GRBs

Are gamma-ray bursts the sources of ultra-high energy cosmic rays?

UHECRs

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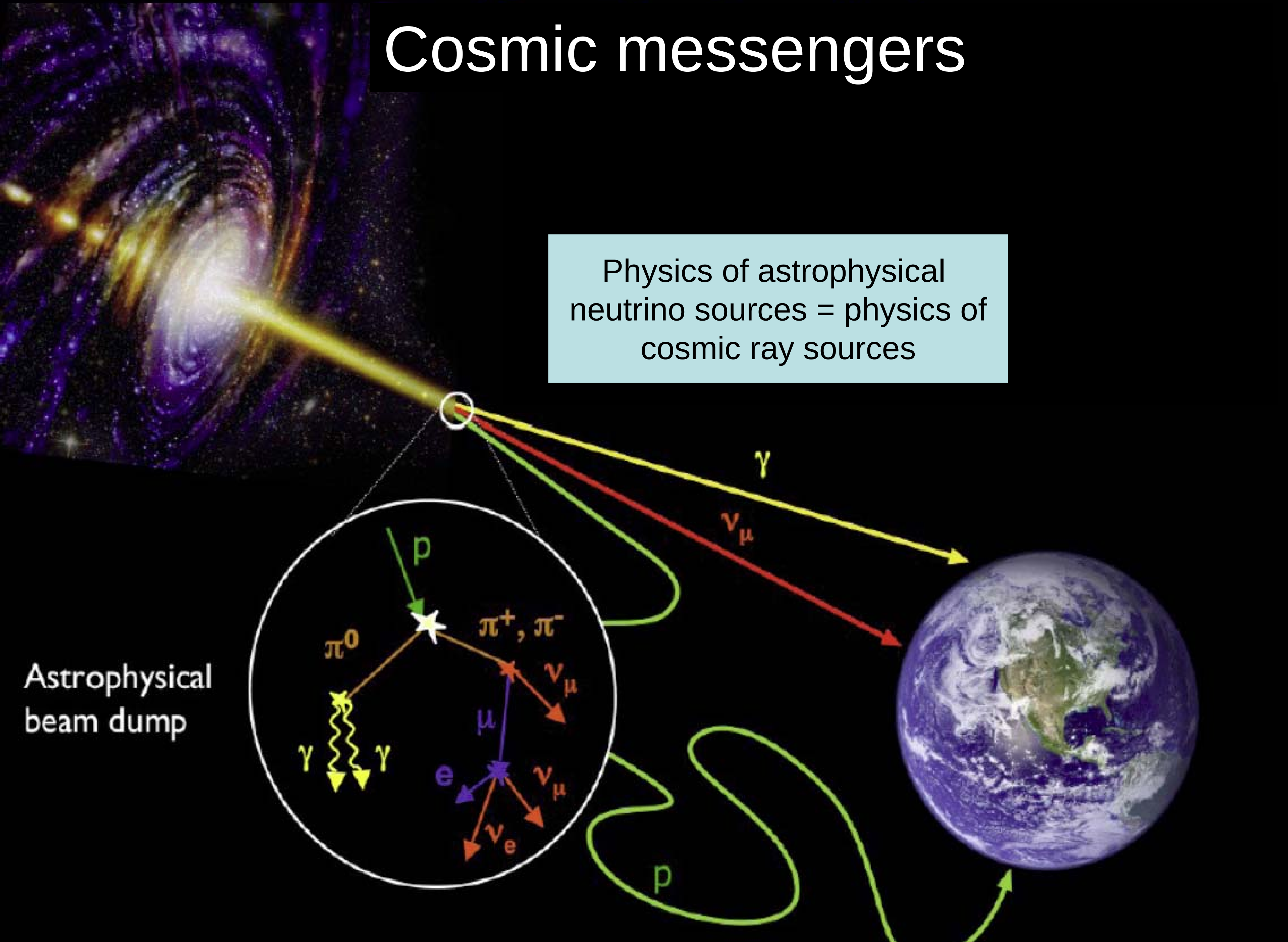
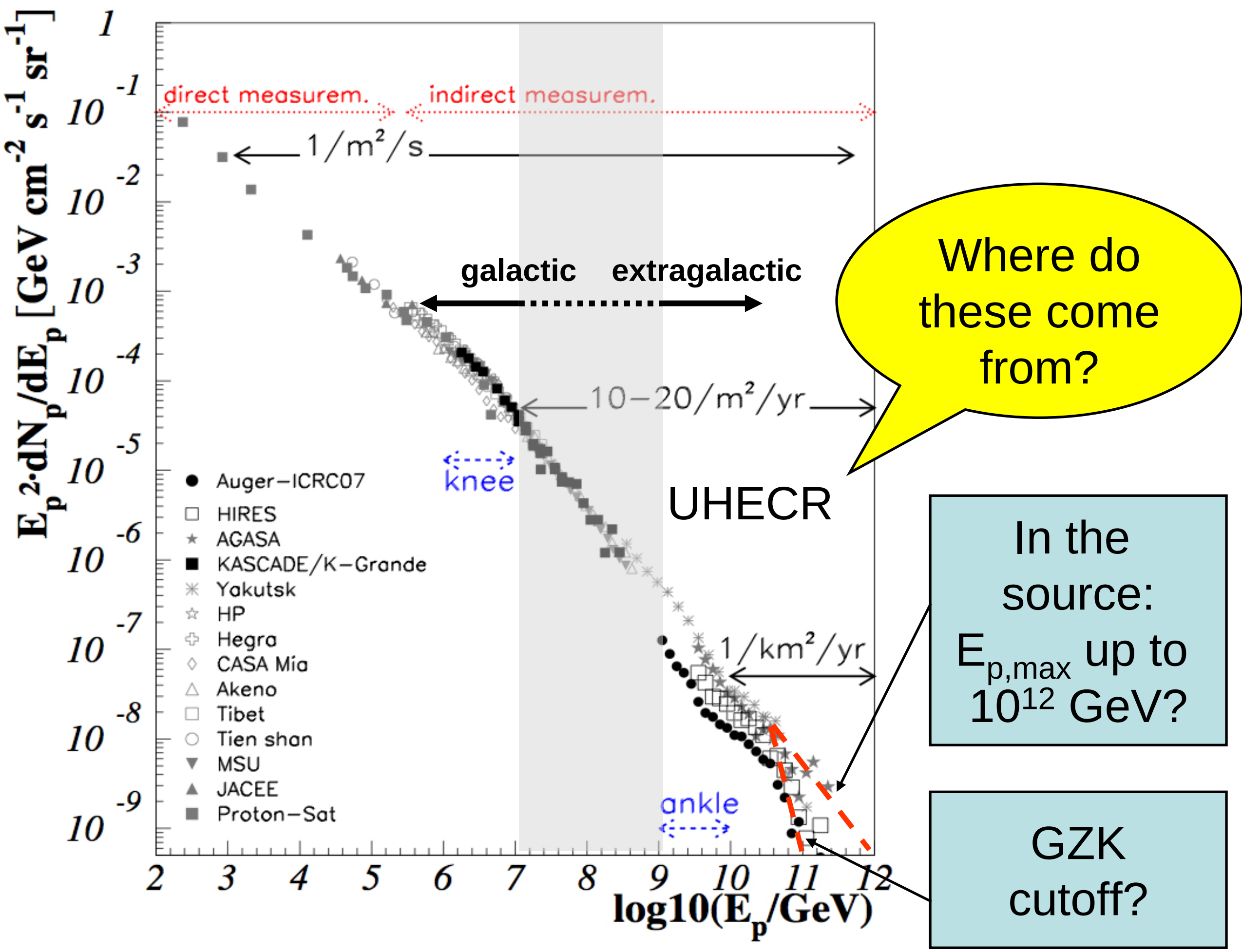
Cosmic messengers

Physics of astrophysical neutrino sources = physics of cosmic ray sources



Cosmic ray observations

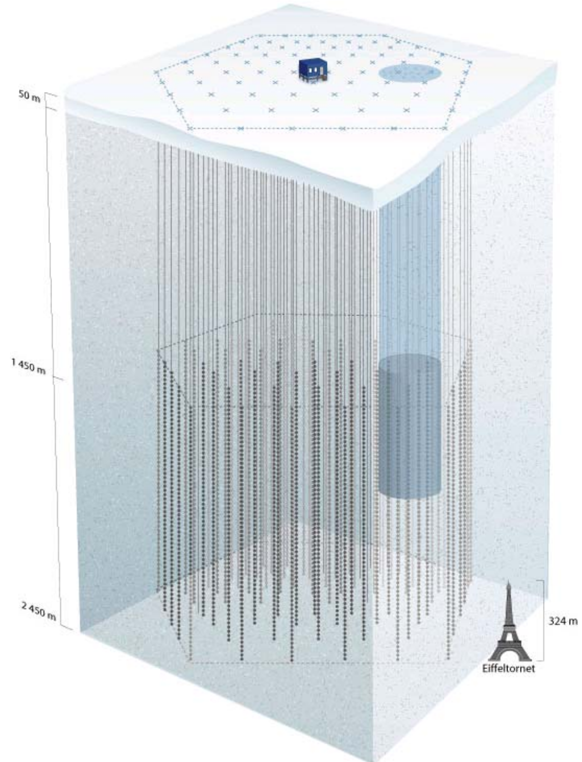
- Observation of cosmic rays: **need to accelerate protons/nuclei somewhere**
- The same sources should produce neutrinos:
  - in the source (pp, pγ interactions)
  - Proton (E > 6 10^{10} GeV) on CMB ⇨ GZK cutoff + cosmogenic neutrino flux

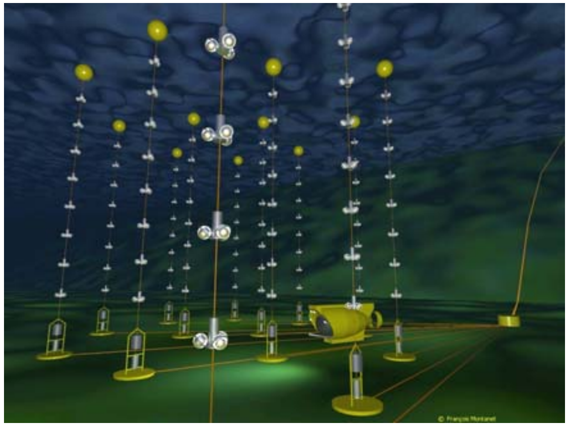


Neutrino detection:

Neutrino telescopes

- Example: IceCube at South Pole. Detector material: ~ 1 km^3 antarctic ice
- Completed 2010/11 (86 strings)
- Recent major successes:
  - Constraints on GRBs Nature 484 (2012) 351
  - 28 events in the TeV-PeV range Science 342 (2013) 1242856
- **Neutrinos established as messengers of the high-energy universe!**





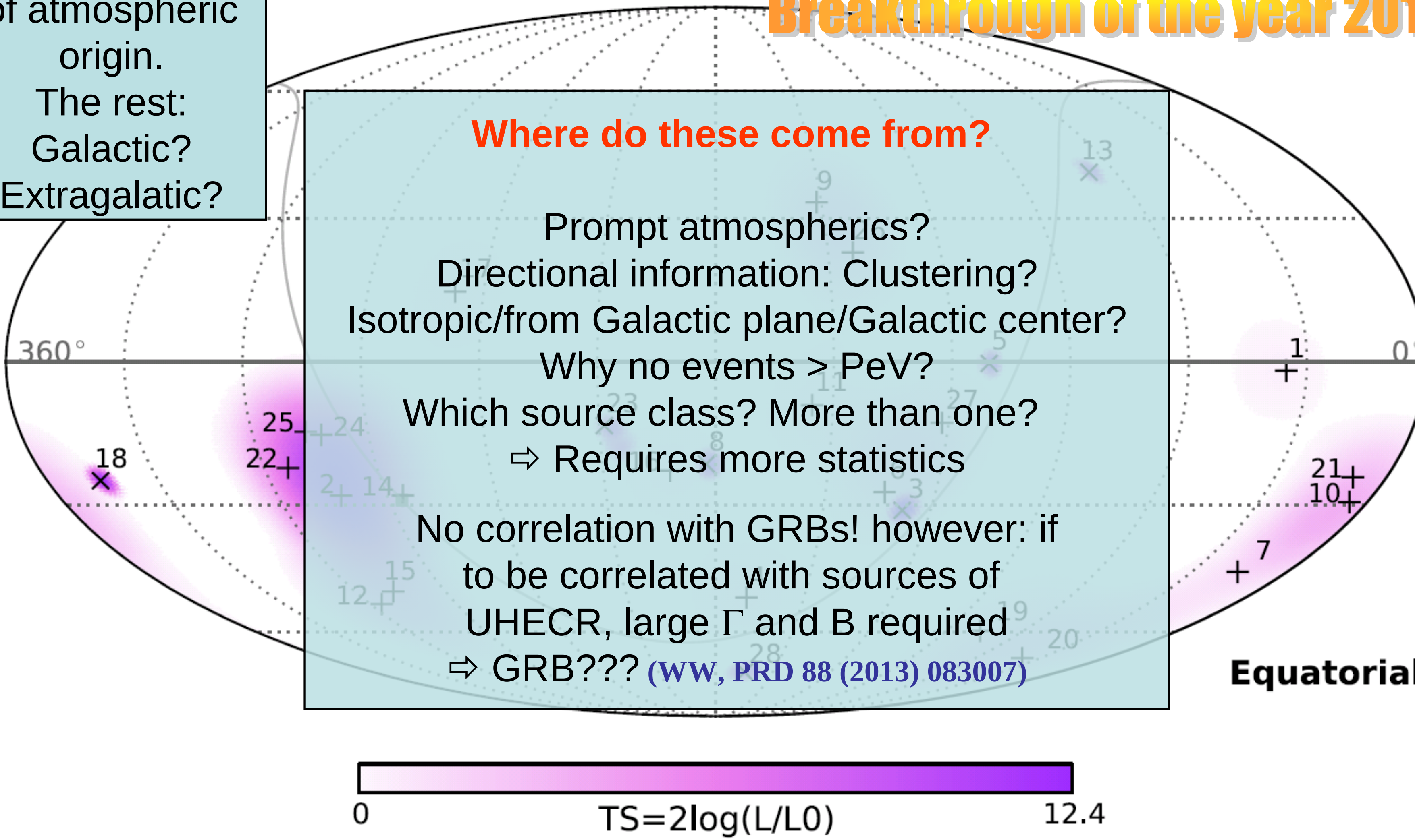
http://icecube.wisc.edu/

http://antares.in2p3.fr/

Neutrinos in the TeV-PeV range

~ 11 events
of atmospheric
origin.
The rest:
Galactic?
Extragalatic?

Physics World
Breakthrough of the year 2013



Where do these come from?

Prompt atmospherics?
Directional information: Clustering?
Isotropic/from Galactic plane/Galactic center?
Why no events > PeV?
Which source class? More than one?
⇨ Requires more statistics

No correlation with GRBs! however: if
to be correlated with sources of
UHECR, large Γ and B required
⇨ GRB??? **(WW, PRD 88 (2013) 083007)**

Science 342 (2013) 1242856

The two paradigms for extragalactic sources:
AGNs and GRBs

- Active Galactic Nuclei (AGN blazars)
  - Relativistic jets ejected from central engine (black hole?)
  - Continuous emission, with time-variability
- Gamma-Ray Bursts (GRBs): transients
  - Relativistically expanding fireball/jet
  - Neutrino production e. g. in prompt phase
    (Waxman, Bahcall, 1997)

Cosmic Rays: 100 years of mystery

2012-04-18

Using data from the IceCube Neutrino Observatory, astrophysicists Nathan Whitehorn and Pete Redl searched for neutrinos coming from the direction of known GRBs. And they found nothing.

