# New Physics with Low Missing Energy: Identification and Discrimination at the LHC

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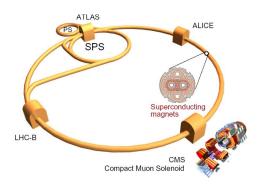
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- K. Ghosh, BM, S. Mukhopadhyaya, JHEP 1010,096 (2010)
  - S. Mukhopadhyay, BM, Phys. Rev. D82, 031501 (2010)
  - S. Mukhopadhyay, BM, arXiv:1108.4921 [hep-ph] (To appear in Phys. Rev. D)

# The LHC...



#### Goals of the LHC....

- To discover the Higgs boson and complete the Standard Model
- To know more about top and bottom quark properties
- To understand strong interaction better
- To look for quark-gluon plasma
- Physics beyond the standard electroweak theory

BUT....

• Why do we think there should be new physics ?

#### BUT....

- Why do we think there should be new physics?
- Why should new laws be manifest at the LHC energy?



# Phenomenological dissatisfactions (unexplained features):

- Large number of unrelated free parameters
- Replication of fermion families
- The pattern of fermion masses
- Maximal P but small CP violation

#### Theoretical/Philosophical questions:

- No way to unify with strong interaction
- No clue on a quantum theory of gravity
- Divergent higher-order contributions to the Higgs boson mass

## Sporadic/seasonal/volatile issues:

- The muon anomalous magnetic moment (3 3.5  $\sigma$  inconsistency)
- ullet PAMELA (excess positrons  $\sim$  10 80 GeV from galactic halo)
- $\bullet$  ATIC (excess galactic cosmic-ray electrons  $\sim$  300 -800 GeV)
- Tevatron multimuon events (excess multimuons, inexplicable from
  - b-decays)
- Top quark forward-backward asymmetry at Tevatron
- Wjj events in the CDF experiment



#### **Concrete and persistent problems:**

- Neutrino masses and mixing
- Cold dark matter (no particle physics explanation)
- Matter-antimatter asymmetry in the universe
- A positive cosmological constant (!)

# Physics beyond the standard model...

Effective energy scale to be probed at the LHC:  $\simeq 1$  -2 TeV

Out of the many motivations listed, which ones definitely suggest 'something new' at this energy?

# Why new physics at the TeV scale?

- The issue of Grand Unification (very indirect!)
- To understand why the Higgs should be within a TeV (relatively pressing!)
- Finding a cold(warm?) dark matter candidate (Quite imperative)

Thus the Dark matter issue is rather central to new physics search at the LHC

# Searches for new physics at the LHC...

Events with large missing- $E_T$  (MET) or resonances are the first to be looked for

Dark matter candidates produced ⇒ Events with MET

A candidate theory: supersymmetry (SUSY)
A lot of progress has taken places in SUSY search—
mSUGRA-based cMSSM ruled out upto about a TeV

Theories often proposed with a  $Z_2$  symmetry to accommodate a stable particle fitting in as dark matter candidate

One also studies how to distinguish among various such scenarios

(SUSY with R-parity, Universal extra dimensions with KK parity, Little Higgs with T-parity)



# Distinguishing among models with $Z_2$ ....

'Inverse problem' within SUSY— mapping from signature space to parameter space
Arkani-Hamed et al. (2006)
Larger number of observables studied

⇒ degeneracy in the parameter space better lifted

SUSY vs. other scenarios with large MET General distinction strategies:

- A.K. Datta, G. Kane, M. Toharia (2005)
- A. K Datta, P. Dey, S. Gupta, BM, A. Nyffeler (2007)
- J. Hubisz et al. (2008)
- W. Ehrenfeld et al. (2009)
- B. Bhattacharjee it et al. (2009)



## But is is also important to remember that....

So long as large MET signals elude us,

One needs to think of scenarios where such  $Z_2$  symmetry is broken

Some of the resulting theories may still accommodate dark matter candidates

We need criteria to point towards them and to discriminate among various low-MET scenarios in general

# Examples....

- (a) SUSY with R-parity violation (SUSY-RPV)
- (b) Little Higgs theories with broken T-parity (LHT-TPV)
- (c) Universal extra dimensions with conserved Kaluza-Klein parity (UED-KKC)
- (d) Universal extra dimensions with Kaluza-Klein parity violated (UED-KKV)
- (e) SUSY with a compressed spectrum

The present discussion includes cases (a) - (d)

#### The MSSM superpotential:

$$W_{MSSM} = Y_{ij}^{l} L_{i} \dot{H}_{1} E_{j}^{c} + Y_{ij}^{d} Q_{i} H_{1} D_{j}^{c} + Y_{ij}^{u} Q_{i} H_{1} U_{j}^{c}$$

When  $R = (-)^{L+3B+2J}$  violated via L or B, one can write

$$W=W_{MSSM}+W_{RPV}$$
, with  $W_{RPV}=\lambda_{ijk}L_{i}L_{j}E_{k}^{c}+\lambda_{ijk}^{'}L_{i}Q_{j}D_{k}^{c}+\epsilon_{i}L_{i}H_{2}+\lambda_{ijk}^{"}U_{i}^{c}D_{j}^{c}D_{k}^{c}$ 

