 v_{\bar{\ell}}]$$

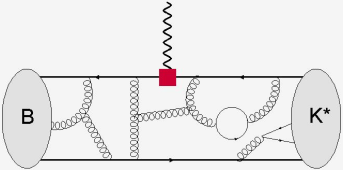
$$\to \sum_{\text{spins}} |\mathcal{A}_{10} + \mathcal{A}_{10}'|^2 = 2 f_{B_s}^2 m_{B_s}^2 m_\ell^2\,|\mathcal{C}_{10} - \mathcal{C}_{10}'|^2$$

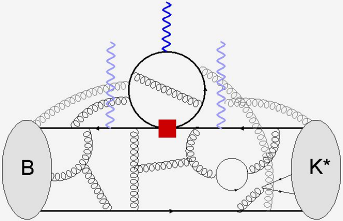
$$\to BR(B_s \to \ell\bar{\ell}) = \frac{\tau_{B_s} f_{B_s}^2 m_{B_s}^3}{8\pi}\,\frac{m_\ell^2}{m_{B_s}^2}\sqrt{1 - \frac{4m_\ell^2}{m_{B_s}^2}}\;|\mathcal{C}_{10} - \mathcal{C}_{10}'|^2$$

Note: Contributions from (pseudo)SCALAR operators are **not helicity suppressed**.

EFT Amplitudes & Observables : $B \to K^{(*)}\gamma^{(*)}$

$$\mathcal{L} = \mathcal{L}_{QED+QCD} + \mathcal{C}_7\,[\bar{s}\sigma^{\mu\nu}P_R b]\,F_{\mu\nu} + \mathcal{C}_2\,[\bar{s}\gamma^\nu P_L c]\,[\bar{c}\gamma^\mu P_L b] + \cdots$$





$\mathcal{C}_7$ **contribution:** $\mathcal{A}_7 = \mathcal{C}_7\,\langle K^*_\lambda|\bar{s}\sigma_{\mu\nu}P_R b|B\rangle q^\mu \epsilon^\nu_\lambda = \mathcal{C}_7\,T_\lambda(q^2)$

$\mathcal{C}_2$ **contribution:** $\mathcal{A}_2 = \mathcal{C}_2 \cdot \epsilon^{*\mu}_\lambda \int dx^4 e^{iq\cdot x}\langle K^*_\lambda|T\{j^{c\bar{c}}_\mu(x)\mathcal{O}_2(0)\}|B\rangle$

Note: There are similar contributions from $\mathcal{O}_8$ and other 4-quark ops. These operators are contained in what we call $\mathcal{H}_{\rm eff}^{\rm had}$.

EFT Amplitudes & Observables : $B \to K\ell\bar{\ell}$

$$\mathcal{A}_9^{(\prime)} = \mathcal{C}_9^{(\prime)}\, \langle \bar{\ell}\ell | \bar{\ell}\gamma_\mu \ell | 0\rangle \langle K | \bar{s}\gamma^\mu P_{L(R)} b | B\rangle \;=\; \mathcal{C}_9^{(\prime)}\, f_+(q^2)\, [\bar{u}_\ell \not{p} v_{\bar{\ell}}]$$

$$\mathcal{A}_{10}^{(\prime)} = \mathcal{C}_{10}^{(\prime)}\, \langle \bar{\ell}\ell | \bar{\ell}\gamma_\mu\gamma_5 \ell | 0\rangle \langle K | \bar{s}\gamma^\mu P_{L(R)} b | B\rangle = \mathcal{C}_{10}^{(\prime)}\, f_+(q^2)\, [\bar{u}_\ell \not{p}\gamma_5 v_{\bar{\ell}}]$$

$$\mathcal{A}_7^{(\prime)} = \mathcal{C}_7^{(\prime)}\, \langle \bar{\ell}\ell | \bar{\ell}\gamma_\mu \ell | 0\rangle \frac{-i}{q^2} \langle K | \bar{s} q_\mu \sigma^{\mu\nu} P_{R(L)} b | B\rangle = \mathcal{C}_7^{(\prime)}\, \frac{f_T(q^2)}{m_B + m_K} [\bar{u}_\ell \not{p} v_{\bar{\ell}}]$$

$$\mathcal{A}_{\rm had} = \mathcal{K}(q^2) [\bar{u}_\ell \not{p} v_{\bar{\ell}}]$$

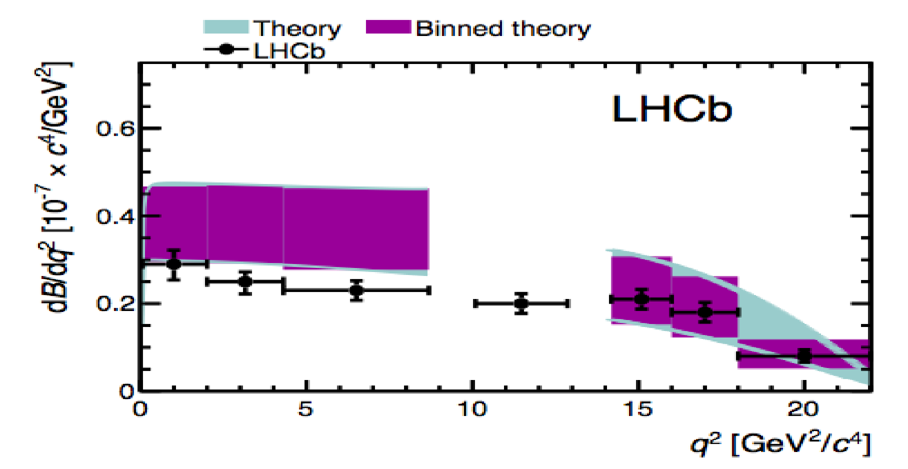
$$\mathcal{A}(B \to K\ell\bar{\ell}) = a_9\, [\bar{u}_\ell \not{p} v_{\bar{\ell}}] + a_{10}\, [\bar{u}_\ell \not{p}\gamma_5 v_{\bar{\ell}}]$$

$$a_9 = (\mathcal{C}_9 + \mathcal{C}_9')\, f_+(q^2) + (\mathcal{C}_7 + \mathcal{C}_7')\, \frac{f_T(q^2)}{m_B + m_K} + \mathcal{K}(q^2)\;; \quad a_{10} = (\mathcal{C}_{10} + \mathcal{C}_{10}')\, f_+(q^2)$$

$$\frac{d\Gamma}{ds_{13}\, ds_{23}} = \frac{m_B^5}{2^8 \pi^3} (|a_9|^2 + |a_{10}|^2) s_{13} s_{23}$$

Where $s_{13} = 2p_\ell \cdot p_K / m_B^2$, $s_{23} = 2p_{\bar{\ell}} \cdot p_K / m_B^2$ – $s_{12} = q^2/m_B^2$

EFT Amplitudes & Observables : $B \to K\ell\bar{\ell}$



$$\frac{d\Gamma}{ds_{13}\,ds_{23}} = \frac{m_B^5}{2^8\pi^3}\left(|a_9|^2 + |a_{10}|^2\right)s_{13}s_{23}$$