Their result, appearing today in the journal Nature, challenges one of the two leading theories for the origin of the highest energy cosmic rays.

Nature 484 (2012) 351

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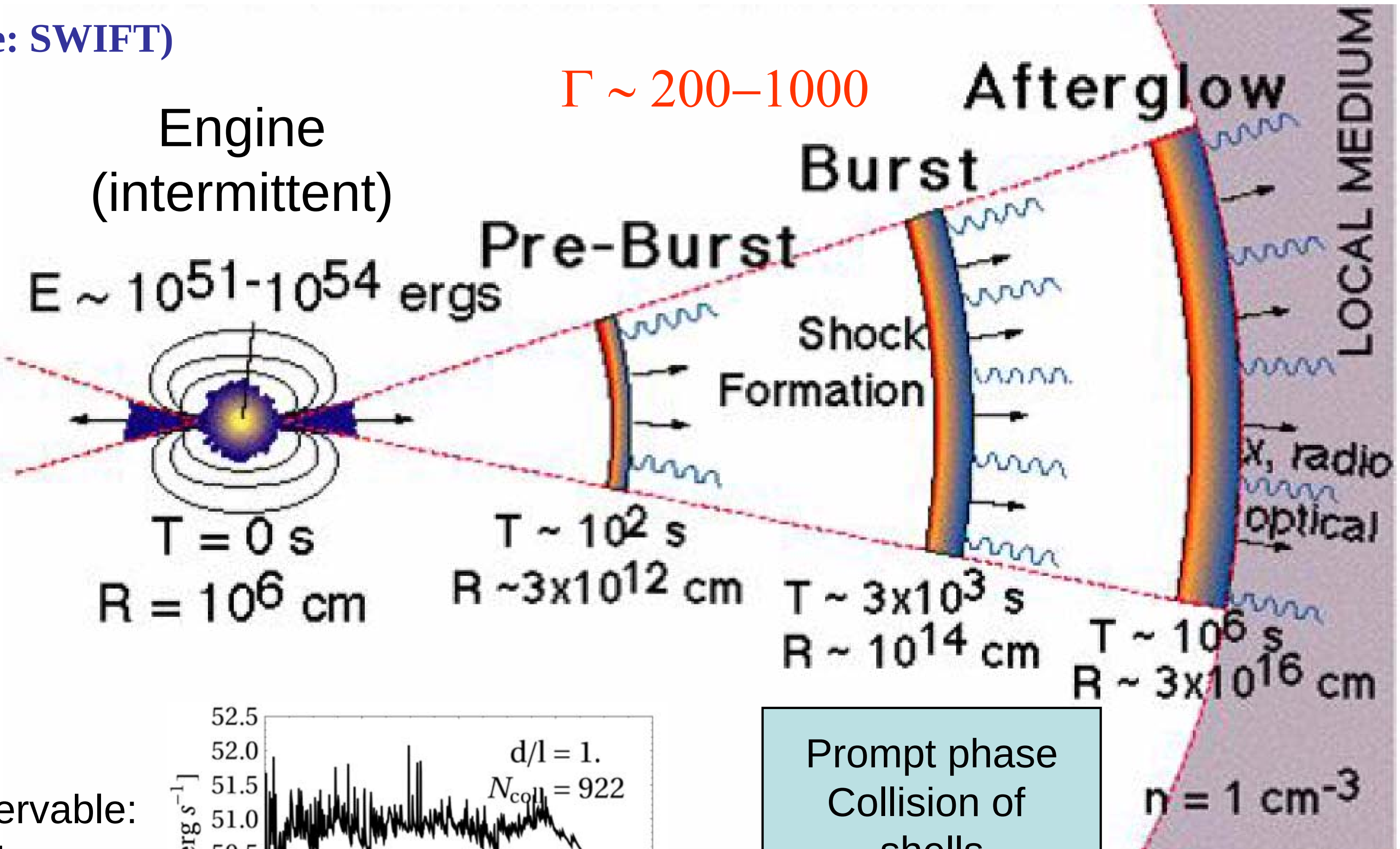
Internal shock model

(Source: SWIFT)

$\Gamma \sim 200–1000$

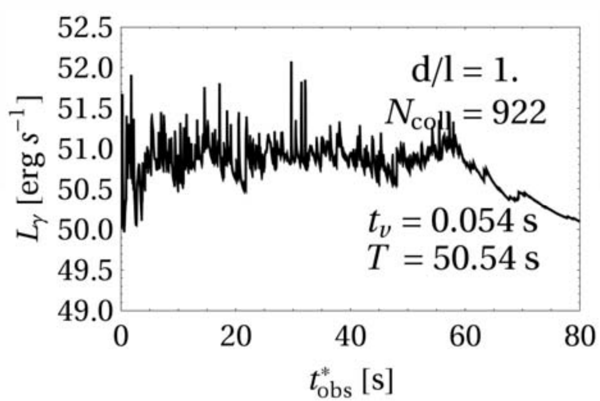
Engine
(intermittent)

“Isotropic
equivalent
energy“



Observable:
Light curves

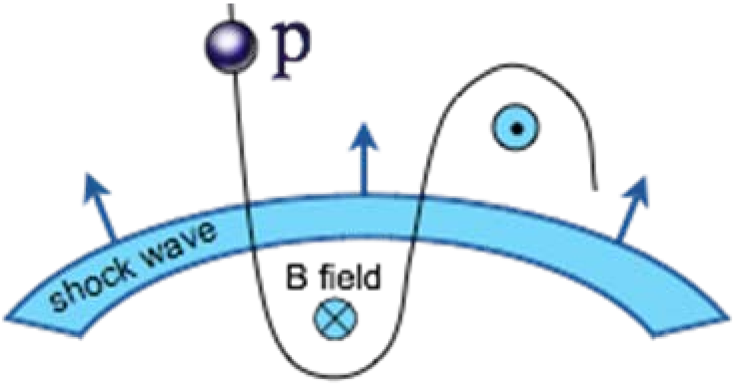
**(Simulation by
M. Bustamante)**



→

Prompt phase
Collision of
shells
⇨ Shocks
⇨ Particle acc.

Fermi shock acceleration



- Fractional energy gain per cycle: η
- Escape probability per cycle: P_{esc}

- Yields a non-thermal power law spectrum $\sim E^{-(P_{esc}/\eta+1)}$
- $P_{esc}/\eta \sim 3/(\chi-1)$ in shock acc., where χ is the compression ratio of the shock
- $\chi \sim 4$ for a strong shock ⇨ $P_{esc}/\eta \sim 1$ and E^{-2} is the typical “textbook“ spectrum

Simulation of cosmic ray and neutrino sources

(focus on proton composition …)

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Cosmic ray source

(illustrative proton-only scenario, pγ interactions)

If neutrons can escape:
Source of cosmic rays

Neutrinos produced in
ratio $(\nu_e:\nu_\mu:\nu_\tau)=(1:2:0)$

$$n \rightarrow p + e^- + \bar{\nu}_e$$

$$\pi^+ \rightarrow \mu^+ + \nu_\mu ,$$

$$p + \gamma_{\mathrm{CMB}} \rightarrow \Delta^+ \rightarrow \text{Cosmogenic neutrinos}$$

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

Delta resonance approximation:

$$p + \gamma \rightarrow \Delta^+ \rightarrow \begin{cases} n + \pi^+ & 1/3 \text{ of all cases} \\ p + \pi^0 & 2/3 \text{ of all cases} \end{cases}$$

π^+/π^0 determines ratio between neutrinos and high-E gamma-rays

$$\pi^0 \longrightarrow \gamma + \gamma$$

High energetic gamma-rays;
typically cascade down to lower E

Cosmic messengers

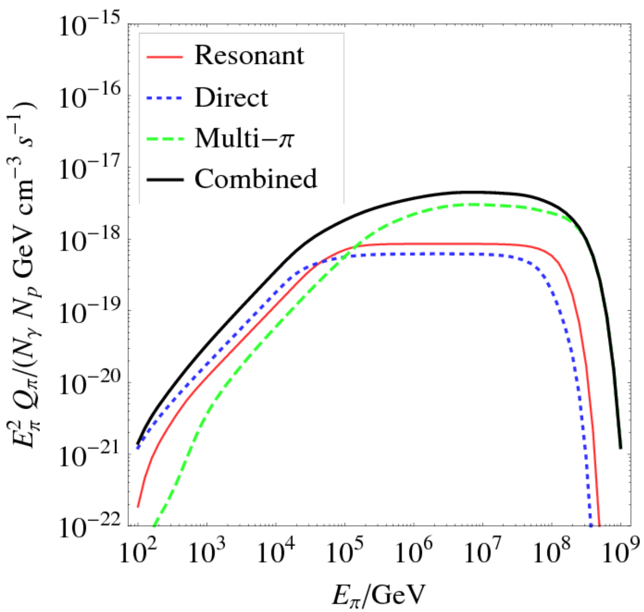
Source simulation: pγ

(particle physics)

- Δ(1232)-resonance approximation:

$$p + \gamma \rightarrow \Delta^+ \rightarrow \begin{cases} n + \pi^+ & 1/3 \text{ of all cases} \\ p + \pi^0 & 2/3 \text{ of all cases} \end{cases}$$

- Limitations:
  - No π^- production; cannot predict π^+/ π^- ratio (Glashow resonance!)
  - High energy processes affect spectral shape (X-sec. dependence!)
  - Low energy processes (t-channel) enhance charged pion production
- Solutions:
  - SC GRB: π^+ iption of physics



from: Hümmer, Rüger, Spanier, Winter, ApJ 721 (2010) 630

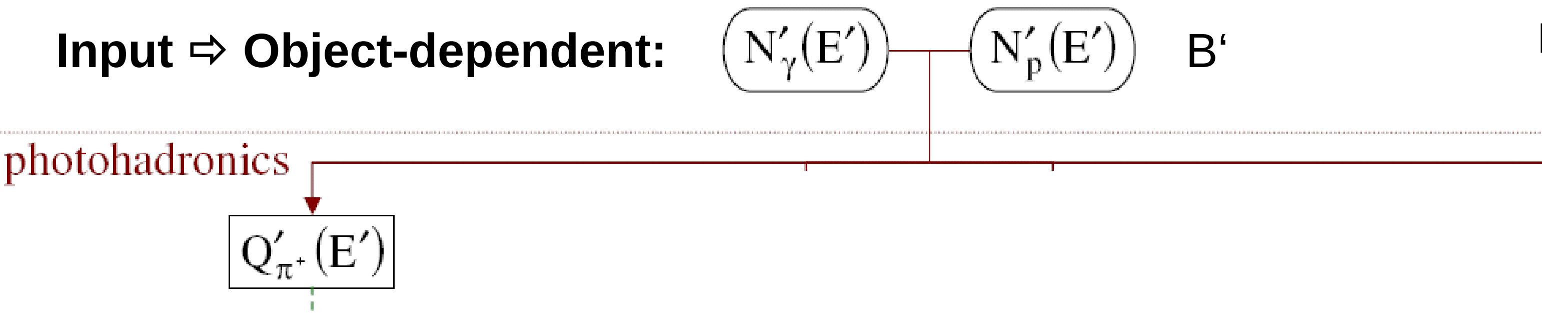
NeuCosmA

“Minimal“ (top down) ν model

Dashed arrows: kinetic equations include cooling and escape

Q(E) [GeV^{-1} cm^{-3} s^{-1}]
per time frame
N(E) [GeV^{-1} cm^{-3}]
steady spectrum

Input ⇨ Object-dependent:



Peculiarity for neutrinos: Secondary cooling

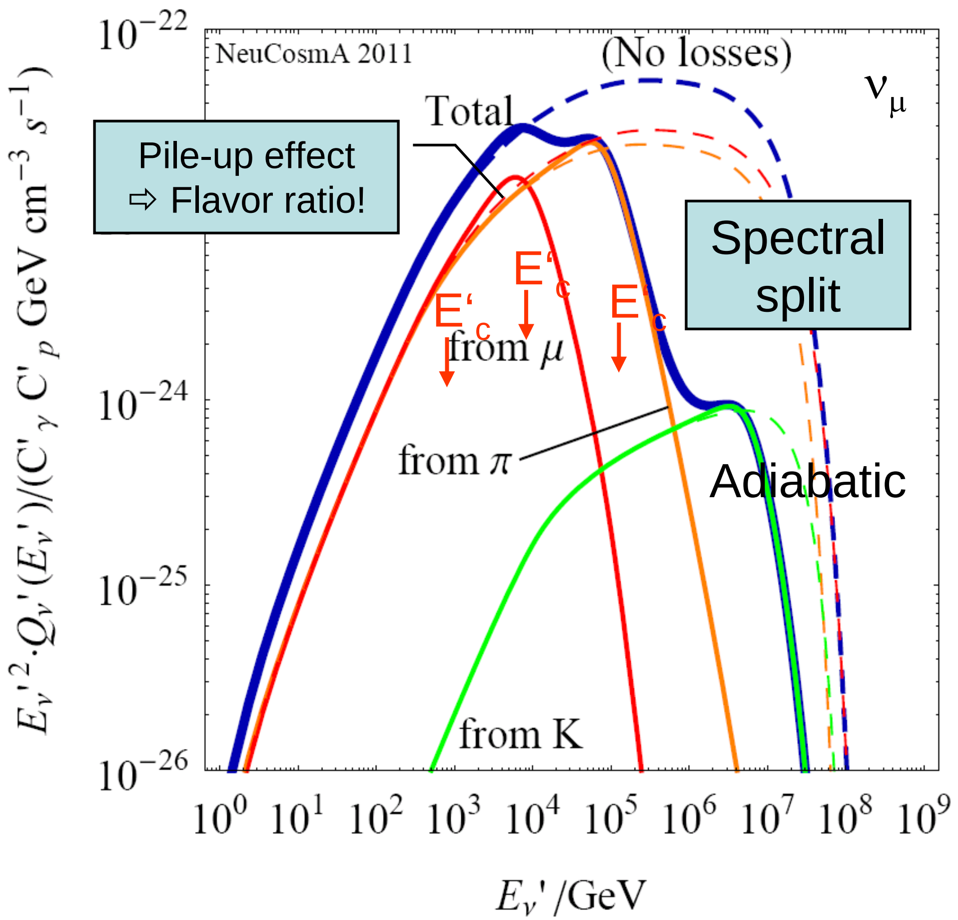
Example: GRB

Secondary spectra (μ, π, K) loss-steepend above critical energy

$$E'_c = \sqrt{\frac{9\pi\epsilon_0 m^5 c^7}{\tau_0 e^4 B'^2}}$$

- E'$_c$ depends on particle physics only (m, τ_0), and **B'**
- Leads to characteristic flavor composition and shape
- Very robust prediction for sources? [e.g. any additional radiation processes mainly affecting the primaries will not affect the flavor composition]

Decay/cooling: charged μ, π, K



Baerwald, Hümmer, Winter, Astropart. Phys. 35 (2012) 508; also: Kashti, Waxman, 2005; Lipari et al, 2007

