# Residual symmetries in the lepton mass matrices 

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## Contents

(1) Introduction
(2) Residual symmetries
(3) Group scan using GAP
(9) Residual symmetries and roots of unity
(6) $\mathrm{TM}_{1}$ and roots of unity
(0) Residual symmetries and caveats

## Introduction

## Tri-bimaximal mixing:

$$
\begin{gathered}
U_{\mathrm{TBM}}=\left(\begin{array}{rrr}
2 / \sqrt{6} & 1 / \sqrt{3} & 0 \\
-1 / \sqrt{6} & 1 / \sqrt{3} & -1 / \sqrt{2} \\
-1 / \sqrt{6} & 1 / \sqrt{3} & 1 / \sqrt{2}
\end{array}\right) \quad \text { ruled out! } \\
\sin ^{2} \theta_{13}=0.0227_{-0.0024}^{+0.0023} \\
\sin ^{2} \theta_{12}=0.302{ }_{-0.012}^{+0.013} \\
U=\left(\begin{array}{ccc}
-s_{12} c_{23}-c_{12} s_{23} s_{13} e^{i \delta} & c_{12} c_{23}-s_{12} s_{23} s_{13} e^{i \delta} & s_{13} e^{-i \delta} \\
c_{12} c_{23} c_{13} \\
s_{12} s_{23}-c_{12} c_{23} s_{13} e^{i \delta} & -c_{12} s_{23}-s_{12} c_{23} s_{13} e^{i \delta} & c_{23} c_{13}
\end{array}\right)
\end{gathered}
$$

## Introduction

$U=\left(U_{\alpha j}\right)=\left(u_{1}, u_{2}, u_{3}\right)$ with columns $u_{j}$
Albright, Rodejohann (2008): $\mathrm{TM}_{1}, \mathrm{TM}_{2}$ still valid!

$$
\mathrm{TM}_{1}: \quad u_{1}=\frac{1}{\sqrt{6}}\left(\begin{array}{c}
2 \\
-1 \\
-1
\end{array}\right), \quad \mathrm{TM} 2: \quad u_{2}=\frac{1}{\sqrt{3}}\left(\begin{array}{l}
1 \\
1 \\
1
\end{array}\right)
$$

$\mathrm{TM}_{1}: \quad s_{12}^{2}=1-\frac{2}{3 c_{13}^{2}}<\frac{1}{3}, \quad \cos \delta \tan 2 \theta_{23} \simeq-\frac{1}{2 \sqrt{2} s_{13}}\left(1-\frac{7}{2} s_{13}^{2}\right)$
$\mathrm{TM}_{2}: \quad s_{12}^{2}=\frac{1}{3 c_{13}^{2}}>\frac{1}{3}, \quad \cos \delta \tan 2 \theta_{23} \simeq \frac{1}{\sqrt{2} s_{13}}\left(1-\frac{5}{4} s_{13}^{2}\right)$

## Residual symmetries

## Fixing the notation:

Mass terms: Majorana neutrinos

$$
\mathcal{L}_{\text {mass }}=-\bar{\ell}_{L} M_{\ell} \ell_{R}+\frac{1}{2} \nu_{L}^{T} C^{-1} \mathcal{M}_{\nu} \nu_{L}+\text { H.c. }
$$

Diagonalization:
$U_{\ell}^{\dagger} M_{\ell} M_{\ell}^{\dagger} U_{\ell}=\operatorname{diag}\left(m_{e}^{2}, m_{\mu}^{2}, m_{\tau}^{2}\right), \quad U_{\nu}^{T} \mathcal{M}_{\nu} U_{\nu}=\operatorname{diag}\left(m_{1}, m_{2}, m_{3}\right)$
Mixing matrix: $U=U_{\ell}^{\dagger} U_{\nu}$

$$
\begin{aligned}
& V_{\ell}(\alpha) \equiv U_{\ell} \operatorname{diag}\left(e^{i \alpha_{1}}, e^{i \alpha_{2}}, e^{i \alpha_{3}}\right) U_{\ell}^{\dagger} \\
& V_{\nu}(\epsilon) \equiv U_{\nu} \operatorname{diag}\left(\epsilon_{1}, \epsilon_{2}, \epsilon_{3}\right) U_{\nu}^{\dagger} \quad \text { with } \quad \epsilon_{j}^{2}=1
\end{aligned}
$$