We consider here L-violating  $W_{RPV}$ :  $\lambda_{ijk}^{"}=0$ 

A relevant mechanism for neutrino mass generation



## SUSY-RPV....

Most phenomenological studies: one type of R-parity violating coupling at a time

Result: the MSSM-LSP (say, the lightest neutralino) has two/three body decays with at least one lepton in the final state

The gravitino or the axino may be the dark matter candidate

## LHT-TPV....

The Higgs is the pesudo-goldstone boson of a broken approximate global symmetry In the minimal (littlest) form, Underlying electroweak gauge group :  $[SU(2) \times U(1)]^2$  , with an exchange symmetry (T-parity) – a  $Z_2$  symmetry Result: a division into T-even (SM) and T-odd (new) particles  $[SU(2) \times U(1)]^2 \longrightarrow SU(2) \times U(1)$  at scale f New particles include heavy T-odd fermions  $(Q_H, L_H)$ , heavy gauge bosons  $(W_H, Z_H, A_h)$ , a Higgs triplet..... The lightest T-odd particle (LTP) is stable: (Usually the  $A_H$ )

## LHT-TPV....

 $Z_2$  symmetry  $\Rightarrow$  LTP is the dark matter candidate (A neutral, weakly interacting particle) The spectrum and the interactions are controlled by  $f = \mathcal{O}(TeV)$ ,  $f \kappa_{ii} = \text{matrix deciding heavy fermion masses}$ But T-parity can be broken... by Wess-Zumino-Witten anomaly terms **Results:** terms  $\sim \epsilon_{\mu\nu\alpha\beta} V^{\mu}_{\mu} V^{\nu} \partial^{\alpha} V^{\beta}$ The LTP becomes unstable: For example,  $A_H \longrightarrow W^{(*)}W^{(*)}$ , leading to tree-level or loop-induced decays

#### UED-KKC....

At least one spacelike compact extra dimension, of radius R, where all fields can propagate

New particles are Kaluza-Klein towers in the 4D projections,

with same spin as in the zero-mode SM states

Two orbiforld fixed points in the extra dimension: a  $Z_2$  symmetry (for ensuring proper fermion chiralities)  $\Rightarrow$  A conserved Kaluza-Klein parity

The lightest KK-odd particle (LKP) is stable due to the  $Z_2$  symmetry: dark matter candidate (A neutral, weakly interacting particle): Usually the excited photon  $A_1$ 

The spectrum is decided by  $R^{-1}$ , and the cut-off scale A highly compressed spectrum in general



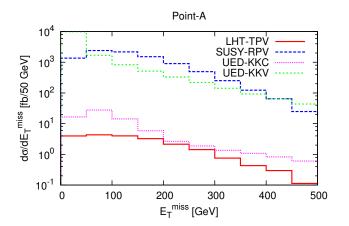
#### UED-KKV...

KK-parity can be broken by additional operators at the orbifold fixed points

⇒ The LKP is again unstable

Claim: all of these four scenarios, including UED-KKC, can lead to similar MET signals at the LHC How to distinguish?

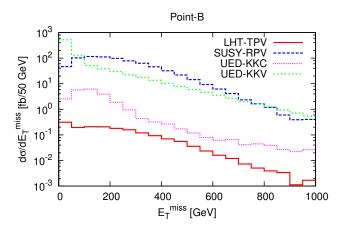
## MET distributions for all four scenarios...



 $M_s = Strongly interacting particle mass \simeq 600 GeV SUSY-RPV: <math>\lambda$ -type with one coupling



## MET distributions for all four scenarios...



 $M_s = Strongly interacting particle mass \simeq 1 TeV SUSY-RPV: <math>\lambda$ -type with one coupling



# Distinguishing among the various scenarios...

In addition to di-and tri-leptons, four-and five-lepton final states can be useful discriminators

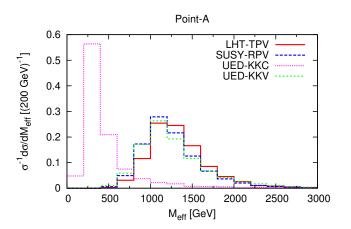
Isolated and central leptons

A  $p_T$ -cut of 10 GeV on the two softest leptons keeps away backgrounds

 $M_{l^+l^-} \geq 20 \, GeV$  for all opposite-sign lepton pairs keep away  $\gamma^*$  backgrounds

For  $M_s=600$  GeV, 5 fb<sup>-1</sup> at 14 TeV is enough For  $M_s=1$  TeV, 30 fb<sup>-1</sup> is required for  $5\sigma$  significance for all scenarios

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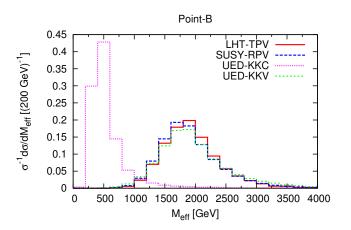