From the source to the detector: UHECR transport

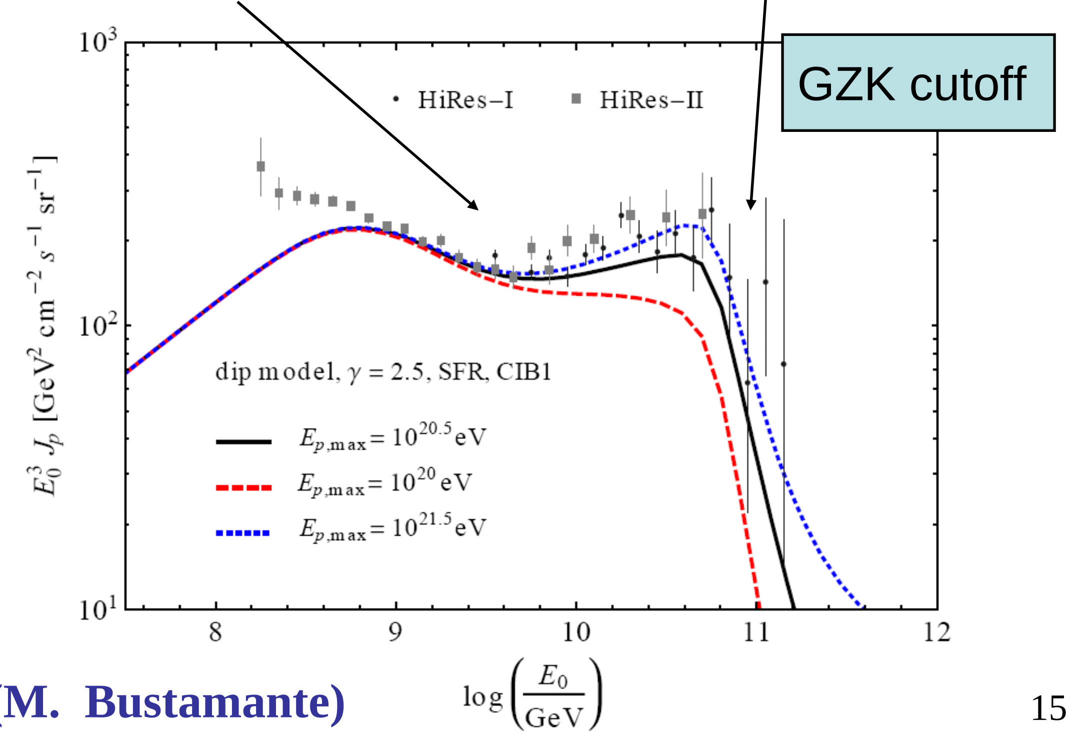
- Kinetic equation for co-moving number density:

$$\dot{Y}_p = \partial_E \left(H E Y_p\right) + \partial_E \left(b_{e^+e^-} Y_p\right) + \partial_E \left(b_{p\gamma} Y_p\right) + \mathcal{L}_{\mathrm{CR}}$$

Expansion of Universe	Pair production **Blumenthal, 1970**	Photohadronics **Hümmer, Rüger, Spanier, Winter, 2010**	CR inj. z-dep!

[here $b=-dE/dt=E\, t^{-1}_{loss}$]

- Energy losses
  ⇨ UHECR must from from our local environment
  (~ 1 Gpc at 10^{10} GeV,
  ~ 50 Mpc at 10^{11} GeV)

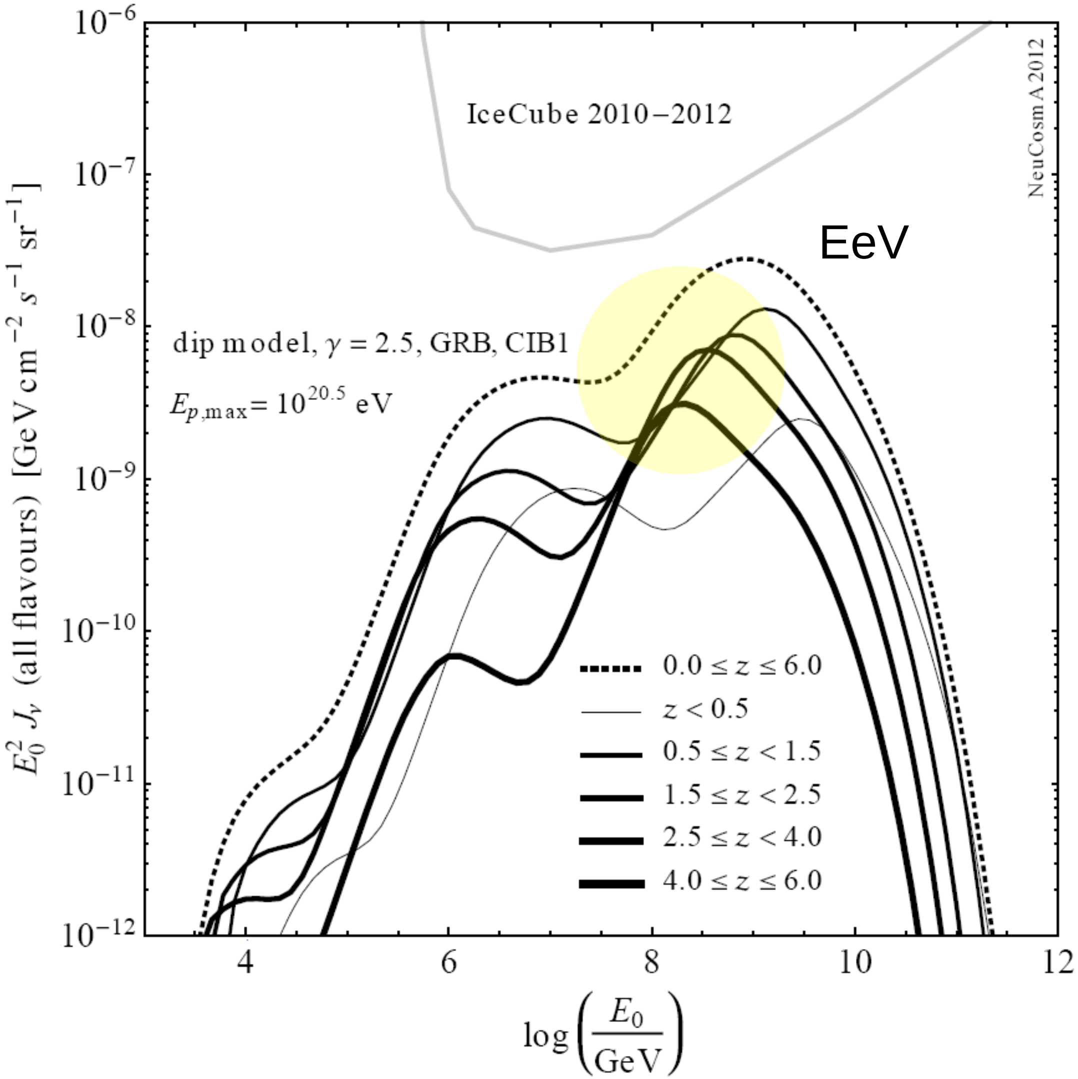


(M. Bustamante)

Cosmogenic neutrinos

$p + \gamma_{CMB} \rightarrow \Delta^{+} \rightarrow$ Cosmogenic neutrinos

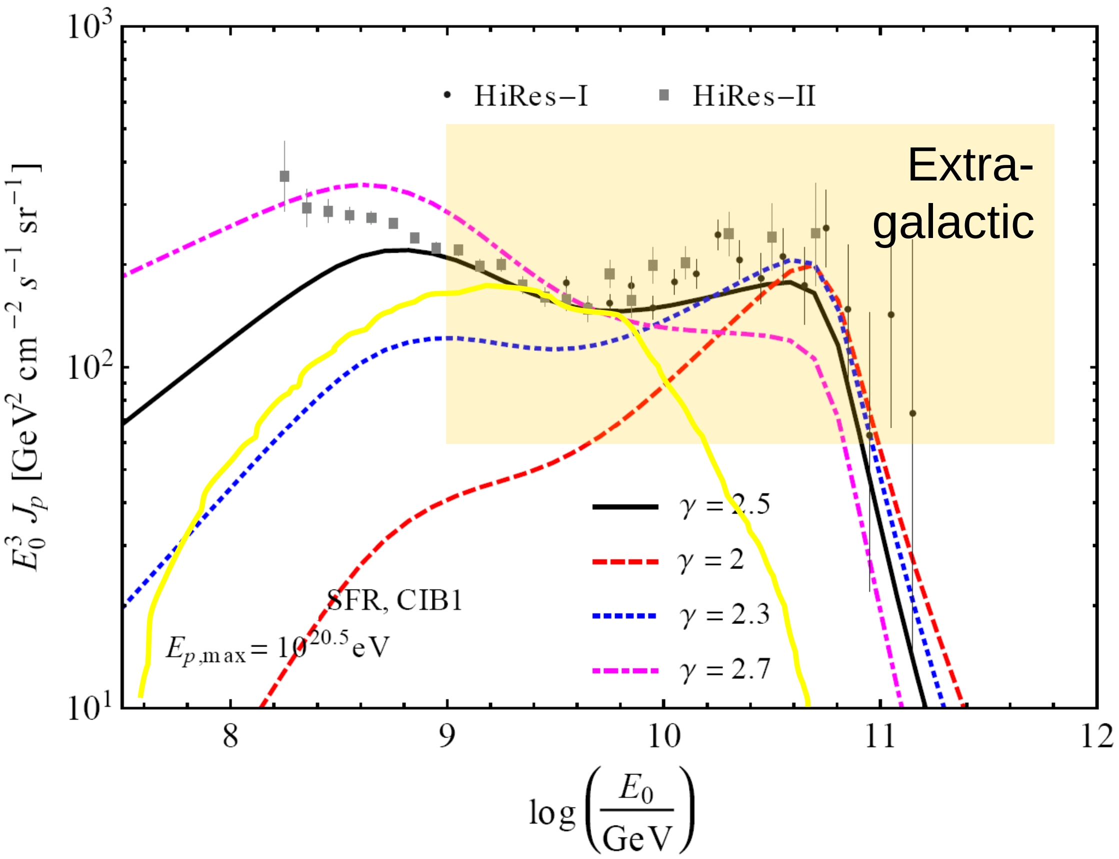
- Prediction depends on maximal proton energy, spectral index γ, source evolution, composition
- Can test UHECR beyond the local environment
- Can test UHECR injection independent of CR production model
  ⇨ constraints on UHECR escape



(courtesy M. Bustamante; see also Kotera, Allard, Olinto, JCAP 1010 (2010) 013)

UHECR transition models

- Transition between Galactic (?) and extragalactic cosmic rays at different energies:



- Ankle model:
  - Injection index γ ~ 2 possible (⇨ Fermi shock acc.)
  - Transition at > 4 EeV
- Dip model:
  - Injection index γ ~ 2.5-2.7 (how?)
  - Transition at ~ 1 EeV
  - Characteristic shape by pair production dip

Figure courtesy M. Bustamante; for a recent review, see Berezinsky, arXiv:1307.4043

Multi-messenger physics with GRBs

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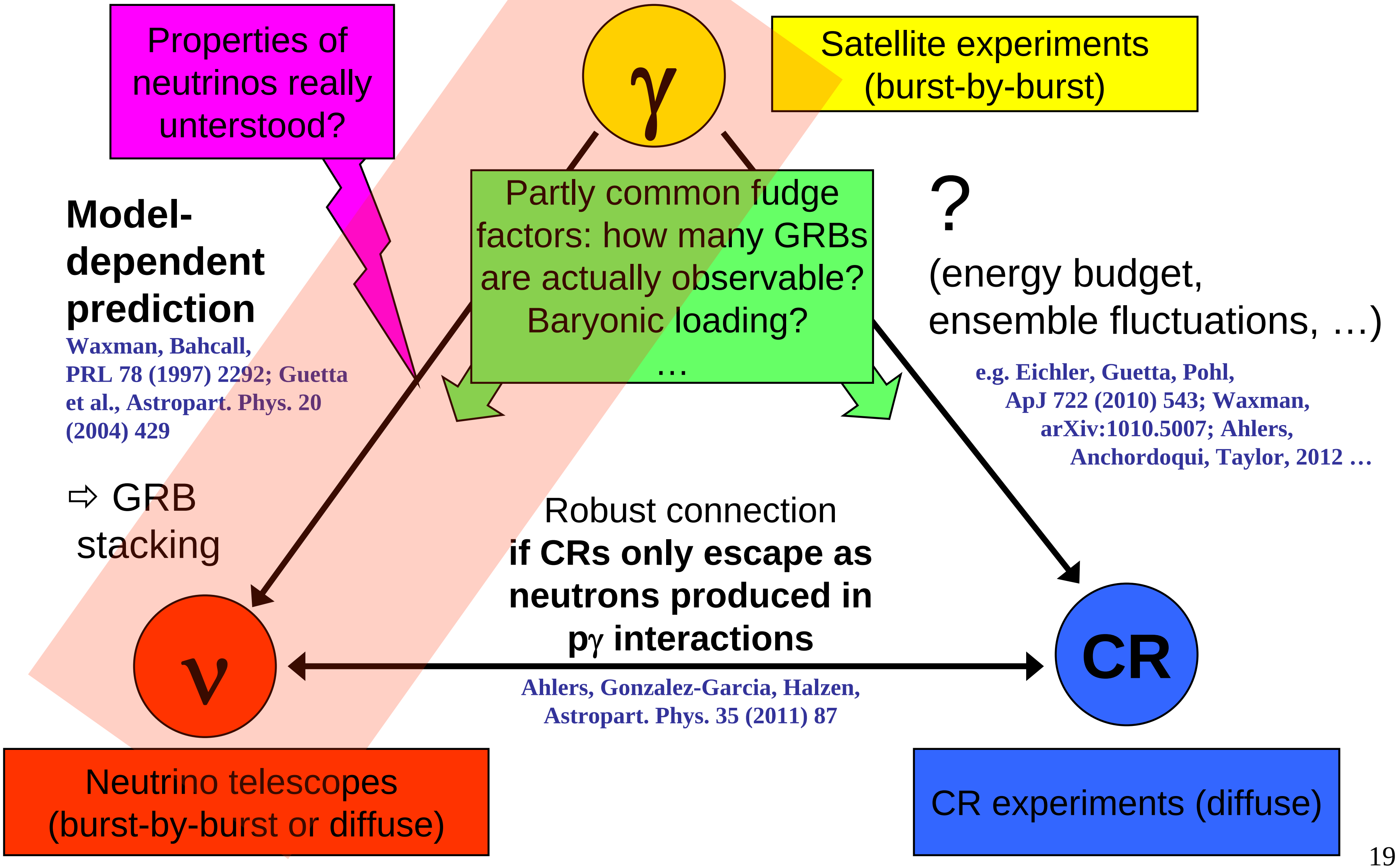
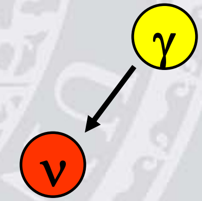
Multi-messenger physics



GRB stacking



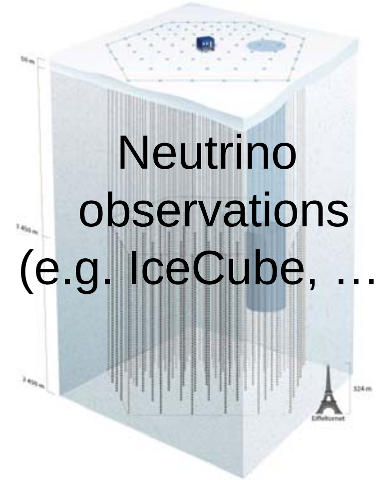
- Idea: Use multi-messenger approach (BG free)

(Source: NASA)

GRB gamma-ray observations
(e.g. Fermi, Swift, etc)

Coincidence!