## Residual symmetries

Invariance of the mass matrices:

$$
V_{\ell}(\alpha)^{\dagger} M_{\ell} M_{\ell}^{\dagger} V_{\ell}(\alpha)=M_{\ell} M_{\ell}^{\dagger}, \quad V_{\nu}(\epsilon)^{T} \mathcal{M}_{\nu} V_{\nu}(\epsilon)=\mathcal{M}_{\nu}
$$

## Remarks:

- $V_{\ell}(\alpha) \in U(1) \times U(1) \times U(1), V_{\nu}(\epsilon) \in \mathbb{Z}_{2} \times \mathbb{Z}_{2} \times \mathbb{Z}_{2}$
- $V_{\ell}(\alpha), V_{\nu}(\epsilon)$ depend on VEVs and Yukawa coupling constants
- Invariance of mass matrices $V_{\ell}, V_{\nu}$ contains no information beyond diagonalizability


## Residual symmetries

## Idea of residual symmetries: C.W. Lam

- Weak basis $\Rightarrow \ell_{L}, \nu_{L}$ in same multiplet of $G$
- $G$ broken to different subgroups in charged-lepton and neutrino sectors:

$$
G_{\ell} \subseteq U(1) \times U(1) \times U(1), \quad G_{\nu} \subseteq \mathbb{Z}_{2} \times \mathbb{Z}_{2} \times \mathbb{Z}_{2}
$$

- For simplicity:

One generator $T$ of $G_{\ell}$, one generator $S$ of $G_{\nu}$ :

$$
T^{\dagger} M_{\ell} M_{\ell}^{\dagger} T=M_{\ell} M_{\ell}^{\dagger}, \quad S^{T} \mathcal{M}_{\nu} S=\mathcal{M}_{\nu}
$$

- For simplicity:
$T$ has three different eigenvalues
- Then $T$ and $S$ determine one column of $U$ !


## Residual symmetries

Why is this so?
(1) $S^{2}=\mathbb{1} \Rightarrow S= \pm\left(2 u u^{\dagger}-\mathbb{1}\right)$ with $S u= \pm u$
(2) $U_{\ell}^{\dagger} T U_{\ell}=\tilde{T}$ diagonal
(3) $U_{\ell}^{\dagger} u$ column in mixing matrix
(9) Two matrices $S_{1}, S_{2}$ with $S_{j}^{T} \mathcal{M}_{\nu} S_{j}=\mathcal{M}_{\nu} \Rightarrow$ mixing matrix $U$ completely determined

## Theorem

$$
\text { If } S^{T} \mathcal{M}_{\nu} S=\mathcal{M}_{\nu} \text { with } S= \pm\left(2 u u^{\dagger}-\mathbb{1}\right) \text {, then } \mathcal{M}_{\nu} u \propto u^{*}
$$

Remark: $U_{\ell}^{\dagger} u$ determined by the group! It does not contain parameters of the model.

## Residual symmetries

Two ways to tackle residual symmetries for the purpose of determination of possible flavour symmetry groups:
(1) Scanning finite groups
(2) Solving relations involving roots of unity

## Group scan using GAP

Holthausen, Lim, Lindner (2013):
$G_{\nu}=\mathbb{Z}_{2} \times \mathbb{Z}_{2}$, group results within $3 \sigma$ of fitted $s_{i j}^{2}$
a) Assumptions: ord $G<1536=3 \times 2^{9}$ (with one exception), $G_{\ell}$ generated by $\widetilde{T}=\operatorname{diag}\left(1, \omega, \omega^{2}\right)$ with $\omega=e^{2 \pi i / 3}$

| $n$ | $G$ | $s_{12}^{2}$ | $s_{13}^{2}$ | $s_{23}^{2}$ |
| :---: | :---: | :---: | :---: | :---: |
| 5 | $\Delta\left(6 \times 10^{2}\right)$ | 0.3432 | 0.0288 | 0.3791 |
|  |  | 0.3432 | 0.0288 | 0.6209 |
| 9 | $\left(\mathbb{Z}_{18} \times \mathbb{Z}_{6}\right) \rtimes S_{3}$ | 0.3402 | 0.0201 | 0.3992 |
| 16 | $\Delta\left(6 \times 16^{2}\right)$ | 0.3402 | 0.0201 | 0.6008 |
|  |  | 0.3420 | 0.0254 | 0.3867 |
|  |  | 0.3420 | 0.0254 | 0.6133 |

b) Assumptions: ord $G<512, G_{\ell}$ Abelian $\Rightarrow$ no candidates!