Where $s_{13} = 2p_\ell \cdot p_K/m_B^2$, $s_{23} = 2p_{\bar{\ell}} \cdot p_K/m_B^2$ – $s_{12} = q^2/m_B^2$

$B \to K^*(\to K\pi)\ell^+\ell^-$ and its Angular Distribution

Structure of the Decay Amplitude

"Semileptonic" contribution $\longrightarrow$ New Physics

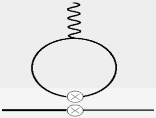
$$\langle K^*\ell\ell|\mathcal{O}_{9^{(\prime)},10^{(\prime)}}|B\rangle = \langle \ell^+\ell^-|\bar{\ell}\gamma_\mu(\gamma_5)\ell|0\rangle\langle K^*|\bar{s}\gamma^\mu P_{L,R}b|B\rangle \sim F_{i,\lambda}^{B\to K*}(q^2)$$

$$\langle K^*\ell\ell|T\{j_{em}^\ell\mathcal{O}_{7^{(\prime)}}\}|B\rangle = \langle \ell^+\ell^-|\bar{\ell}\gamma_\mu\ell|0\rangle\frac{q_\nu}{q^2}\langle K^*|\bar{s}\sigma^{\mu\nu}P_{R,L}b|B\rangle \sim T_{i,\lambda}^{B\to K*}(q^2)$$

$$\mathcal{A}^{\rm sl} = \sum_i f_i(\mathcal{C}_{7^{(\prime)}},\mathcal{C}_{9^{(\prime)}},\mathcal{C}_{10^{(\prime)}}) \times (\text{Form Factor})_i$$

"Hadronic" contribution $\longrightarrow$ QCD $[\mathcal{C}_{1,2},\mathcal{C}_8,\mathcal{C}_{3,4,5,6}]$

$$\mathcal{A}^{\rm had} = i\frac{e^2}{q^2}\langle \ell^+\ell^-|\bar{\ell}\gamma_\mu\ell|0\rangle\int d^4x e^{iq\cdot x}\langle K^*|T\{j_{em}^\mu(x)\mathcal{H}_{\rm eff}^{\rm had}(0)\}|B\rangle$$

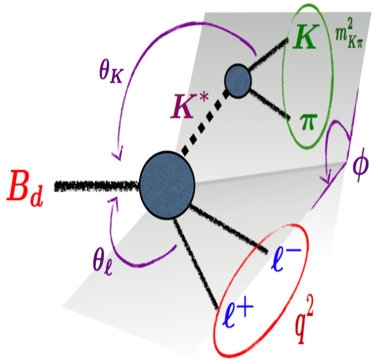


2 main problems:

1. Precise determination of Form Factors (LCSRs, LQCD, ...)
2. Computation of the hadronic contribution (SCET/QCDF, OPE, ...)

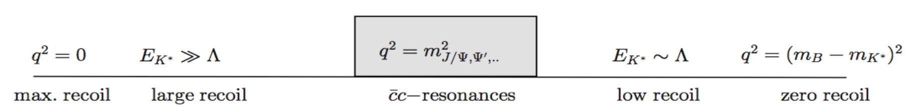
$B \to K^*(\to K\pi)\ell^+\ell^-$ and its Angular Distribution

Kinematics and angular distribution



$$\frac{d^4\Gamma}{dq^2\, d\cos\theta_K\, d\cos\theta_l\, d\phi} = \frac{9}{32\pi} \times$$

$$\Big[\mathbf{J}_{1s}\sin^2\theta_K + \mathbf{J}_{1c}\cos^2\theta_K + \mathbf{J}_{2s}\sin^2\theta_K\cos 2\theta_l$$
$$+\mathbf{J}_{2c}\cos^2\theta_K\cos 2\theta_l + \mathbf{J}_3\sin^2\theta_K\sin^2\theta_l\cos 2\phi$$
$$+\mathbf{J}_4\sin 2\theta_K\sin 2\theta_l\cos\phi + \mathbf{J}_5\sin 2\theta_K\sin\theta_l\cos\phi$$
$$+\mathbf{J}_{6s}\sin^2\theta_K\cos\theta_l + \mathbf{J}_{6c}\cos^2\theta_K\cos\theta_l$$
$$+\mathbf{J}_7\sin 2\theta_K\sin\theta_l\sin\phi + \mathbf{J}_8\sin 2\theta_K\sin 2\theta_l\sin\phi$$
$$+\mathbf{J}_9\sin^2\theta_K\sin^2\theta_l\sin 2\phi\Big]$$



$B \to K^*(\to K\pi)\ell^+\ell^-$ and its Angular Distribution

Optimized Observables

Several Form Factor ratios can be predicted:

- At large recoil $\longrightarrow$ SCET — Charles et.al. hep-ph/9812358, Beneke, Feldmann, hep-ph/0008255
- At low recoil $\longrightarrow$ HQET — Grinstein, Pirjol, hep-ph/0404250, Bobeth, Hiller, van Dyk

Example — SCET relation at large recoil

$$\frac{\epsilon_-^{*\mu} q^\nu \langle K_-^* | \bar{s}\sigma_{\mu\nu} P_R b | B\rangle}{i m_B \langle K_-^* | \bar{s}\not{\epsilon}_-^* P_L b | B\rangle} = 1 + \mathcal{O}(\alpha_s, \Lambda/m_b)$$

This allows to build observables with **reduced dependence on FFs**.

Optimized observables at large recoil — Matias, Mescia, Ramon, JV – 1202.4266; Descotes-Genon, Matias, Ramon, JV – 1207.2753

$$P_1 = \frac{J_3}{2J_{2s}} \qquad P_2 = \frac{J_{6s}}{8J_{2s}} \qquad P_4' = \frac{J_4}{\sqrt{-J_{2s}J_{2c}}}$$

$$P_5' = \frac{J_5}{2\sqrt{-J_{2s}J_{2c}}} \qquad P_6' = \frac{-J_7}{2\sqrt{-J_{2s}J_{2c}}} \qquad P_8' = \frac{-J_8}{\sqrt{-J_{2s}J_{2c}}}$$

Form Factor SCET Relations and Clean Observables

SCET Heavy-Light Currents

Charles et.al. hep-ph/9812358, Beneke, Feldmann, hep-ph/0008255
Bauer, Fleming, Pirjol, Stewart, hep-ph/0011336

⋆ χ_n, h_v are 2-comp spinors : $(\bar{q}\Gamma b)_{QCD} \longrightarrow \bar{\chi}_n h_v,\ \bar{\chi}_n\gamma_5 h_v, \bar{\chi}_n\gamma_\perp^\mu h_v$

⋆ Only 2 "soft" Form Factors:

$$\langle K_n^*|\chi_n h_v|B_v\rangle = 0\ ;\ \langle K_n^*|\chi_n\gamma_5 h_v|B_v\rangle = -2m_{K^*}\xi_\parallel(E)v\cdot\epsilon^*\ ;\ \langle K_n^*|\chi_n\gamma_\perp^\mu h_v|B_v\rangle = 2E\xi_\perp(E)i\epsilon_\perp^{\mu\nu}\epsilon_\nu^*$$