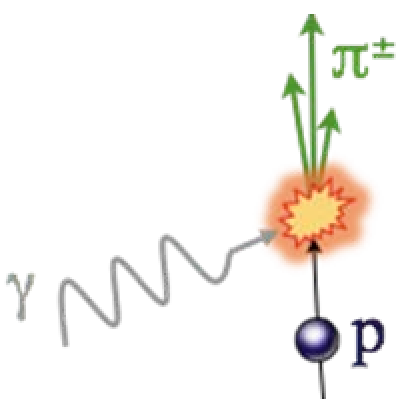
(Source: IceCube)



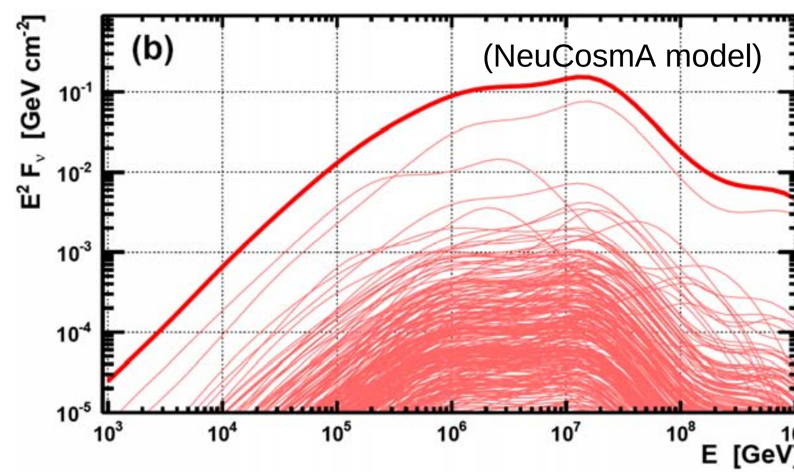
Neutrino
observations
(e.g. IceCube, …)

- Predict neutrino flux from observed photon fluxes event by event

Observed:
broken power law
(Band function)



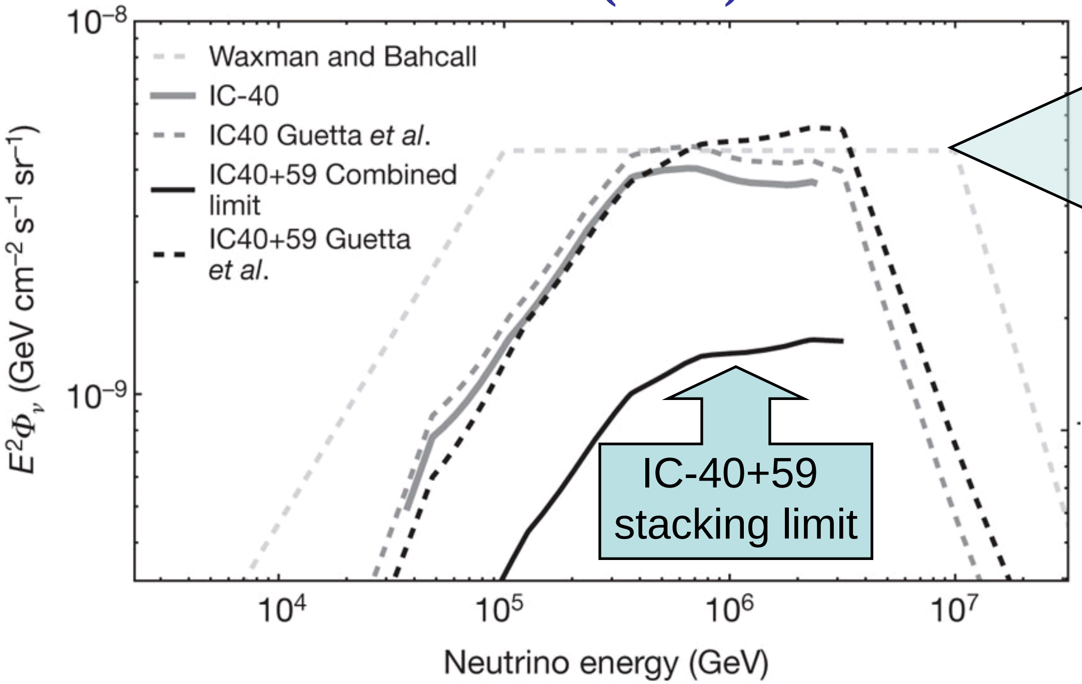
E^{-2} injection



(Example: ANTARES, arXiv:1307.0304)

Gamma-ray burst fireball model:
IC-40+59 data meet generic bounds

Nature 484 (2012) 351



Generic flux based on the assumption that GRBs are the sources of (highest energetic) cosmic rays
(Waxman, Bahcall, 1999; Waxman, 2003; spec. bursts: Guetta et al, 2003)

- Does IceCube really rule out the paradigm that GRBs are the sources of the ultra-high energy cosmic rays?

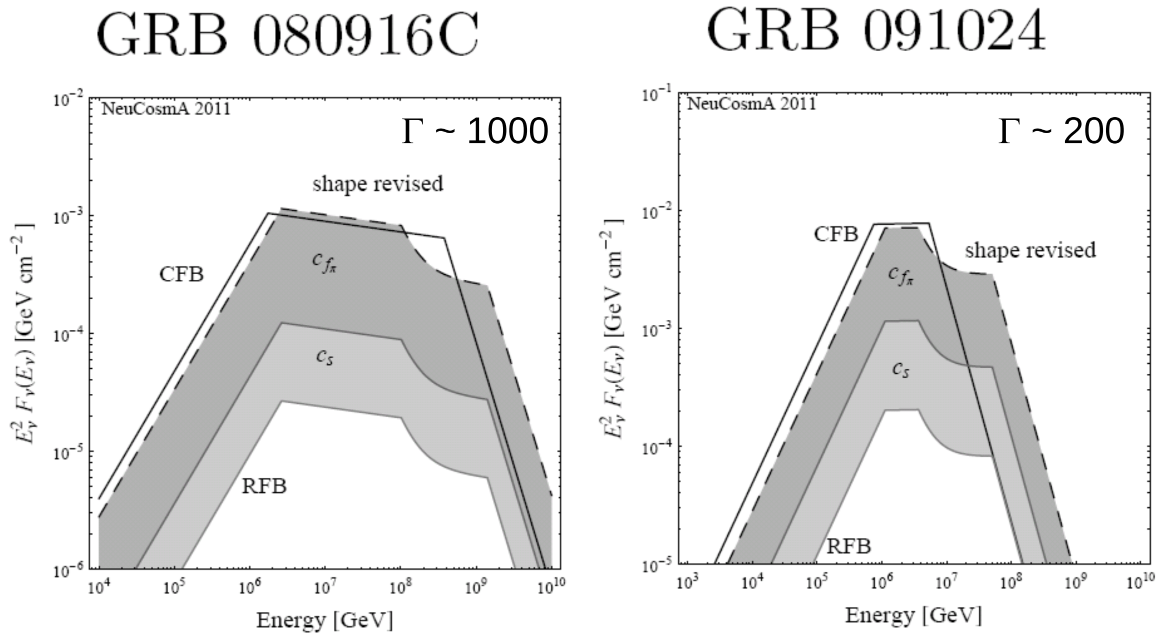
Revision of neutrino flux predictions

Analytical recomputation of IceCube method (CFB):

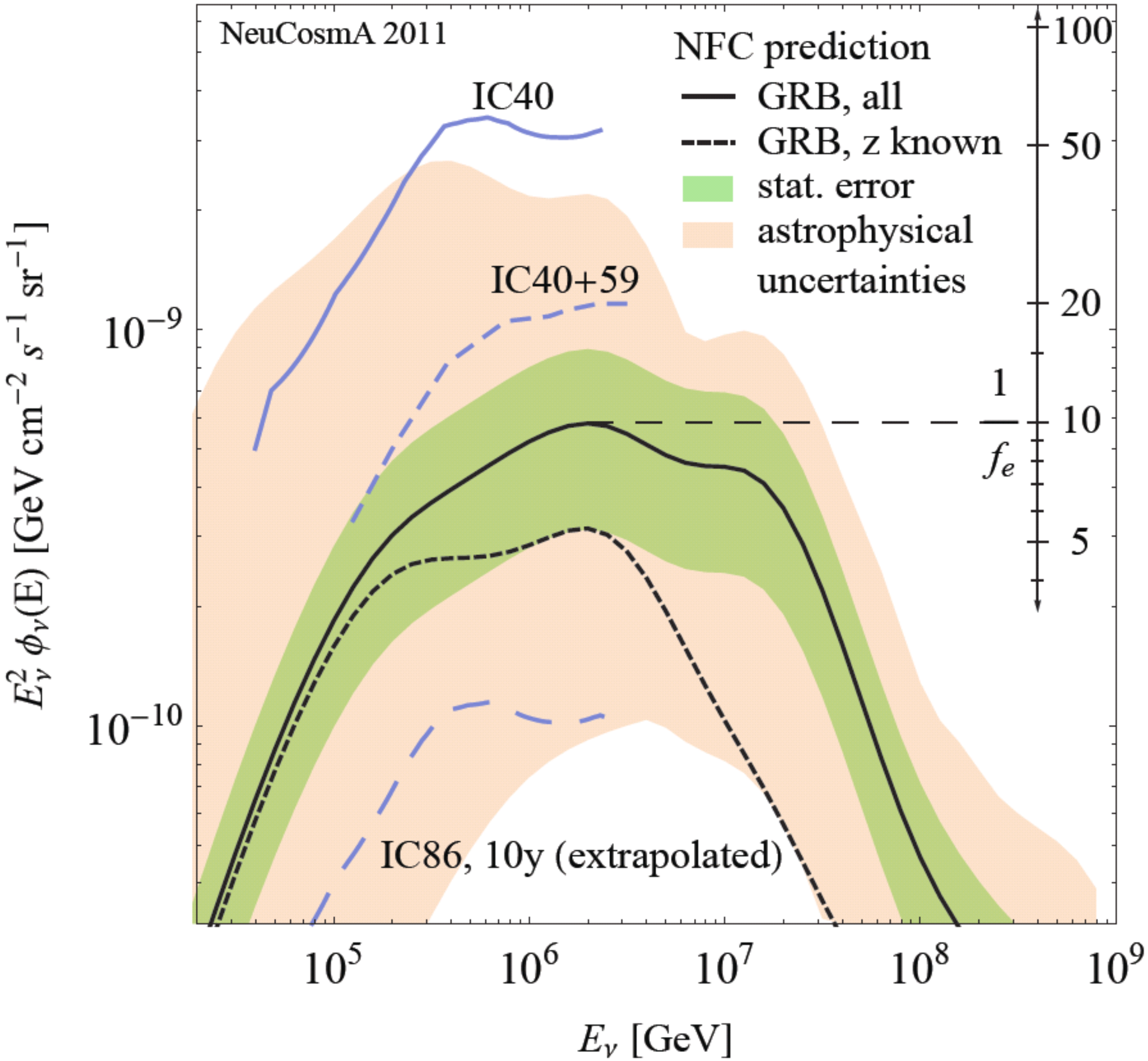
$c_{f\pi}$: corrections to pion production efficiency

c_S: secondary cooling and energy-dependence of proton mean free path

(see also Li, 2012, PRD)



Quasi-diffuse prediction
(NeuCosmA model)

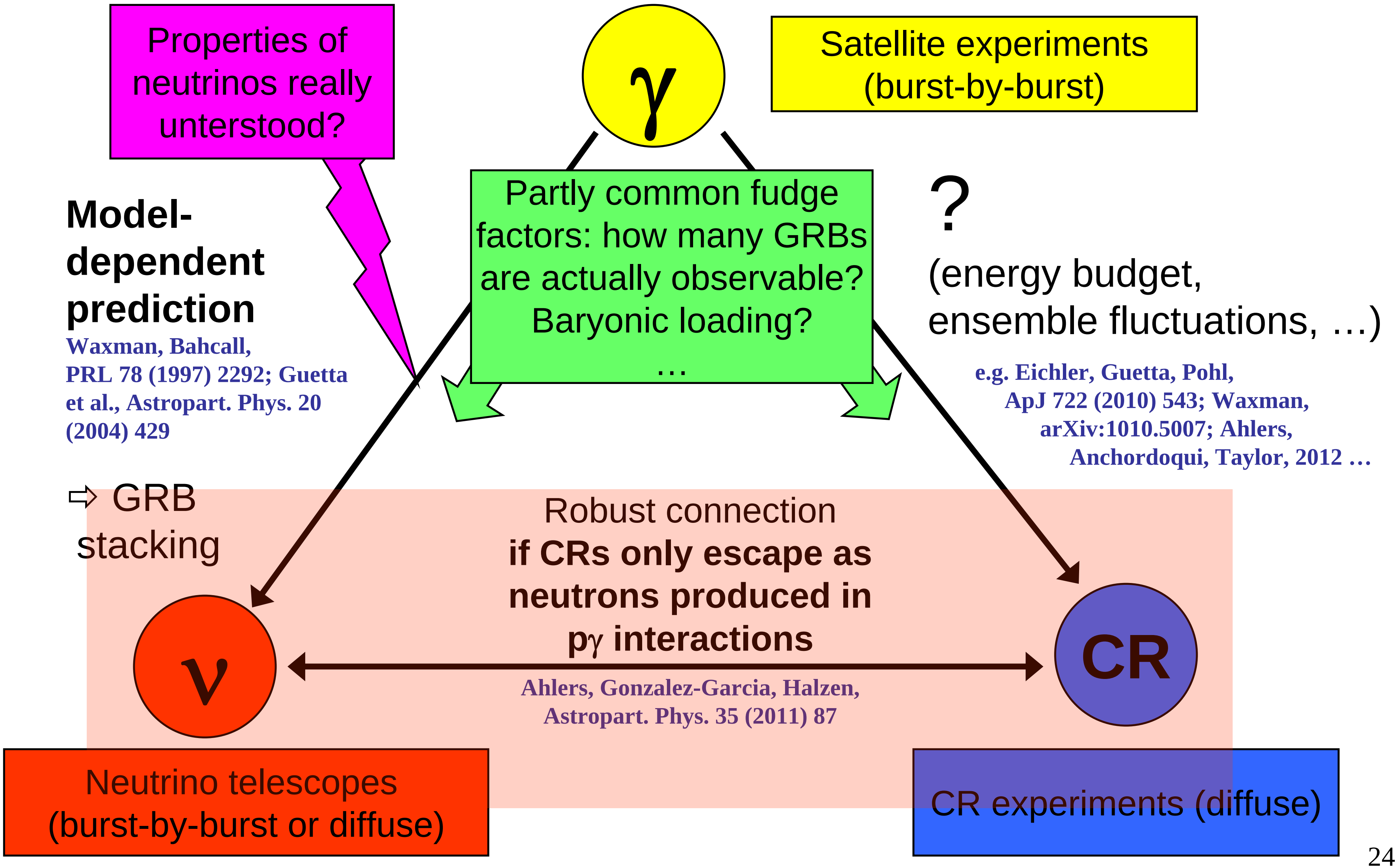


- Numerical fireball model cannot be ruled out yet with IC40+59 for same parameters, bursts, assumptions
- "Fudge" factors:
  - f_e^{-1} Baryonic loading (energy in protons vs. electrons) ⇨ 10?
  - $\dot{N}$ Total number of GRBs in observable universe ⇨ chosen to be 667/year