## Group scan using GAP

Model-building addendum to group scan (Grimus, Lavoura (2013)):

\[

\]

## Residual symmetries and roots of unity

Basic assumption: Flavour group $G$ finite! (finitely generated) Mixing matrix: $U=\left(U_{\alpha j}\right)(\alpha=e, \mu, \tau, j=1,2,3)$

- $G_{\ell}$ generated by $T, G_{\nu}$ generated by $S$
- $\operatorname{det} S=1 \Rightarrow S=2 u u^{\dagger}-\mathbb{1}$
- Finiteness $\Rightarrow \exists m, n \in \mathbb{N}$ such that $T^{m}=S^{2}=(S T)^{n}=\mathbb{1}$ $T$ has eigenvalues $e^{i \phi_{\alpha}}$, $S T$ has eigenvalues $\lambda_{j} \Rightarrow$ $\operatorname{Tr}(S T)=\lambda_{1}+\lambda_{2}+\lambda_{3}$


## Trace and determinant of ST

Hernandez, Smirnov (2012)
$u$ i-th column of $U \Rightarrow$ two equations for 6 roots of unity:

$$
\sum_{\alpha=e, \mu, \tau}\left(2\left|U_{\alpha i}\right|^{2}-1\right) e^{i \phi_{\alpha}}=\lambda_{1}+\lambda_{2}+\lambda_{3} \quad \text { and } \quad \prod_{\alpha} e^{i \phi_{\alpha}}=\lambda_{1} \lambda_{2} \lambda_{3}
$$

$\mathrm{TM}_{1}$ and roots of unity

Which finite group can enforce $\mathrm{TM}_{1}$ ? Grimus (2013)

$$
\mathrm{TM}_{1}: \quad u_{1}=\frac{1}{\sqrt{6}}\left(\begin{array}{c}
2 \\
-1 \\
-1
\end{array}\right) \Rightarrow \begin{aligned}
& 2\left|U_{e 1}\right|^{2}-1=\frac{1}{3} \\
& 2\left|U_{\mu 1}\right|^{2}-1=-\frac{2}{3} \\
& 2\left|U_{\tau 1}\right|^{2}-1=-\frac{2}{3}
\end{aligned}
$$

Vanishing sum of roots of unity:

$$
-e^{i \phi_{e}}+2 e^{i \phi_{\mu}}+2 e^{i \phi_{\tau}}+3 \lambda_{1}+3 \lambda_{2}+3 \lambda_{2}=0
$$

Solution by theorem of Conway and Jones (1976)

## $\mathrm{TM}_{1}$ and roots of unity

Formal sums of roots of unity: ring over rational numbers $\omega=e^{2 \pi i / 3}, \beta=e^{2 \pi i / 5}, \gamma=e^{2 \pi i / 7}$

## Theorem (Conway and Jones (1976))

Let $\mathcal{S}$ be a non-empty vanishing sum of length at most 9 .
Then either $\mathcal{S}$ involves $\theta, \theta \omega, \theta \omega^{2}$ for some root $\theta$, or $\mathcal{S}$ is similar to one of

$$
\begin{aligned}
& 1+\beta+\beta^{2}+\beta^{3}+\beta^{4}, \\
& -\omega-\omega^{2}+\beta+\beta^{2}+\beta^{3}+\beta^{4}, \\
& 1+\beta+\beta^{2}-\left(\omega+\omega^{2}\right)\left(\beta^{2}+\beta^{3}\right), \\
& 1+\gamma+\gamma^{2}+\gamma^{3}+\gamma^{4}+\gamma^{5}+\gamma^{6}, \\
& -\omega-\omega^{2}+\gamma+\gamma^{2}+\gamma^{3}+\gamma^{4}+\gamma^{5}+\gamma^{6}, \\
& \beta+\beta^{4}-\left(\omega+\omega^{2}\right)\left(1+\beta^{2}+\beta^{3}\right), \\
& 1+\gamma^{2}+\gamma^{3}+\gamma^{4}+\gamma^{5}-\left(\omega+\omega^{2}\right)\left(\gamma+\gamma^{6}\right), \\
& 1-\left(\omega+\omega^{2}\right)\left(\beta+\beta^{2}+\beta^{3}+\beta^{4}\right) .
\end{aligned}
$$

## $\mathrm{TM}_{1}$ and roots of unity

## Solution:

$e^{i \phi_{e}}=\eta, e^{i \phi_{\mu}}=\eta \omega, e^{i \phi_{\tau}}=\eta \omega^{2}, \lambda_{1}=\epsilon, \lambda_{2}=-\epsilon, \lambda_{3}=\eta$
$\eta$ is an arbitrary root of unity, $\epsilon= \pm i \eta$