⋆ Large-recoil Form-Factor relations: (up to (known) $\mathcal{O}(\alpha_s)$ + (unknown) PCs)

$$\frac{m_B}{m_B+m_{K^*}}V(q^2) = \frac{m_B+m_{K^*}}{2E}A_1(q^2) = T_1(q^2) = \frac{m_B}{2E}T_2(q^2) = \xi_\perp(E)$$

$$\frac{m_{K^*}}{E}A_0(q^2) = \frac{m_B+m_{K^*}}{2E}A_1(q^2) - \frac{m_B-m_{K^*}}{m_B}A_2(q^2) = \frac{m_B}{2E}T_2(q^2) - T_3(q^2) = \xi_\parallel(E)$$

Amplitudes and Clean Observables

Kruger, Matias – hep-ph/0502060
Matias, Mescia, Ramon, JV – 1202.4266

$$A_\perp^{L,R} = a_\perp[(\mathcal{C}_9+\mathcal{C}_9')\mp(\mathcal{C}_{10}+\mathcal{C}_{10}')]V(q^2) + b_\perp(\mathcal{C}_7+\mathcal{C}_7')T_1(q^2)+\cdots = f_\perp^{L,R}\xi_\perp+\cdots$$

$$A_\parallel^{L,R} = a_\parallel[(\mathcal{C}_9-\mathcal{C}_9')\mp(\mathcal{C}_{10}-\mathcal{C}_{10}')]A_1(q^2) + b_\parallel(\mathcal{C}_7+\mathcal{C}_7')T_2(q^2)+\cdots = f_\parallel^{L,R}\xi_\perp+\cdots$$

$$P_1 = \frac{J_3}{2J_2s} = \frac{|A_\perp^L|^2-|A_\parallel^L|^2+|A_\perp^R|^2-|A_\parallel^R|^2}{|A_\perp^L|^2+|A_\parallel^L|^2+|A_\perp^R|^2+|A_\parallel^R|^2} = \frac{f_\perp^L-f_\parallel^L+f_\perp^R-f_\parallel^R}{f_\perp^L+f_\parallel^L+f_\perp^R+f_\parallel^R}+\cdots$$

$B \to K^*(\to K\pi)\ell^+\ell^-$ and its Angular Distribution

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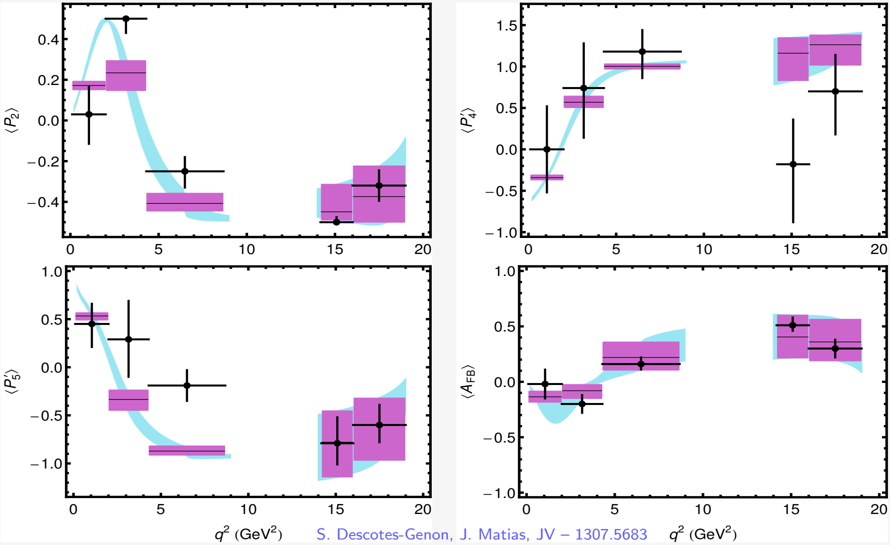
Theory vs. Experiment (LHCb: April'13 + July'13)



S. Descotes-Genon, J. Matias, JV – 1307.5683

Fitting the data: Set-up

Strategy:

We fit to **47** observables by means of a frequentist χ^2 approach.

Observables included in the analysis

$$BR(B \to X_s\gamma), \quad BR(B \to X_s\mu^+\mu^-)_{Low\ q^2}$$

$$BR(B_s \to \mu^+\mu^-), \quad A_I(B \to K^*\gamma), \quad S(B \to K^*\gamma)$$

$$B \to K^*\mu^+\mu^- : \ \langle P_1\rangle, \langle P_2\rangle, \langle P_4'\rangle, \langle P_5'\rangle, \langle P_6'\rangle, \langle P_8'\rangle, \langle A_{\rm FB}\rangle$$

in several different bins

Observables *not* included in the analysis

$$B \to K\mu^+\mu^-, \quad B_s \to \phi\mu^+\mu^-, \quad B \to X_s\mu^+\mu^- \text{ @ } Large\ q^2, \ \ldots$$