"Astrophysical uncertainties":
t_v: 0.001s … 0.1s
Γ: 200 …500
α: 1.8 … 2.2
$\varepsilon_e/\varepsilon_B$: 0.1 … 10

(Hümmer, Baerwald, Winter, Phys. Rev. Lett. 108 (2012) 231101)

Neutrinos-cosmic rays



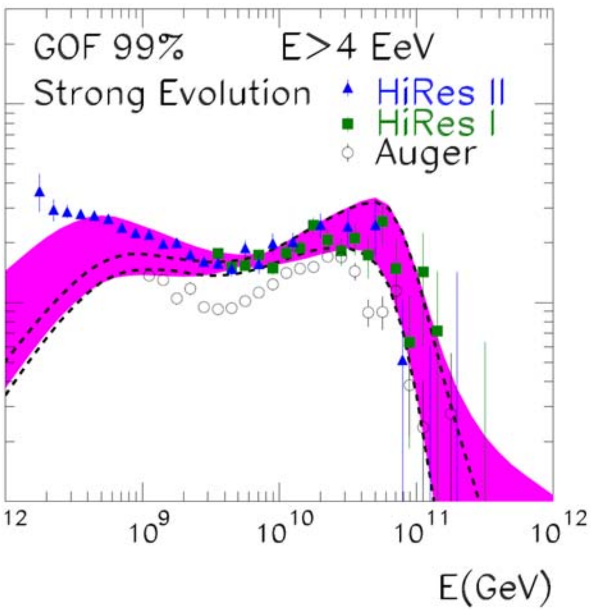
The “neutron model“

ν ⟷ CR

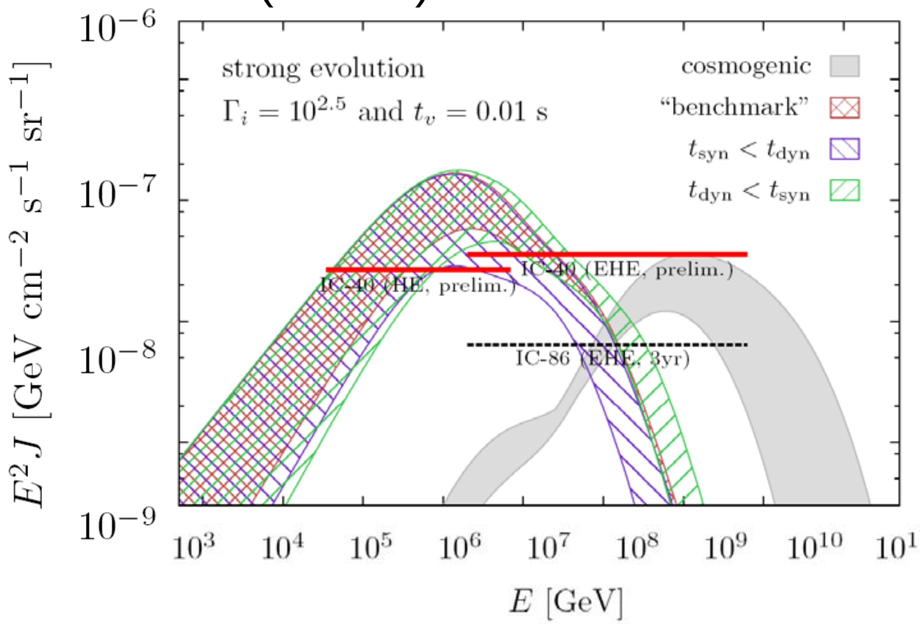
- If charged π and n produced together:

$$p + \gamma \rightarrow \Delta^+ \rightarrow \begin{cases} n + \pi^+ & 1/3 \text{ of all cases} \\ p + \pi^0 & 2/3 \text{ of all cases} \end{cases}$$

Fit to UHECR spectrum



Consequences for (diffuse) neutrino fluxes



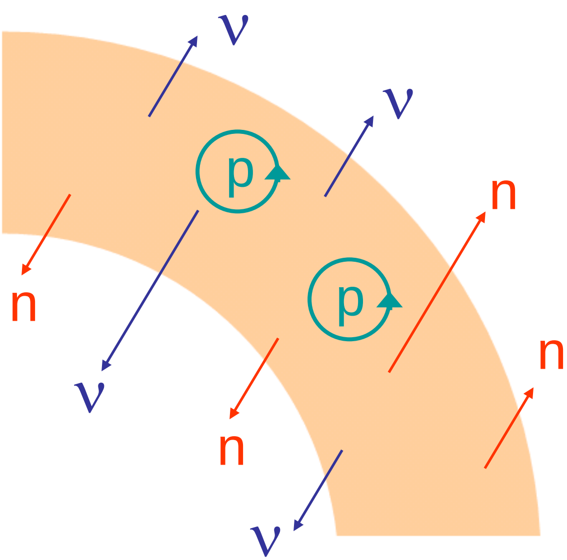
Ahlers, Gonzalez-Garcia, Halzen, Astropart. Phys. 35 (2011) 87

➢ Baryonic loading? CR leakage? Ensemble fluctuations? (Ahlers, Anchordoqui, Taylor, 2012; Kistler, Stanev, Yuksel, 2013; …)

CR escape mechanisms

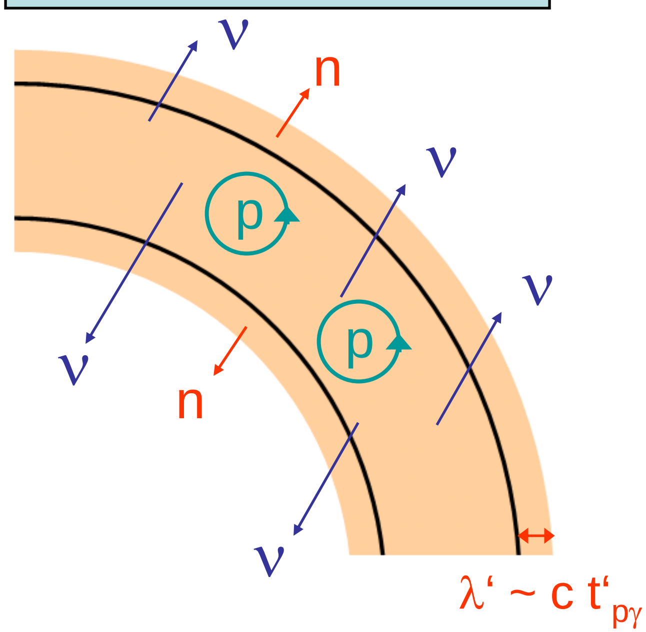
Baerwald, Bustamante, Winter, Astrophys. J. 768 (2013) 186

Optically thin (to neutron escape)



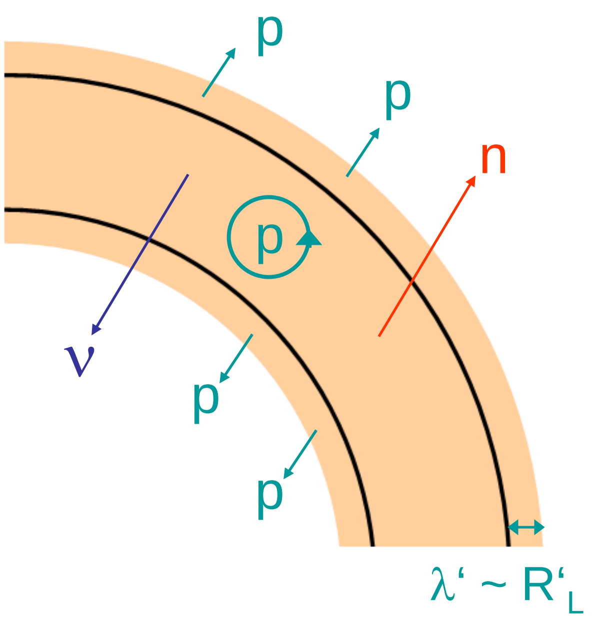
- One neutrino per cosmic ray
- Protons magnetically confined

Optically thick (to neutron escape)



- Neutron escape limited to edge of shells
- Neutrino prod. relatively enhanced

Direct proton escape (UHECR leakage)



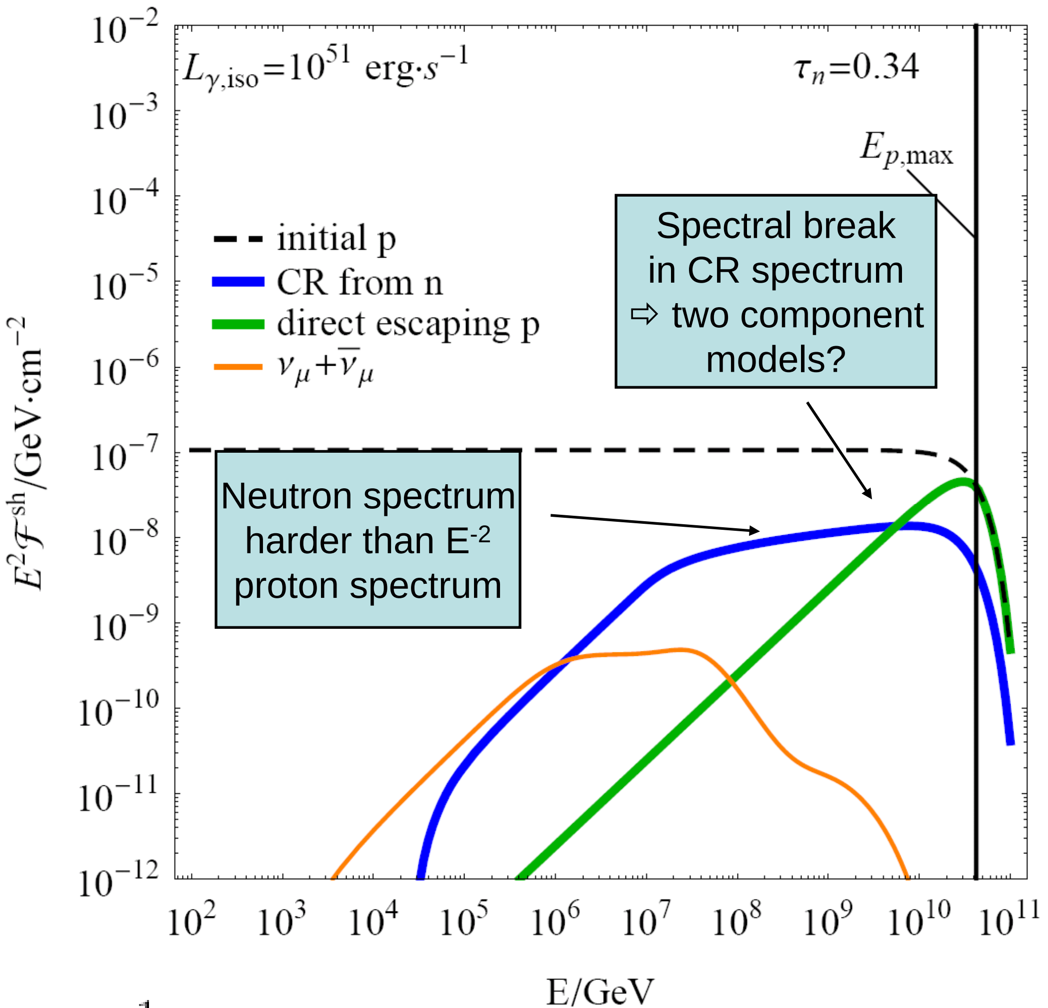
- pγ interaction rate relatively low
- Protons leaking from edges dominate

A typical (?) example

- For high acceleration efficiencies: R'_L can reach shell thickness at highest energies (if $E'_{p,max}$ determined by t'_{dyn})

➢ UHECR from optically thin GRBs will be direct escape-dominated

(Baerwald, Bustamante, Winter, Astrophys.J. 768 (2013) 186)

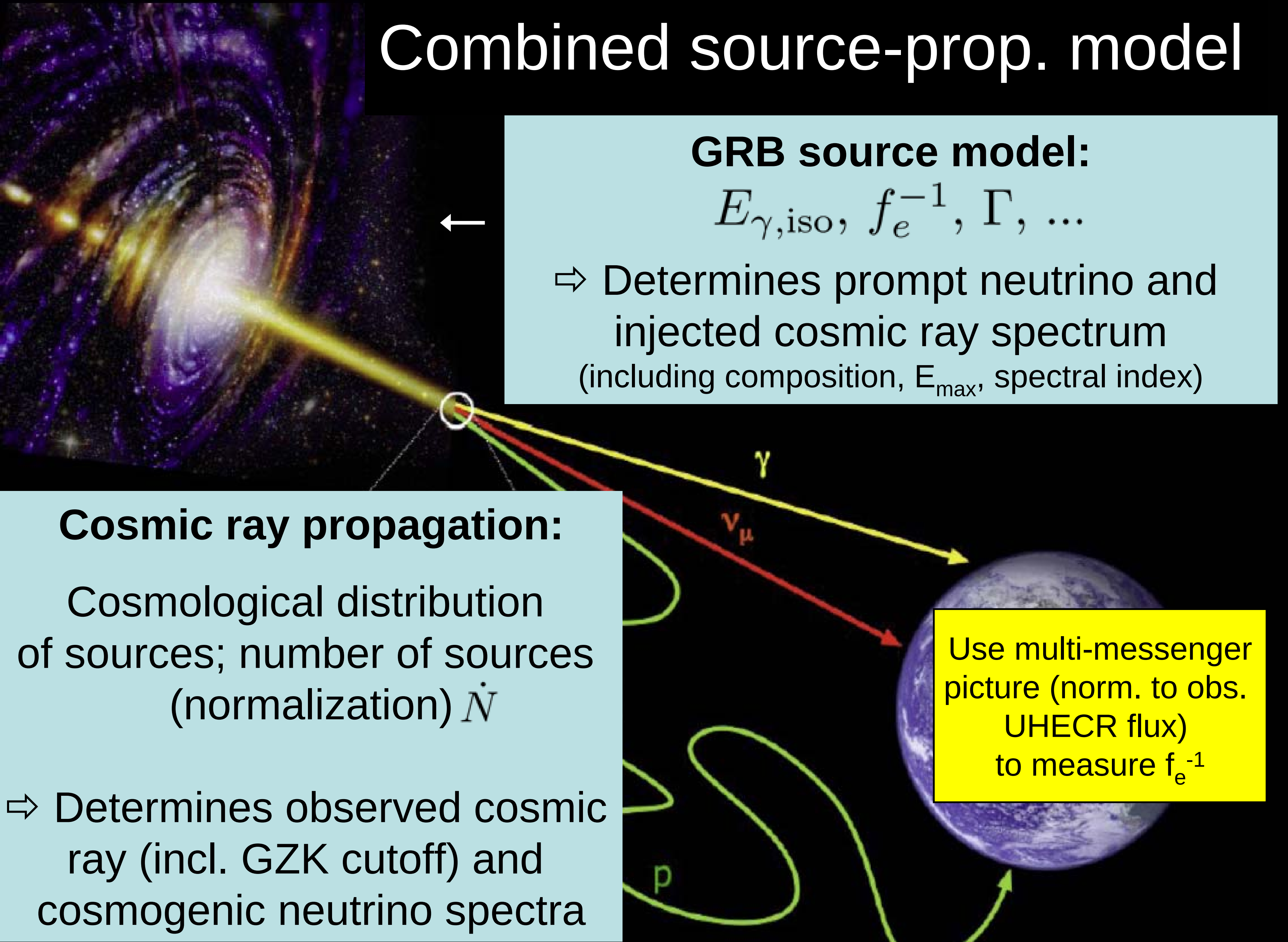


$\Gamma = 300,\ t_v = 0.01\,\mathrm{s},\ T_{90} = 10\,\mathrm{s},\ \eta = 1,$
$\epsilon_e/\epsilon_B = 1,\ f_e = 0.1,\ \alpha_\gamma = 1,\ \beta_\gamma = 2,\ \varepsilon'_{\gamma,b} = 1\,\mathrm{keV}$, and $z = 2$.