In basis where charged lepton mass matrix is diagonal:

$$
\begin{aligned}
\tilde{T} & =\eta \operatorname{diag}\left(1, \omega, \omega^{2}\right) \\
u_{1} & =\frac{1}{\sqrt{6}}\left(\begin{array}{c}
2 \\
-1 \\
-1
\end{array}\right) \Rightarrow \widetilde{S}=\frac{1}{3}\left(\begin{array}{rrr}
1 & -2 & -2 \\
-2 & -2 & 1 \\
-2 & 1 & -2
\end{array}\right)
\end{aligned}
$$

$\widetilde{T}$ and $\widetilde{S}$ generate group $\mathbb{Z}_{q} \times S_{4}$ with $\eta$ being a primitive root of order $q$.

Another basis:

$$
\begin{gathered}
U_{\omega}=\frac{1}{\sqrt{3}}\left(\begin{array}{ccc}
1 & 1 & 1 \\
1 & \omega & \omega^{2} \\
1 & \omega^{2} & \omega
\end{array}\right) \text { with } \omega=e^{2 \pi i / 3}=\frac{-1+i \sqrt{3}}{2} \\
S=U_{\omega} \tilde{S} U_{\omega}^{\dagger}=\left(\begin{array}{rrr}
-1 & 0 & 0 \\
0 & 0 & 1 \\
0 & 1 & 0
\end{array}\right), \quad T=U_{\omega} \tilde{T} U_{\omega}^{\dagger}=\eta\left(\begin{array}{lll}
0 & 1 & 0 \\
0 & 0 & 1 \\
1 & 0 & 0
\end{array}\right) \\
E^{\dagger}\left(M_{\ell} M_{\ell}^{\dagger}\right) E=M_{\ell} M_{\ell}^{\dagger} \Rightarrow U_{\omega}^{\dagger}\left(M_{\ell} M_{\ell}^{\dagger}\right) U_{\omega} \text { is diagonal }
\end{gathered}
$$

## $\mathrm{TM}_{1}$ and roots of unity

$$
S u=u \quad \Rightarrow \quad u=\frac{1}{\sqrt{2}}\left(\begin{array}{l}
0 \\
1 \\
1
\end{array}\right)
$$

## Mechanism for $\mathrm{TM}_{1}$ :

Lavoura, de Madeiros Varzielas (2012); Grimus (2013)
$U_{\omega}$ diagonalizes $M_{\ell} M_{\ell}^{\dagger}$ and $u$ eigenvector of $\mathcal{M}_{\nu} \Rightarrow$

$$
U_{\omega}^{\dagger} u=\frac{1}{\sqrt{6}}\left(\begin{array}{c}
2 \\
-1 \\
-1
\end{array}\right) \text { is column in mixing matrix }
$$

Example: $S_{4}$ and type II seesaw mechanism
Needs 7 scalar gauge doublets in $\mathbf{1} \oplus \mathbf{3} \oplus \mathbf{3}^{\prime}$
and 4 gauge triplets in $\mathbf{1} \oplus \mathbf{3}^{\prime}+\mathrm{VEV}$ alignment

$$
M_{\ell}=\left(\begin{array}{ccc}
a & b+c & b-c \\
b-c & a & b+c \\
b+c & b-c & a
\end{array}\right), \quad \mathcal{M}_{\nu}=\left(\begin{array}{ccc}
A & B & -B \\
B & A & C \\
-B & C & A
\end{array}\right)
$$

## Residual symmetries and caveats

Notation:
$G=$ flavour symmetry group of the Lagrangian
$\bar{G}=$ group determined by residual symmetries in $M_{\ell} M_{\ell}^{\dagger}$ and $\mathcal{M}_{\nu}$

- Restriction:
- Symmetry group $G$ of Lagrangian is finitely generated
- Neutrinos have Majorana nature
- Possible relationship between $G$ and $\bar{G}$ :
- $\bar{G} \subset U(3)$ due to 3 families
- Method is purely group-theoretical and uses only information contained in the mass matrices $\Rightarrow \bar{G}$ can at most yield $D(G)$
- Accidental symmetries in the mass matrices $\Rightarrow$ $\bar{G}$ not even a subgroup of $D(G)$
- Total breaking of $G$ :

Method not applicable, though model might be predictive

## Thank you for your attention!