not considered for different reasons

Fitting the data: Set of data and pulls

Observable	Experiment	SM prediction	Pull
$\langle P_1\rangle_{[0.1,2]}$	$-0.19^{+0.40}_{-0.35}$	$0.007^{+0.043}_{-0.044}$	-0.5
$\langle P_1\rangle_{[2,4.3]}$	$-0.29^{+0.65}_{-0.46}$	$-0.051^{+0.046}_{-0.046}$	-0.4
$\langle P_1\rangle_{[4.3,8.68]}$	$0.36^{+0.30}_{-0.31}$	$-0.117^{+0.056}_{-0.052}$	$+1.5$
$\langle P_1\rangle_{[1,6]}$	$0.15^{+0.39}_{-0.41}$	$-0.055^{+0.041}_{-0.043}$	$+0.5$
$\langle P_2\rangle_{[0.1,2]}$	$0.03^{+0.14}_{-0.15}$	$0.172^{+0.020}_{-0.021}$	-1.0
$\langle P_2\rangle_{[2,4.3]}$	$0.50^{+0.00}_{-0.07}$	$0.234^{+0.060}_{-0.086}$	$+2.9$
$\langle P_2\rangle_{[4.3,8.68]}$	$-0.25^{+0.07}_{-0.08}$	$-0.407^{+0.049}_{-0.037}$	$+1.7$
$\langle P_2\rangle_{[1,6]}$	$0.33^{+0.11}_{-0.12}$	$0.084^{+0.060}_{-0.078}$	$+1.8$
$\langle P_4'\rangle_{[0.1,2]}$	$0.00^{+0.52}_{-0.52}$	$-0.342^{+0.031}_{-0.026}$	$+0.7$
$\langle P_4'\rangle_{[2,4.3]}$	$0.74^{+0.54}_{-0.60}$	$0.569^{+0.073}_{-0.063}$	$+0.3$
$\langle P_4'\rangle_{[4.3,8.68]}$	$1.18^{+0.26}_{-0.32}$	$1.003^{+0.028}_{-0.032}$	$+0.6$
$\langle P_4'\rangle_{[1,6]}$	$0.58^{+0.32}_{-0.36}$	$0.555^{+0.067}_{-0.058}$	$+0.1$
$\langle P_5'\rangle_{[0.1,2]}$	$0.45^{+0.21}_{-0.24}$	$0.533^{+0.033}_{-0.041}$	-0.4
$\langle P_5'\rangle_{[2,4.3]}$	$0.29^{+0.40}_{-0.39}$	$-0.334^{+0.097}_{-0.113}$	$+1.6$
$\langle P_5'\rangle_{[4.3,8.68]}$	$-0.19^{+0.16}_{-0.16}$	$-0.872^{+0.053}_{-0.041}$	$+4.0$
$\langle P_5'\rangle_{[1,6]}$	$0.21^{+0.20}_{-0.21}$	$-0.349^{+0.088}_{-0.100}$	$+2.5$
$\langle P_6'\rangle_{[0.1,2]}$	$0.24^{+0.23}_{-0.20}$	$-0.084^{+0.034}_{-0.044}$	$+1.6$
$\langle P_6'\rangle_{[2,4.3]}$	$-0.15^{+0.38}_{-0.36}$	$-0.098^{+0.043}_{-0.056}$	-0.1
$\langle P_6'\rangle_{[4.3,8.68]}$	$0.04^{+0.16}_{-0.16}$	$-0.027^{+0.060}_{-0.063}$	$+0.4$
$\langle P_6'\rangle_{[1,6]}$	$0.18^{+0.21}_{-0.21}$	$-0.089^{+0.042}_{-0.052}$	$+1.3$
$\langle P_8'\rangle_{[0.1,2]}$	$-0.12^{+0.56}_{-0.56}$	$0.037^{+0.037}_{-0.030}$	-0.3
$\langle P_8'\rangle_{[2,4.3]}$	$-0.30^{+0.60}_{-0.58}$	$0.070^{+0.045}_{-0.034}$	-0.6
$\langle P_8'\rangle_{[4.3,8.68]}$	$0.58^{+0.34}_{-0.38}$	$0.020^{+0.054}_{-0.055}$	$+1.5$
$\langle P_8'\rangle_{[1,6]}$	$0.46^{+0.36}_{-0.38}$	$0.063^{+0.042}_{-0.033}$	$+1.0$
$\langle A_{\rm FB}\rangle_{[0.1,2]}$	$-0.02^{+0.13}_{-0.13}$	$-0.136^{+0.051}_{-0.048}$	$+0.8$
$\langle A_{\rm FB}\rangle_{[2,4.3]}$	$-0.20^{+0.08}_{-0.08}$	$-0.081^{+0.055}_{-0.069}$	-1.1
$\langle A_{\rm FB}\rangle_{[4.3,8.68]}$	$0.16^{+0.06}_{-0.05}$	$0.220^{+0.138}_{-0.113}$	-0.5
$\langle A_{\rm FB}\rangle_{[1,6]}$	$-0.17^{+0.06}_{-0.06}$	$-0.035^{+0.037}_{-0.034}$	-2.0

Observable	Experiment	SM prediction	Pull
$\langle P_1\rangle_{[14.18,16]}$	$0.07^{+0.26}_{-0.28}$	$-0.352^{+0.697}_{-0.468}$	$+0.6$
$\langle P_1\rangle_{[16,19]}$	$-0.71^{+0.36}_{-0.26}$	$-0.603^{+0.589}_{-0.315}$	-0.2
$\langle P_2\rangle_{[14.18,16]}$	$-0.50^{+0.03}_{-0.00}$	$-0.449^{+0.136}_{-0.041}$	-1.1
$\langle P_2\rangle_{[16,19]}$	$-0.32^{+0.08}_{-0.08}$	$-0.374^{+0.151}_{-0.126}$	$+0.3$
$\langle P_4'\rangle_{[14.18,16]}$	$-0.18^{+0.54}_{-0.70}$	$1.161^{+0.190}_{-0.332}$	-2.1
$\langle P_4'\rangle_{[16,19]}$	$0.70^{+0.44}_{-0.52}$	$1.263^{+0.119}_{-0.248}$	-1.1
$\langle P_5'\rangle_{[14.18,16]}$	$-0.79^{+0.27}_{-0.22}$	$-0.779^{+0.328}_{-0.363}$	$+0.0$
$\langle P_5'\rangle_{[16,19]}$	$-0.60^{+0.21}_{-0.18}$	$-0.601^{+0.282}_{-0.367}$	$+0.0$
$\langle P_6'\rangle_{[14.18,16]}$	$0.18^{+0.24}_{-0.25}$	$0.000^{+0.000}_{-0.000}$	$+0.7$
$\langle P_6'\rangle_{[16,19]}$	$-0.31^{+0.38}_{-0.39}$	$0.000^{+0.000}_{-0.000}$	-0.8
$\langle P_8'\rangle_{[14.18,16]}$	$-0.40^{+0.60}_{-0.50}$	$-0.015^{+0.009}_{-0.013}$	-0.6
$\langle P_8'\rangle_{[16,19]}$	$0.12^{+0.52}_{-0.54}$	$-0.008^{+0.005}_{-0.007}$	$+0.2$
$\langle A_{\rm FB}\rangle_{[14.18,16]}$	$0.51^{+0.07}_{-0.05}$	$0.404^{+0.199}_{-0.191}$	$+0.5$
$\langle A_{\rm FB}\rangle_{[16,19]}$	$0.30^{+0.08}_{-0.08}$	$0.360^{+0.205}_{-0.172}$	-0.3
$10^4\,\mathcal{B}_{B\to X_s\gamma}$	3.43 ± 0.22	3.15 ± 0.23	$+0.9$
$10^6\,\mathcal{B}_{B\to X_s\mu^+\mu^-}$	1.60 ± 0.50	1.59 ± 0.11	$+0.0$
$10^9\,\mathcal{B}_{B_s\to\mu^+\mu^-}$	2.9 ± 0.8	3.56 ± 0.18	-0.8
$A_I(B\to K^*\gamma)$	0.052 ± 0.026	0.041 ± 0.025	$+0.3$
$S_{K^*\gamma}$	-0.16 ± 0.22	-0.03 ± 0.01	-0.6