Combined source-propagation model

… as a first step towards completely self-consistent picture

Combined source-prop. model



GRB source model:

$$E_{\gamma,\mathrm{iso}},\ f_e^{-1},\ \Gamma,\ \ldots$$

⇨ Determines prompt neutrino and injected cosmic ray spectrum
(including composition, E_{max}, spectral index)

Cosmic ray propagation:

Cosmological distribution of sources; number of sources (normalization) $\dot{N}$

⇨ Determines observed cosmic ray (incl. GZK cutoff) and cosmogenic neutrino spectra

Use multi-messenger picture (norm. to obs. UHECR flux) to measure f_e^{-1}

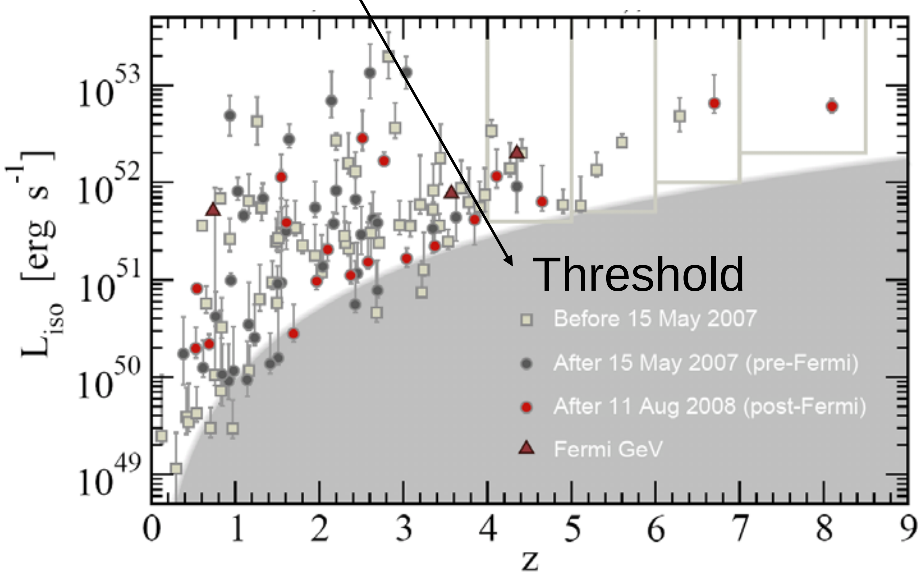
Gamma-ray observables?

- Redshift distribution

~ $(1+z)^{\alpha}$ → $\mathcal{E}(z)$

Star formation rate → $\dot{\rho}_*(z)$

$$\frac{d\dot{N}}{dz} = F(z)\frac{\mathcal{E}(z)\,\dot{\rho}_*(z)}{\langle f_{\rm beam}\rangle}\frac{dV/dz}{1+z}$$



(Kistler et al, Astrophys.J. 705 (2009) L104)

- Total number of observable bursts

$$\dot{N} = \int_0^{\infty} \frac{d\dot{N}}{dz}\,dz$$

  - Can be directly determined (counted)!
  - **Order 1000 yr^{-1}**

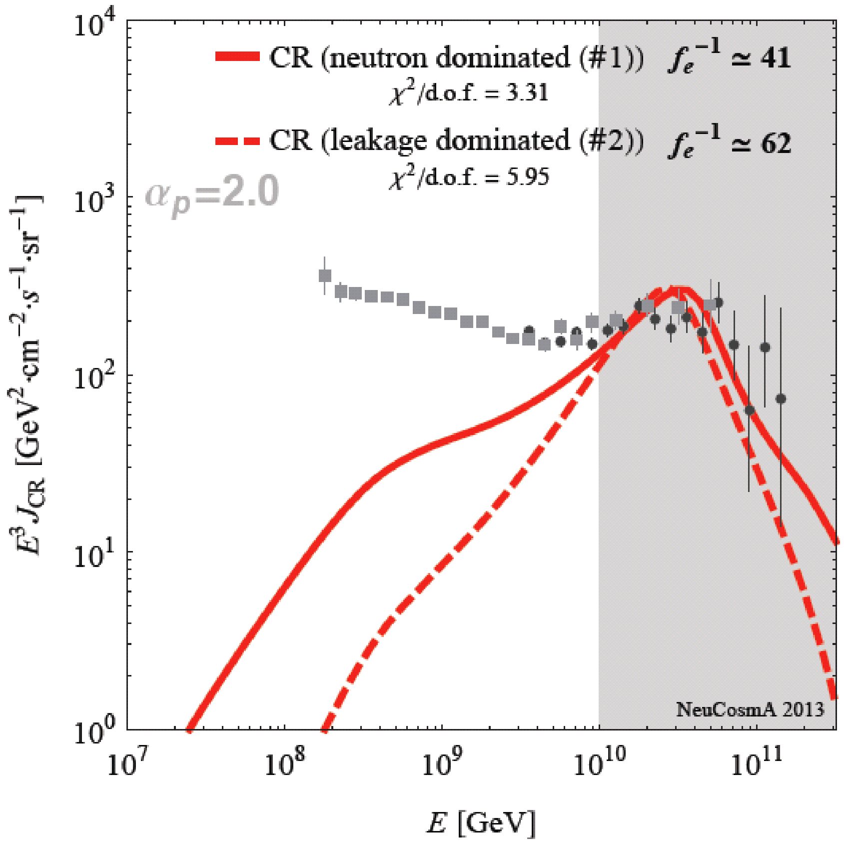
- Provides information about cosmic distribution of sources as well (α)

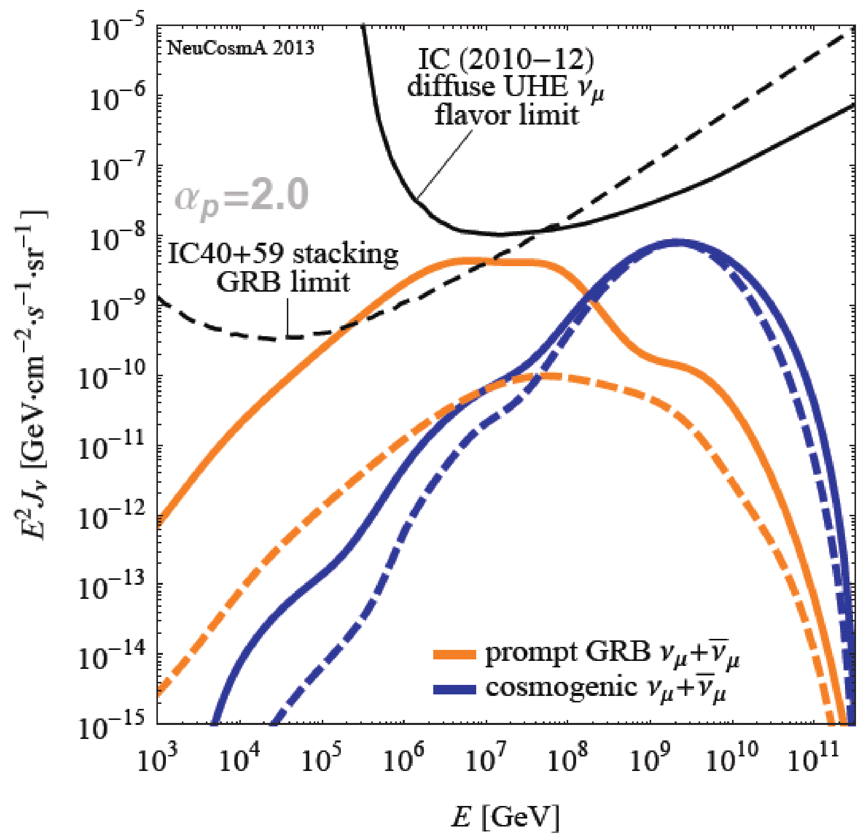
Combined source-prop. model fit

(cosmic ray ankle model transition, α_p ~ 2)

- Cosmic ray leakage (dashed) can evade prompt neutrino bound with comparable f_e^{-1}:

Fit ⇨ Consequence, incl. f_e^{-1}



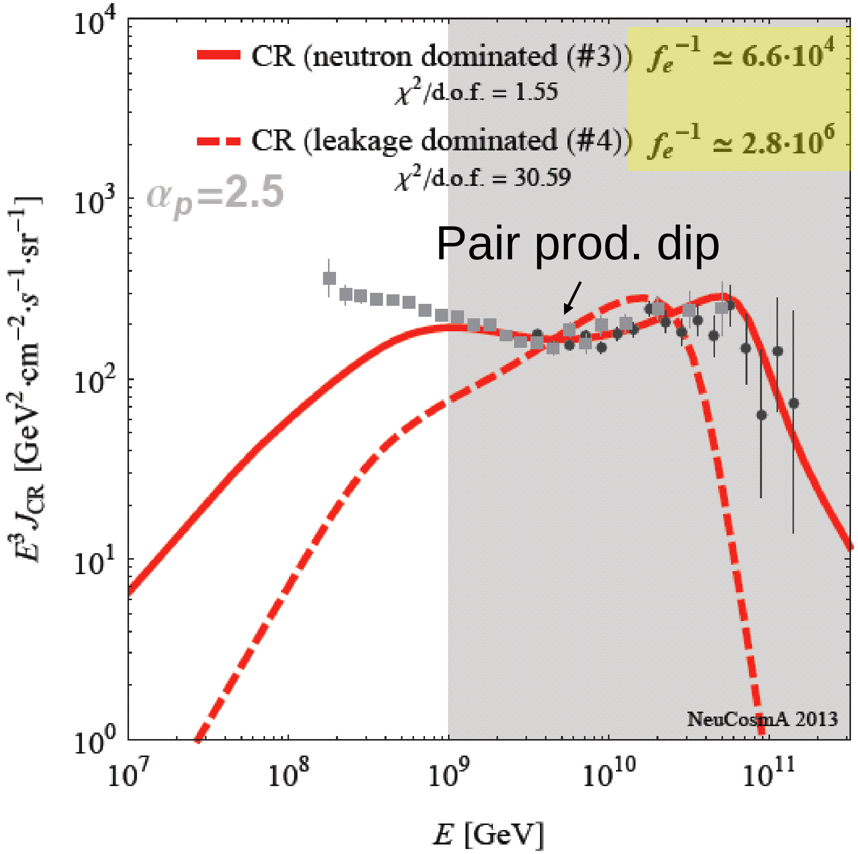


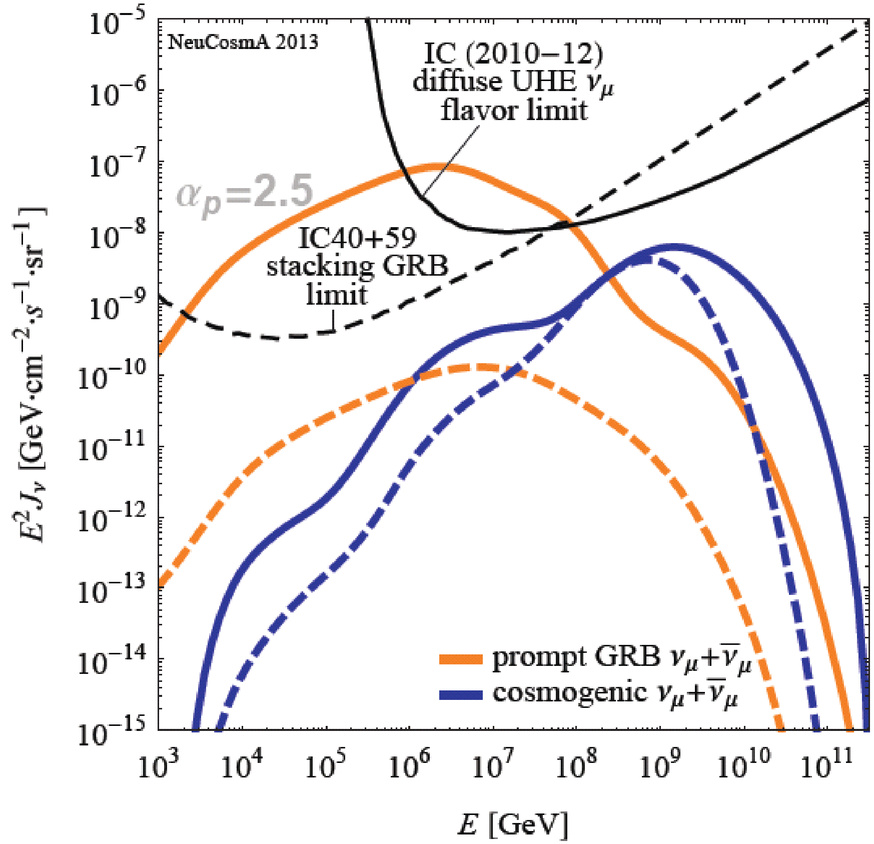
(Baerwald, Bustamante, Winter, arXiv:1401.XXXX)

Combined source-prop. model fit

(cosmic ray dip model transition, α_p ~ 2.5)

- Dip-model transition requires extremely large baryonic loadings (bol. correction!):

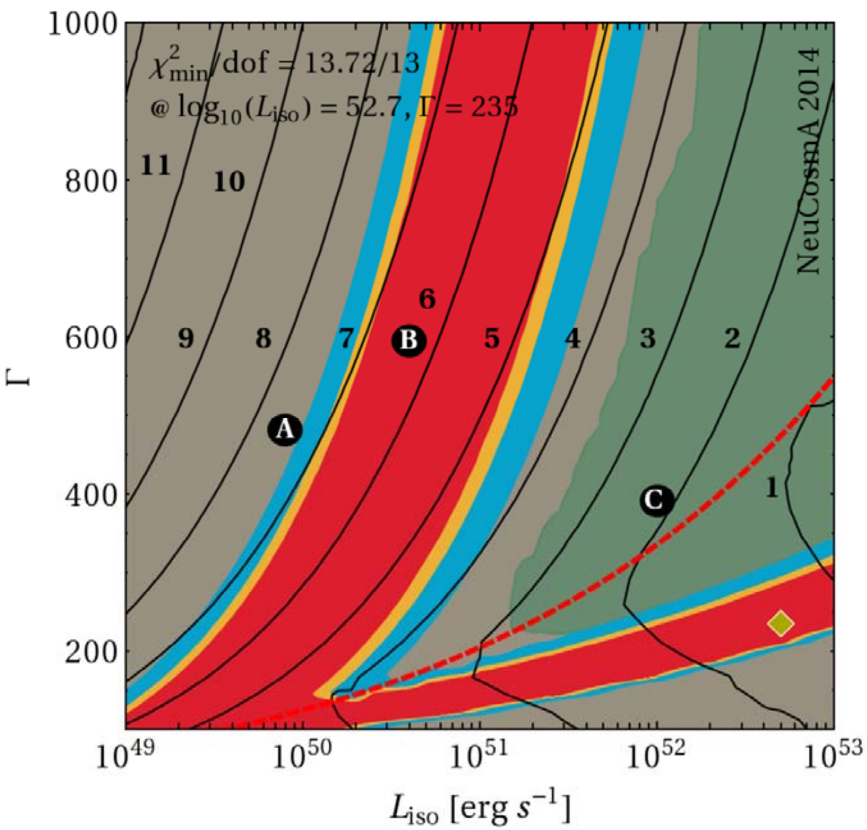




(Baerwald, Bustamante, Winter, , arXiv:1401.XXXX)

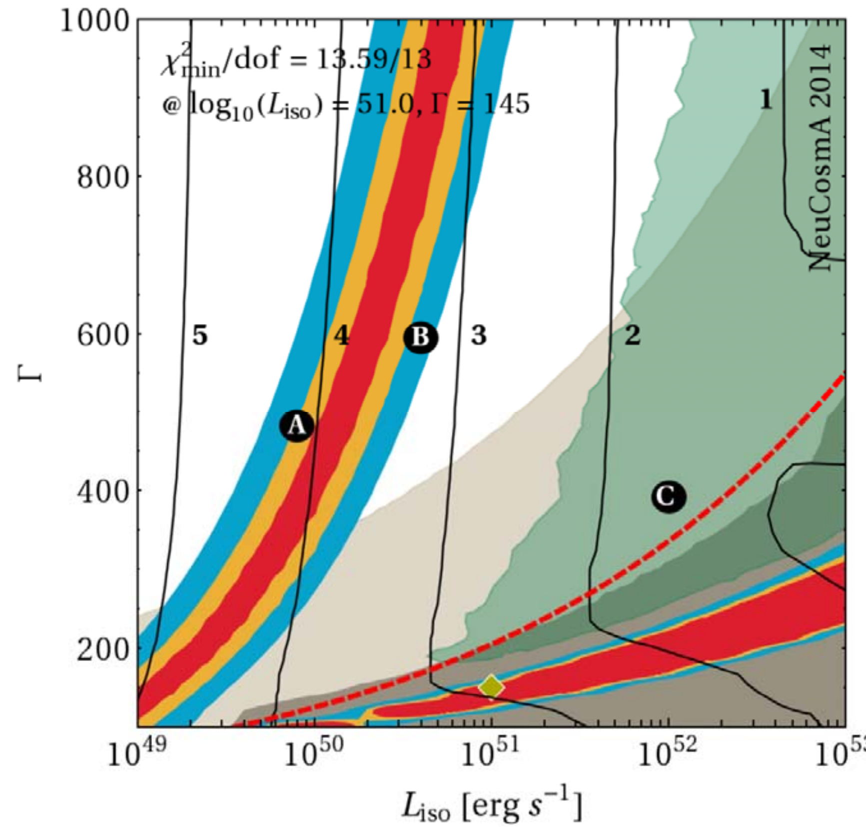
Parameter space?
[ankle model, high acceleration efficiency]

- Neutron model



➢ Ruled out by current stacking bounds on prompt neutrinos (dark gray)

- Diffusive escape

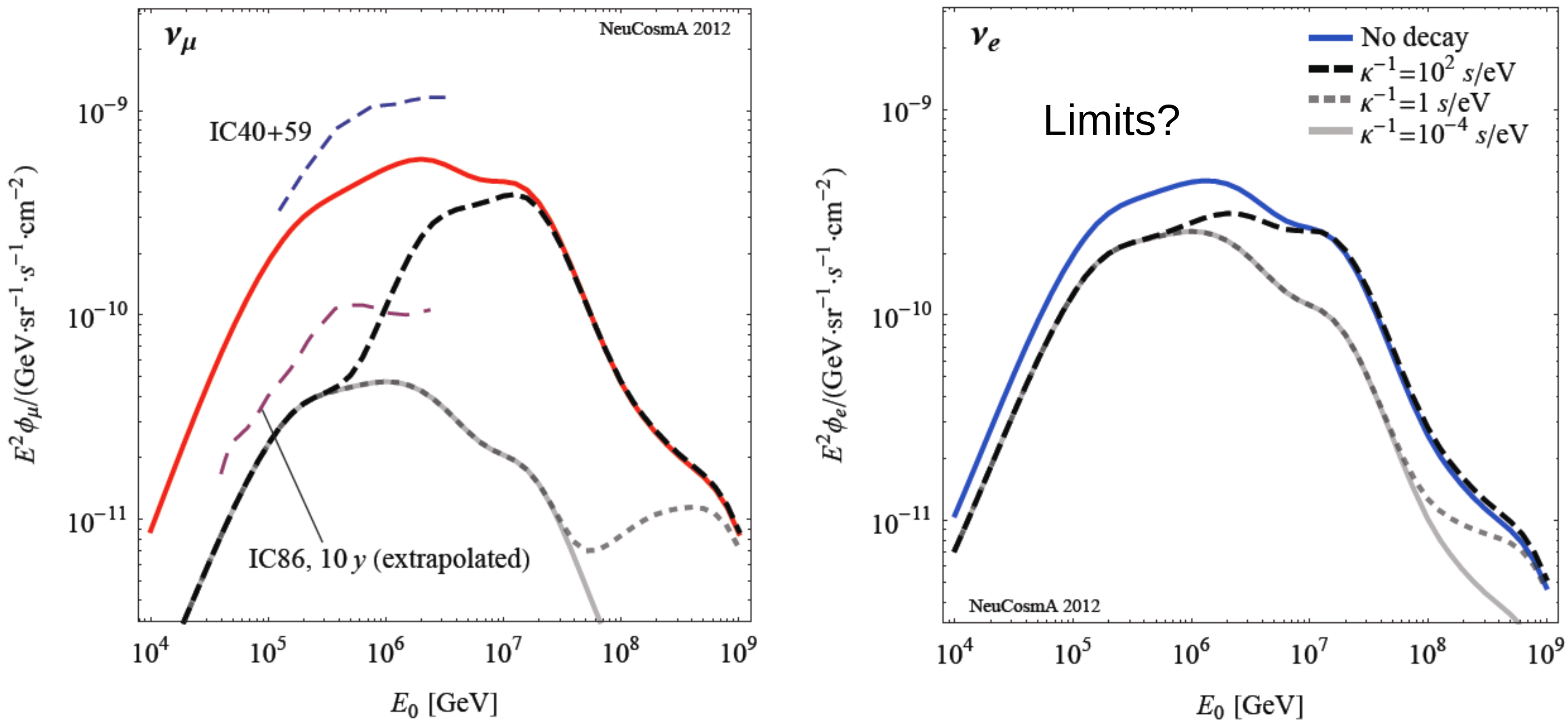


➢ Remaining parameter space requires extremely large (unrealistic?) baryonic loadings

(Baerwald, Bustamante, Winter, arXiv:1401.XXXX)

What if: Neutrinos decay?

Decay hypothesis: ν_2 and ν_3 decay with lifetimes compatible with SN 1987A bound



- Reliable conclusions from flux bounds require cascade (ν_e) measurements!

Baerwald, Bustamante, Winter, JCAP 10 (2012) 20

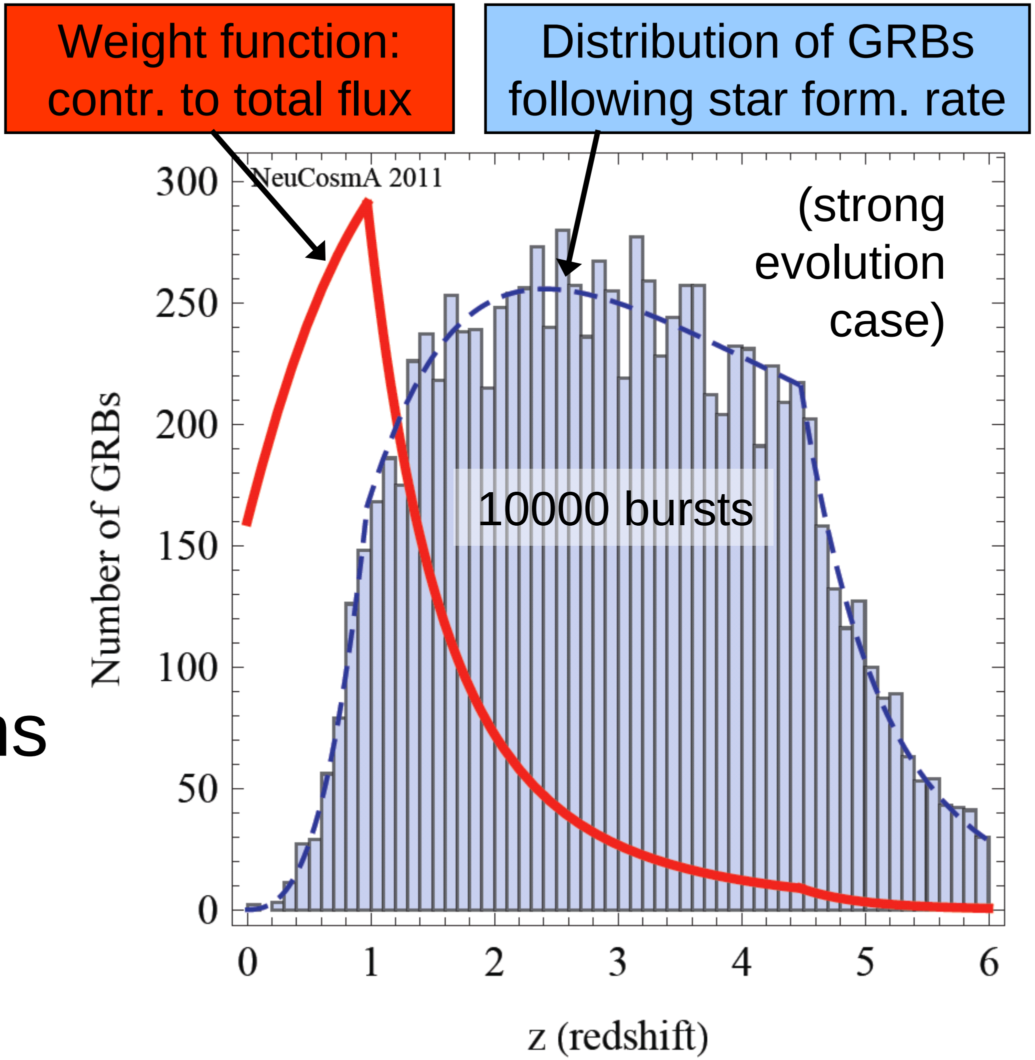
Summary

- GRB explanation of UHECR still possible (plausible?) in “ankle model“ for UHECR transition
- Neutron model for UHECR escape already excluded by current neutrino data
- Future neutrino bounds will strongly limit parameter space where pion production efficiency is large
- Possible ways out:
  - GRBs are *not* the exclusive sources of the UHECR
  - Extremely large baryonic loadings + direct/diffusive escape [applies not only to internal shock scenario …]
  - Model too simple? Do the cosmic rays and neutrinos come from different collision radii?

Backup

Systematics in aggregated fluxes

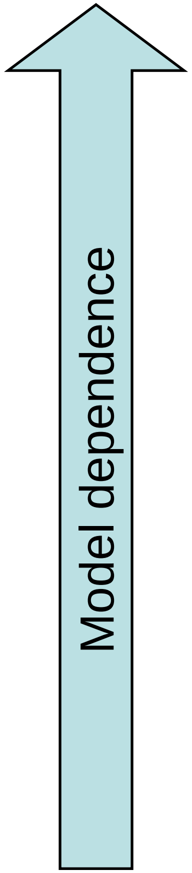
- z ~ 1 “typical“ redshift of a GRB
- Peak contribution in a region of low statistics
  - ➢Ensemble fluctuations of quasi-diffuse flux

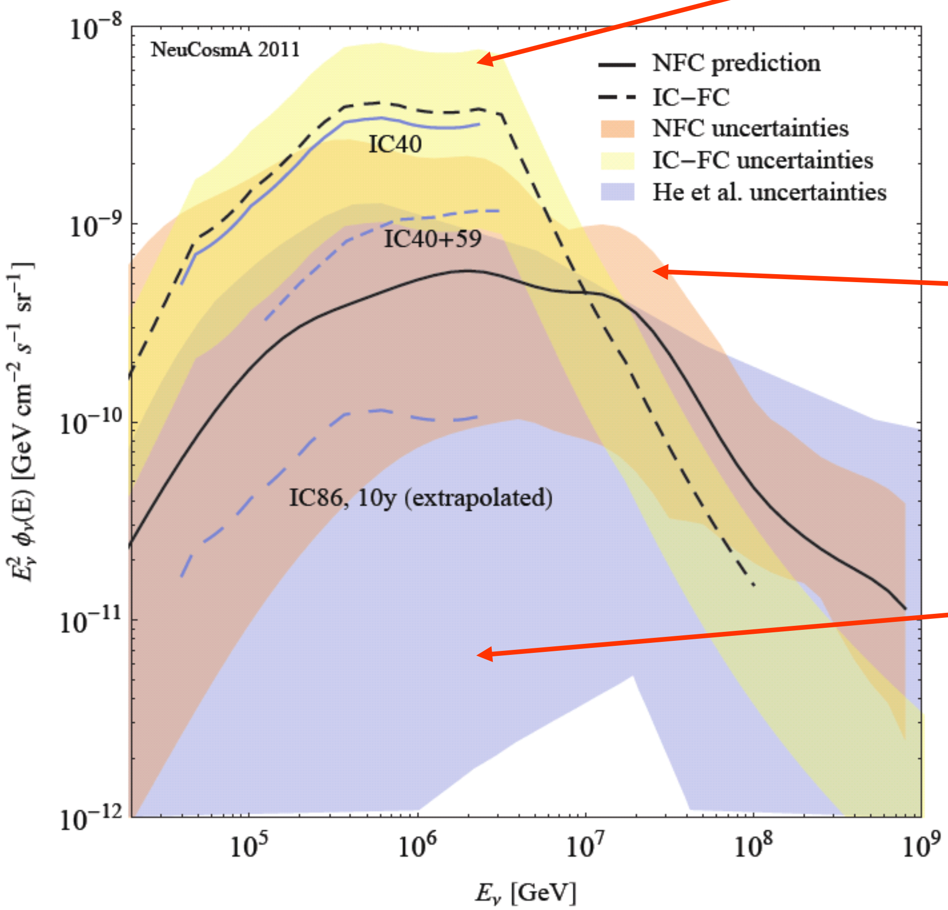


(Baerwald, Hümmer, Winter, Astropart. Phys. 35 (2012) 508)