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Fitting the data: Patterns

Simplified Linearized expressions:

$$\delta\langle P_2\rangle_{[0.1,2]} \simeq +0.37\,\mathcal{C}_7^{\rm NP} \qquad -0.03\,\mathcal{C}_{10}^{\rm NP} \quad \ominus$$

$$\delta\langle P_2\rangle_{[2,4.3]} \simeq -2.48\,\mathcal{C}_7^{\rm NP} -0.17\,\mathcal{C}_9^{\rm NP} +0.03\,\mathcal{C}_{10}^{\rm NP} \quad \oplus$$

$$\delta\langle P_2\rangle_{[4.3,8.68]} \simeq -0.71\,\mathcal{C}_7^{\rm NP} -0.09\,\mathcal{C}_9^{\rm NP} -0.04\,\mathcal{C}_{10}^{\rm NP} \quad \oplus$$

$$\delta\langle P_4'\rangle_{[0.1,2]} \simeq +0.59\,\mathcal{C}_7^{\rm NP} -0.08\,\mathcal{C}_9^{\rm NP} -0.13\,\mathcal{C}_{10}^{\rm NP} \quad \oplus$$

$$\delta\langle P_4'\rangle_{[2,4.3]} \simeq +2.45\,\mathcal{C}_7^{\rm NP} +0.06\,\mathcal{C}_9^{\rm NP} -0.14\,\mathcal{C}_{10}^{\rm NP} \quad \oplus$$

$$\delta\langle P_4'\rangle_{[4.3,8.68]} \simeq +0.33\,\mathcal{C}_7^{\rm NP} +0.01\,\mathcal{C}_9^{\rm NP} \quad \oplus$$

$$\delta\langle P_5'\rangle_{[0.1,2]} \simeq -0.91\,\mathcal{C}_7^{\rm NP} -0.12\,\mathcal{C}_9^{\rm NP} -0.03\,\mathcal{C}_{10}^{\rm NP} \quad \ominus$$

$$\delta\langle P_5'\rangle_{[2,4.3]} \simeq -3.04\,\mathcal{C}_7^{\rm NP} -0.29\,\mathcal{C}_9^{\rm NP} -0.03\,\mathcal{C}_{10}^{\rm NP} \quad \oplus$$

$$\delta\langle P_5'\rangle_{[4.3,8.68]} \simeq -0.52\,\mathcal{C}_7^{\rm NP} -0.08\,\mathcal{C}_9^{\rm NP} -0.03\,\mathcal{C}_{10}^{\rm NP} \quad \oplus$$

Fitting the data: Results

General Fit

Coefficient	1σ	2σ	3σ
$\mathcal{C}_7^{\rm NP}$	$[-0.05, -0.01]$	$[-0.06, 0.01]$	$[-0.08, 0.03]$
$\mathcal{C}_9^{\rm NP}$	$[-1.6, -0.9]$	$[-1.8, -0.6]$	$[-2.1, -0.2]$
$\mathcal{C}_{10}^{\rm NP}$	$[-0.4, 1.0]$	$[-1.2, 2.0]$	$[-2.0, 3.0]$
$\mathcal{C}_{7'}^{\rm NP}$	$[-0.04, 0.02]$	$[-0.09, 0.06]$	$[-0.14, 0.10]$
$\mathcal{C}_{9'}^{\rm NP}$	$[-0.2, 0.8]$	$[-0.8, 1.4]$	$[-1.2, 1.8]$
$\mathcal{C}_{10'}^{\rm NP}$	$[-0.4, 0.4]$	$[-1.0, 0.8]$	$[-1.4, 1.2]$

- Negative values for $(\mathcal{C}_7^{\rm NP}, \mathcal{C}_9^{\rm NP})$ favoured at $> (1\sigma,$**3**$\sigma)$.
- Large-recoil only $\longrightarrow$ effect enhanced ($\mathcal{C}_9^{\rm NP} \sim -1.6$).
- Only [1-6] bin: Same pattern, less significance.

Pulls : Most Economic Scenarios

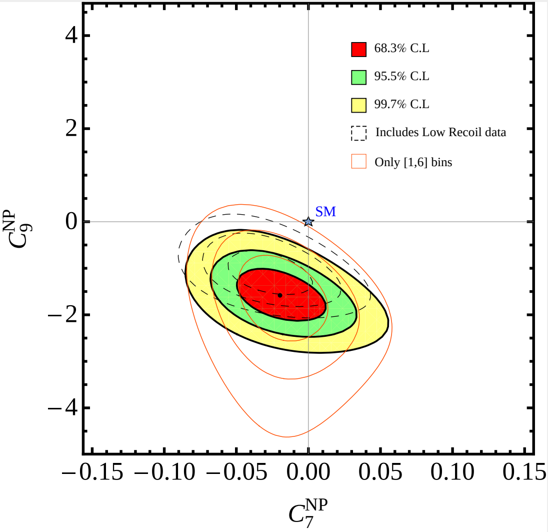
ONLY LARGE RECOIL

	C7	C7'	C9	C9'	C10	C10'
	3.09484	1.3451	4.92649	2.48843	0.43637	2.03476
	------	------	------	------	------	------
C7	*	2.9282	0.937218	2.74639	3.09598	2.84969
C7'	0.897572	*	0.227127	0.364298	1.27602	1.48412
C9	3.94597	4.74474	*	4.37932	4.90937	4.48669
C9'	2.03886	2.12502	1.049	*	2.48784	2.75007
C10	0.444375	0.0966814	0.148416	0.43296	*	3.18616
C10'	1.63798	2.12922	0.0174349	2.34752	3.75519	*

LARGE + LOW RECOIL

	C7	C7'	C9	C9'	C10	C10'
	3.06539	1.18839	4.2551	0.902501	0.238176	1.23639
	------	------	------	------	------	------
C7	*	2.9255	1.34744	2.98177	3.06376	2.93211
C7'	0.757771	*	0.0671292	0.932908	1.16507	1.17631
C9	3.2442	4.08633	*	4.19246	4.32275	4.14197
C9'	0.555753	0.522075	0.534224	*	0.870537	0.659162
C10	0.216167	0.0429785	0.798155	0.00738279	*	0.887756
C10'	0.853995	1.22479	0.760724	1.07175	1.50335	*

$\mathcal{C}_7^{\rm NP}$- $\mathcal{C}_9^{\rm NP}$ Scenario



- At 68.5% CL:

  $\mathcal{C}_7^{\rm NP} \in [-0.035, 0.000]$

  $\mathcal{C}_9^{\rm NP} \in [-1.9, -1.3]$

- Pulls for SM Hyp.:

  Large-recoil: **4.5**σ

  Large + Low-recoil: **3.9**σ

  Only [1-6] GeV bin: **3.2**σ

- The overall quality of the fit is very good.