Model dependence

Not only normalization, but also uncertainties depend on assumptions:





(figure courtesy of Philipp Baerwald)

Internal shock model, target photons from synchrotron emission/inverse Compton

from Fig. 3 of IceCube, Nature 484 (2012) 351; uncertainties from Guetta, Spada, Waxman, Astrophys. J. 559 (2001) 2001

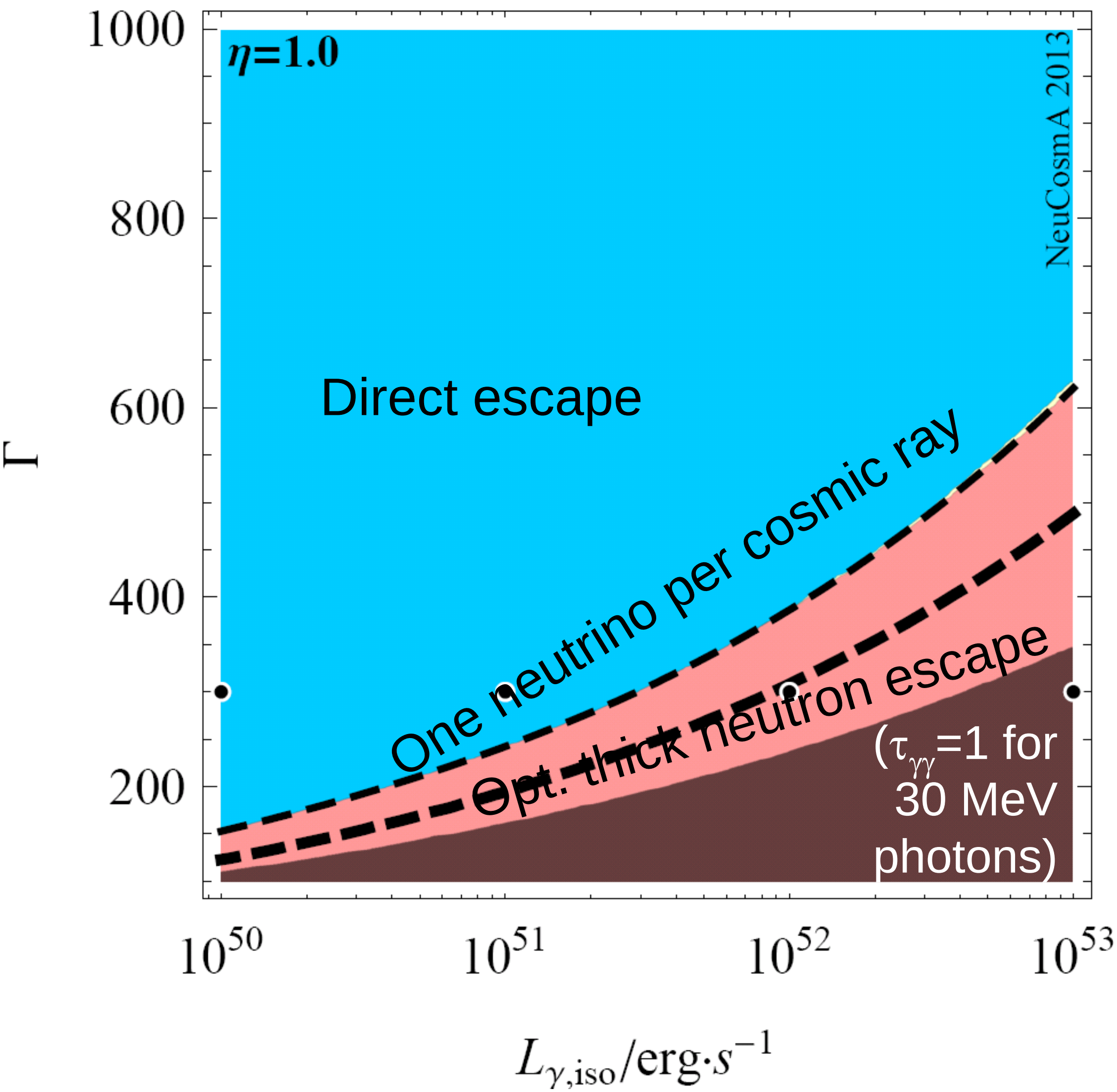
Internal shock model, target photons from observation, origin not specified

from Fig. 3 of Hümmer et al, PRL 108 (2012) 231101

Dissipation radius not specified (e. g. magnetic reconnection models), target photons from observation, origin not specified

from Fig. 3 of He, Murase, Nagataki, et al, ApJ. 752 (2012) 29

UHECR escape:
Parameter space?



- The challenge: need high enough E_p to describe observed UHECR spectrum
  - The acceleration efficiency η has to be high
  - Can evade the "one neutrino per cosmic ray" paradigm

(Baerwald, Bustamante, Winter, Astrophys. J. 768 (2013) 186)

Local GRB rate

- The local GRB rate can be written as

$$\dot{n}_{\mathrm{GRB}} = \frac{1}{\mathrm{Gpc}^3\,\mathrm{yr}} \, \frac{\dot{N}_{\mathrm{tot}}\,[\mathrm{yr}^{-1}]}{968} \, \frac{1}{f_z}$$

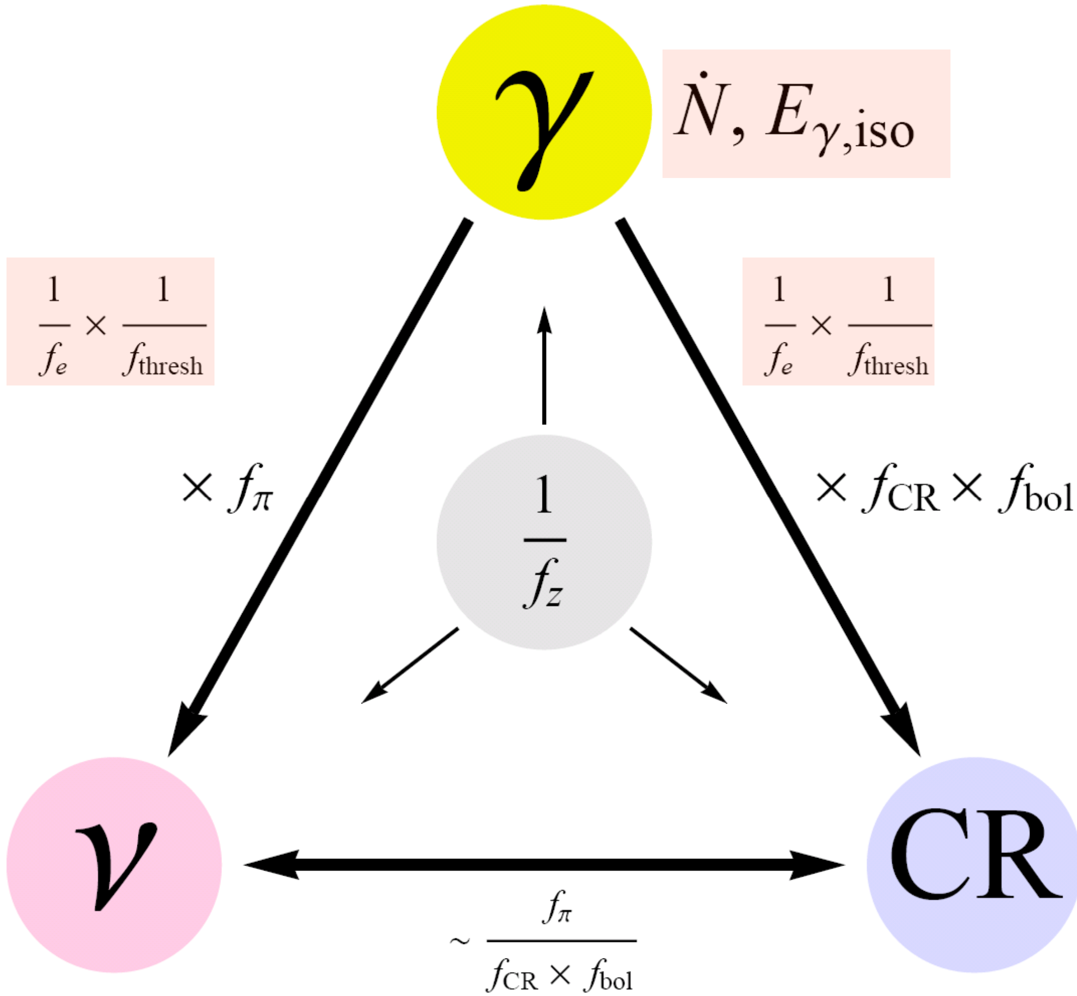
where f_z is a cosmological correction factor:

SFR model	α	f_z	$\dot{\tilde{n}}_{\mathrm{GRB}}\big\vert_{z=0}$ $[\mathrm{Gpc}^{-3}\,\mathrm{yr}^{-1}]$
Hopkins & Beacom (2006)	1.2	25.15	0.08
	0.0	5.65	0.35
Wanderman & Piran (2010)	0.0	7.70	0.26
Madau & Porciani (2000)			
SF1	0.0	9.89	0.21
SF2	0.0	14.42	0.14
SF3	0.0	14.36	0.14

(for 1000 observable GRBs per year)

(Baerwald, Bustamante, Halzen, Winter, to appear)

Impact factors



Same.
Focus on $1/f_e$ in following

Baryonic loading

Depends on
model for UHECR escape

(Baerwald, Bustamante, Winter, to appear)

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Required UHECR injection

- Required energy ejected in UHECR per burst:

$$E_{\mathrm{CR}}^{[10^{10},10^{12}]} = 10^{53}\,\mathrm{erg} \cdot \frac{\dot{\varepsilon}_{\mathrm{CR}}^{[10^{10},10^{12}]}}{10^{44}\,\mathrm{erg\,Mpc^{-3}\,yr^{-1}}} \cdot \frac{968\,\mathrm{yr}^{-1}}{\dot{N}_{\mathrm{tot}}} \cdot f_z$$

~1.5 to fit UHECR observations

~5-25

- In terms of γ-ray energy:

Fraction of energy in CR production?

How much energy in UHECR?

$$E_{\mathrm{CR}}^{[10^{10},10^{12}]} = f_{\mathrm{CR}} \frac{f_{\mathrm{bol}}}{f_e} E_{\gamma,\mathrm{iso}}$$

Energy in protons vs. electrons (IceCube def.)

➢ Baryonic loading f_e^{-1}~50-100 for E^{-2} inj. spectrum (f_{bol} ~ 0.2), $E_{\gamma,iso}$ ~ 10^{53} erg, neutron model (f_{CR} ~ 0.4)
[IceCube standard assumption: f_e^{-1}~10]

IceCube method ...normalization

- Connection γ-rays – neutrinos

½ (charged pions) x
¼ (energy per lepton)

Energy in protons

$$\int_0^\infty \mathrm{d}E_\nu\ E_\nu F_\nu(E_\nu) = \frac{1}{8}\ \frac{1}{f_e}\ \left(1 - (1 - \langle x_{p\to\pi}\rangle)^{\Delta R/\lambda_{p\gamma}}\right) \int_{1\,\mathrm{keV}}^{10\,\mathrm{MeV}} \mathrm{d}E_\gamma\ E_\gamma F_\gamma(E_\gamma)$$

Energy in neutrinos

Fraction of p energy converted into pions f_π

Energy in electrons/ photons

- Optical thickness to pγ interactions:

$$\frac{\Delta R}{\lambda_{p\gamma}} = \left(\frac{L_\gamma^{\mathrm{iso}}}{10^{52}\,\mathrm{erg\,s^{-1}}}\right) \left(\frac{0.01\,\mathrm{s}}{t_{\mathrm{var}}}\right) \left(\frac{10^{2.5}}{\Gamma_{\mathrm{jet}}}\right)^4 \left(\frac{\mathrm{MeV}}{\epsilon_\gamma}\right)$$

[in principle, $\lambda_{p\gamma} \sim 1/(n_\gamma\ \sigma)$; need estimates for n_γ, which contains the size of the acceleration region]

(Description in arXiv:0907.2227; see also Guetta et al, astro-ph/0302524; Waxman, Bahcall, astro-ph/9701231)

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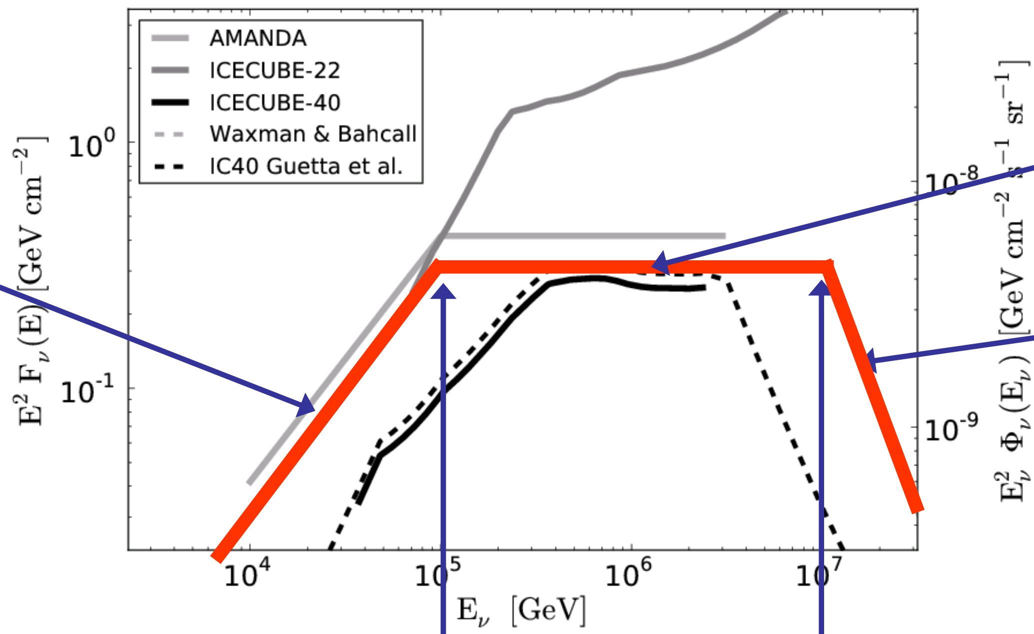
IceCube method ... spectral shape

- Example:

$$F_\gamma(E_\gamma) = \frac{\mathrm{d}N(E_\gamma)}{\mathrm{d}E_\gamma} = f_\gamma \times \begin{cases} \left(\frac{\epsilon_\gamma}{\mathrm{MeV}}\right)^{\alpha_\gamma} \left(\frac{E_\gamma}{\mathrm{MeV}}\right)^{-\alpha_\gamma} & \text{for } E_\gamma < \epsilon_\gamma \\ \left(\frac{\epsilon_\gamma}{\mathrm{MeV}}\right)^{\beta_\gamma} \left(\frac{E_\gamma}{\mathrm{MeV}}\right)^{-\beta_\gamma} & \text{for } E_\gamma \geq \epsilon_\gamma \end{cases}$$

$3\text{-}\beta_\gamma$



$3\text{-}\alpha_\gamma$

$3\text{-}\alpha_\gamma+2$

First break from break in photon spectrum (here: $E^{-1} \Rightarrow E^{-2}$ in photons)

Second break from pion cooling (simplified)