Fitting the data: Results

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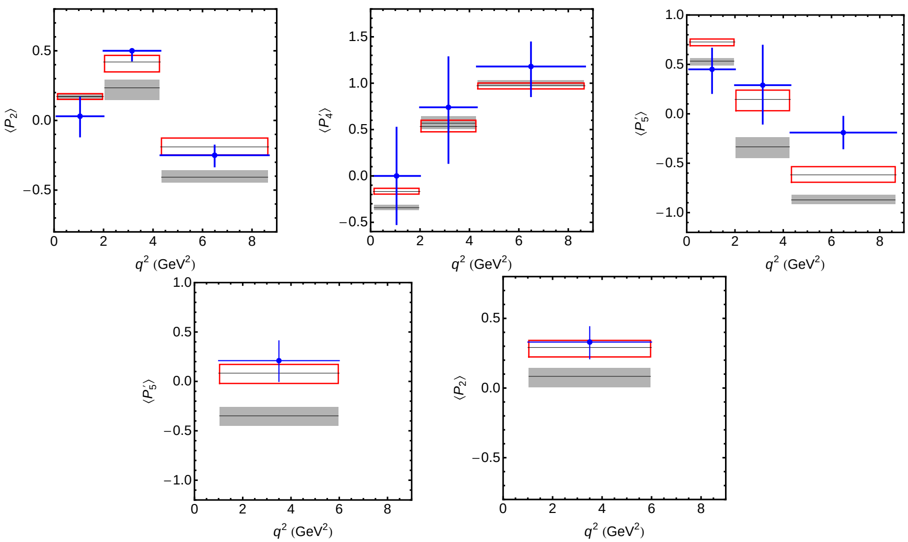
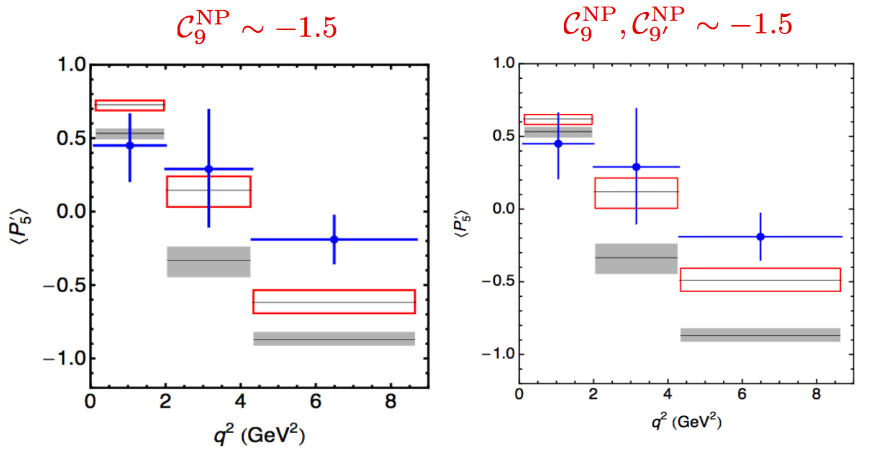
$\mathcal{C}_7^{\rm NP}$- $\mathcal{C}_9^{\rm NP}$ **Scenario:** Best-Fit point vs SM



$\mathcal{C}_9$ - $\mathcal{C}_9'$ Scenario

⋆ $\mathcal{C}_{9'}^{\rm NP} < 0$ can improve the situation with P_5'.

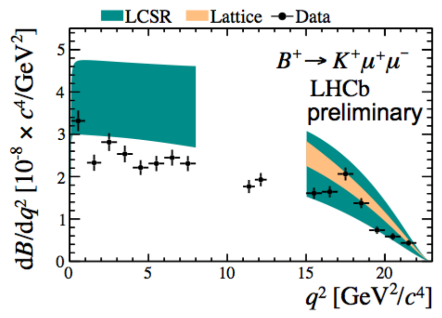


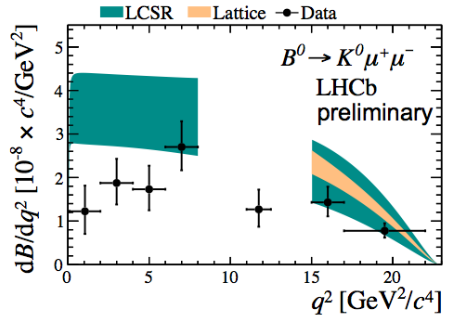
But increases tension with $BR(B \to K\mu\mu)$...

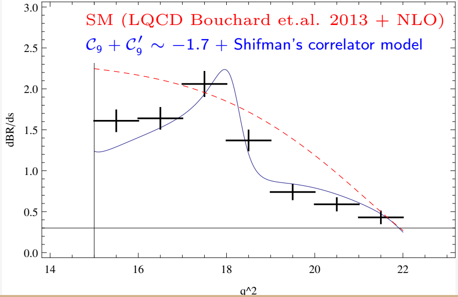
Altmannshofer, Straub 1308.1501,
Beaujean, Bobeth, van Dyk 1310.2478

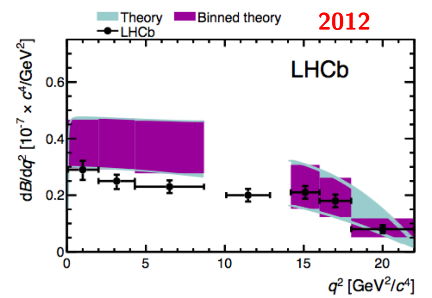
$B \to K\mu\mu$ Branching Ratios

LHCb-PAPER-2014-007







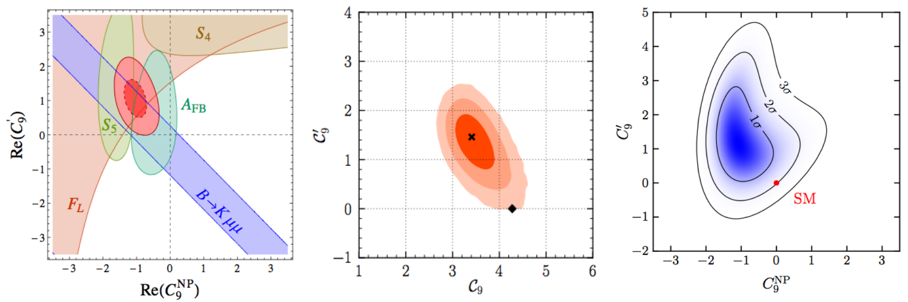


Summary / Remarks

- A global fit to $b \to s\gamma$, $b \to s\mu\mu$ observables including the latest data on $B \to K^*\mu\mu$ angular observables shows some level of tension w.r.t the SM, pointing (mostly) to a large NP contribution to $\mathcal{C}_9$.

  S. Descotes-Genon, J. Matias, JV – 1307.5683

- This has been later confirmed by other groups



Altmannshofer, Straub 1308.1501, Beaujean, Bobeth, van Dyk 1310.2478, Horgan et al. 1310.3887

Summary / Remarks

- A global fit to $b \to s\gamma$, $b \to s\mu\mu$ observables including the latest data on $B \to K^*\mu\mu$ angular observables shows some level of tension w.r.t the SM, pointing (mostly) to a large NP contribution to $\mathcal{C}_9$.

  S. Descotes-Genon, J. Matias, JV – 1307.5683

- This has been later confirmed by other groups

- New experimental analyses with the full 3 fb^{-1} of data will clarify a bit more the situation. Also new experimental initiatives:
  - Fit for the q^2-dependent amplitudes within some ansatz.
  - Fit directly for the WCs.
  - Improve on the binning.

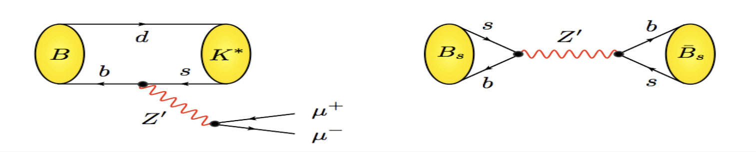
- Still a lot to do from the theory side:
  - FFs, hadronic contributions, PCs, resonance tails, etc.
  - New modes & observables: $B_s \to \phi\mu\mu$, $\Lambda_b \to \Lambda\mu\mu$ ...
  - Implications on NP models...

Epilogue: NP scale?

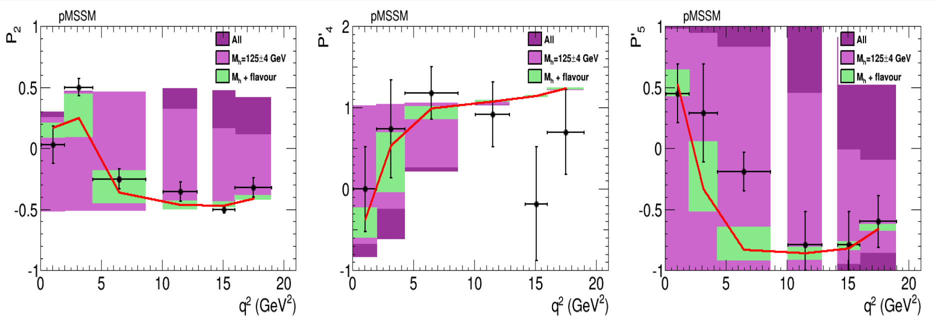
$$\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{\alpha}{4\pi} \mathcal{C}_9 = \frac{c}{\Lambda^2}$$

For $\mathcal{C}_9 \sim 1$ the NP scale Λ would be:

Type of NP	coupling	c	NP scale
Tree-level flavor-generic	$g \sim 1$	~ 1	$\Lambda \sim 38\,\mathrm{TeV}$
Tree-level flavor-CKMish	$g \sim 1$	$\sim V_{tb} V_{ts}^*$	$\Lambda \sim 8\,\mathrm{TeV}$
Tree-level flavor-generic	$g \sim 0.1$	~ 0.01	$\Lambda \sim 3.8\,\mathrm{TeV}$
Loop-level flavor-generic	$g \sim 1$	$\sim \frac{1}{(4\pi)^2}$	$\Lambda \sim 3\,\mathrm{TeV}$
Loop-level flavor-CKMish	$g \sim 1$	$\sim \frac{V_{tb} V_{ts}^*}{(4\pi)^2}$	$\Lambda \sim 600\,\mathrm{GeV}$



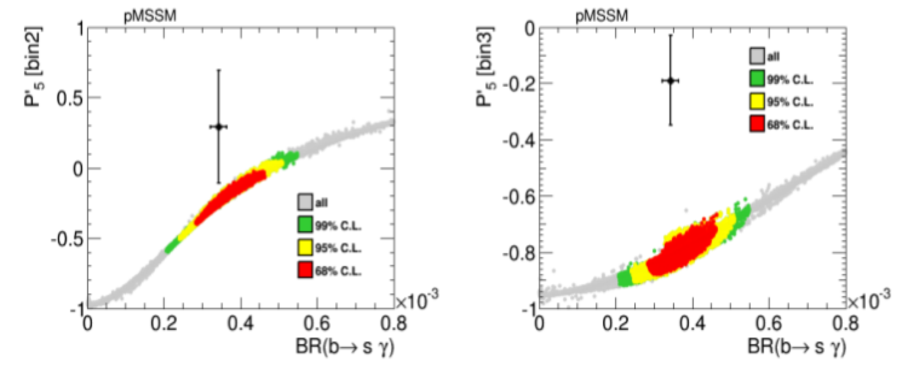
$B \to K^* \mu\mu$ Optimized Observables in the MSSM



⋆ The MSSM has enough freedom to account for all $B \to K^* \mu\mu$ measurements individually and M_H and direct exclusion limits.

⋆ The tension comes exclusively from correlation among flavor observables

F. Mahmoudi, S. Neshatpour, JV, 1401.2145



P_5' is extremely difficult to reproduce in the MSSM, because:

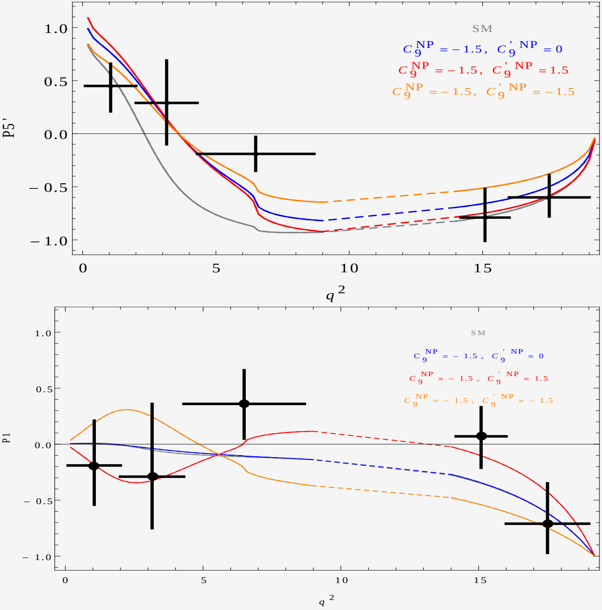
- Large values of $\mathcal{C}_9$ are correlated to large values of other coefficients.
- Large values of $\mathcal{C}_7$ can do it, but are excluded by $B \to X_s\gamma$.

Some References

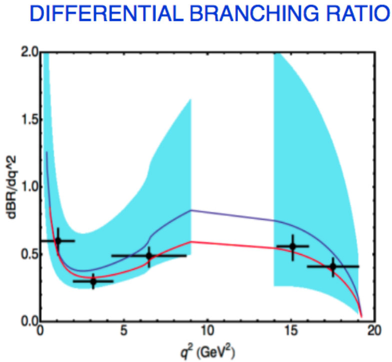
- The LHCb papers on the $B \to K^* \mu\mu$ angular analysis:
  LHCb collaboration, 1304.6325[hep-ex], 1308.1707[hep-ex]
- Statement and analysis of the "$B \to K^* \mu\mu$ Anomaly":
  Descotes-Genon, Matias, JV, 1307.5683[hep-ph]
- Definition of the *Optimised observables* and SM predictions:
  Matias, Mescia, Ramon, JV, 1202.4266[hep-ph]
  Descotes-Genon, Matias, Ramon, JV, 1207.2753[hep-ph]
  Descotes-Genon, Hurth, Matias, JV, 1303.5794[hep-ph]
- Further papers addressing the Anomaly:
  Altmannshofer, Straub, 1308.1501[hep-ph]
  Buras, Girrbach, 1309.2466[hep-ph]
  Beaujean, Bobeth, van Dyk, 1310.2478[hep-ph]
  Gauld, Goertz, Haisch, 1308.1959[hep-ph], 1310.1082[hep-ph]
  Horgan, Liu, Meinel, Wingate, 1310.3887[hep-ph]
  Datta, Duraisamy, Gosh, 1310.1937[hep-ph]
  Mahmoudi, Neshatpour, JV, 1401.2145[hep-ph]
  Altmannshofer, Gori, Pospelov, Yavin, 1403.1269[hep-ph]
  Biancofiore, Collangelo, De Fazio, 1403.2944[hep-ph]
- Form factors and charm-loop effects:
  Khodjamirian, Mannel, Pivovarov, Wang, 1006.4945[hep-ph]
  Horgan, Liu, Meinel, Wingate, 1310.3722[hep-lat]
- The theory of $B \to K^* \ell\ell$ at large and low recoil:
  Beneke, Feldmann, Seidel, 0106067[hep-ph], 0412400[hep-ph]
  Beylich, Buchalla, Feldmann, 1101.5118[hep-ph]
  Grinstein, Pirjol, 0404250[hep-ph]

Backup Slides

$\mathcal{C}_9$ - $\mathcal{C}_9'$ Scenario

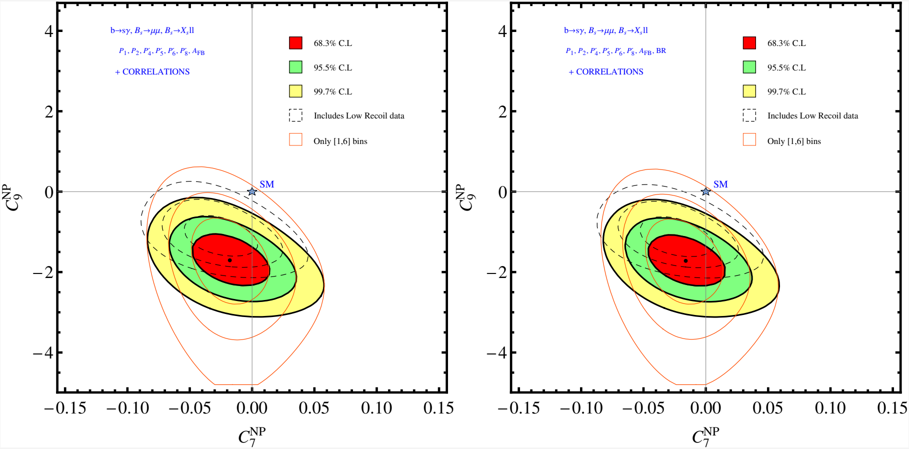


$B \to K^*\mu^+\mu^-$ Branching Ratio

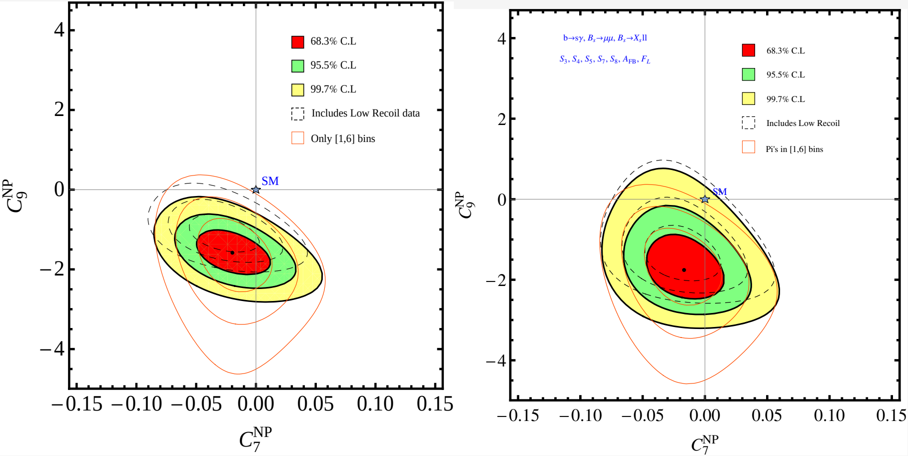


where the blue curve is SM and the red curve corresponds to $C_9^{NP} = -1.5$. Interestingly the central value it goes in the right direction, but given the error bars all is consistent with data.

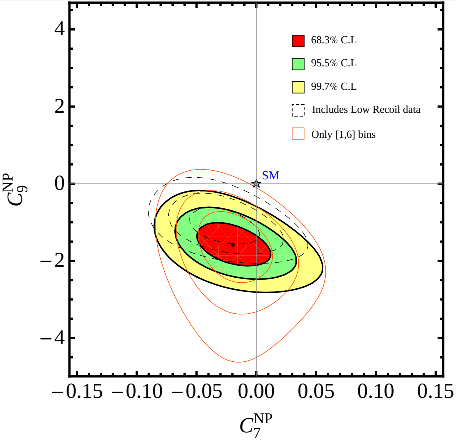
Experimental correlations & Branching Ratio

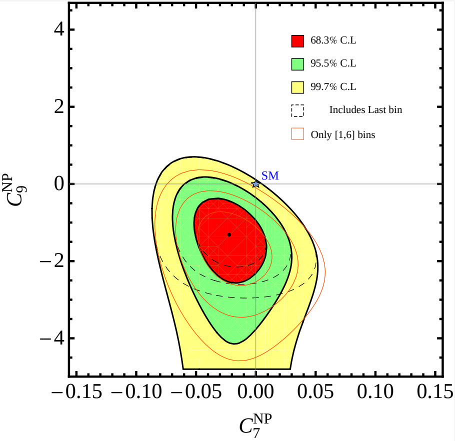


Fit to Form-Factor-dependent observables S_i



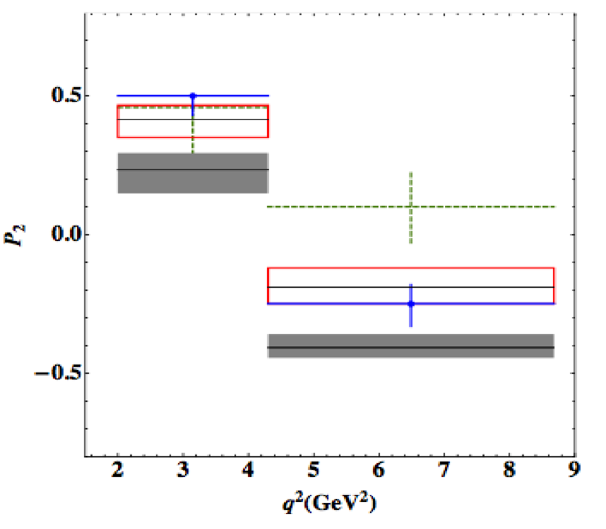
Excluding the [4.3,8.68] bin





Relationship between P_2 and P_5'

$$P_2 = \frac{1}{2}\left[P_4' P_5' + \frac{1}{\beta}\sqrt{(-1 + P_1 + P_4'^2)(-1 - P_1 + \beta^2 P_5'^2)}\right]$$



$B^+ \to K^+\mu\mu$ @ 2012