Effective mass distribution with  $M_s = 600$  GeV: UED-KKC stands out

$$(M_{eff} = \Sigma_i p_T^i + MET)$$



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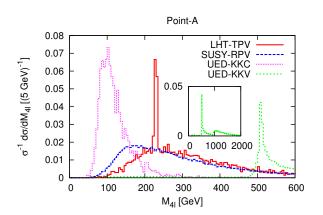


Effective mass distribution with  $M_s = 1$  TeV: UED-KKC stands out

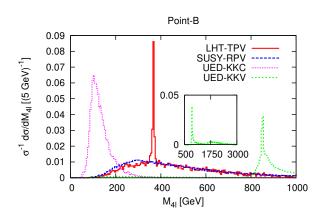
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# Four-lepton invariant mass...



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• Pairwise opposite-sign lepton invariant masses and their correlations in 4ℓ events

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- Pairwise opposite-sign lepton invariant masses and their correlations in 4ℓ events
- Angular correlation of each lepton with the nearest jet (for SUSY-RPV, there is often peaking in the forward direction)
- $N(5\ell)/N(4\ell)$ : SUSY-RPV and LHT-TPV show higher ratios than the two other cases

# Same-sign trileptons(SS3 $\ell$ ): unexplored potential...

Lepton sign: seriously used in the search for same-sign dilepton (SSD) events

Majorana fermions enhance SSD rates,  $p_T$  + isolation cuts reduce backgrounds (mostly from  $t\bar{t}$ )

Leptons of higher multiplicity: SM backgrounds extremely small

Theories with L-violation + self-conjugate fields: unsuppressed signals

A very discriminating check on scenarios with low-MET

SUSY-RPV stands out by contributing to SS3 $\ell$  (and also SS4 $\ell$ ) (Even in the early run)

Also, the dynamics of R-parity violation can be probed thereby S. Mukhopadhyay, BM (2010, 2011)

# Same-sign trileptons(SS3 $\overline{\ell}$ ): unexplored potential...

Standard model contribution to  $\sigma(\text{SS}3\ell)$ : with appropriate cuts,  $\simeq 2.5 \times 10^{-3}$  fb (  $\simeq 7.0 \times 10^{-4}$  fb ) at 14 (7) TeV Even smaller backgrounds for SS4 $\ell$ 

If high-MET new physics signals continue to elude us, Low-MET ones must be looked for  $SS3\ell \Rightarrow$  a discriminating signature of specific scenario(s) In SUSY-RPV, LSP-pair decays (with no branching ratio suppression) can yield two same-sign leptons, and one more comes from the cascade

• SUSY with R-parity: rates are always very low due to lack of L-violation and branching fraction suppression

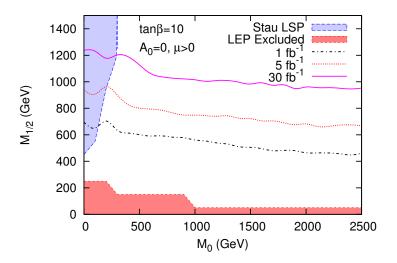
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- LHT-TPV: Decay of the LKP to one charged lepton has low branching ratio
- UED: SS3ℓ occurs even more rarely
- Thus SUSY-RPV stands out

# Coverage in mSUGRA parameter space at 14 TeV



5-event contours with background << 1: One  $\lambda$ -type coupling included

### Rates at 7 TeV

Case	$\sigma_{SS3\ell}$
	(fb)
$\lambda$ -type: $m_{\tilde{g}} \simeq 660\text{GeV}$ , neutralinoLSP	19.82
$\lambda$ -type: $m_{\widetilde{g}} \simeq 1  TeV$ , neutralinoLSP	4.29
$\lambda$ -type: $m_{\widetilde{g}} \simeq 770 GeV$ , stauLSP	30.74
$\lambda$ -type: $ extit{m}_{ ilde{ extit{g}}} \simeq 1 extit{TeV}$ , $ extit{stauLSP}$	3.35
$\lambda'$ -type: $m_{\tilde{g}} \simeq 660  GeV$ , neutralinoLSP	2.07

## In a purely phenomenological R-parity violating SUSY

### Some general conclusions about SS3 $\ell$ :

- $M_{2(1)} \ge 2M_{1(2)} \Rightarrow Signal \ rate \ enhanced$
- $M_1 \simeq M_2 \Rightarrow Signal suppressed unless M_3 >> M_{1,2}$
- For  $m_{\tilde{g}}$ ,  $m_{\tilde{q}} \simeq 1$  TeV, 5 background-free events possible with 0.6 5.5 fb<sup>-1</sup> at 7 TeV, and 0.1 3.0 fb<sup>-1</sup> at 14 TeV
- The nature of SUSY-RPV can be extracted for one type of RPV coupling present at a time

### Some event rates

BP	$M_1$	$M_2$	$M_{\chi_1^0}$	$M_{\chi_1^{\pm}}$	$M_{ ilde{e_L}}$	$M_{ ilde{ au}}$
	(GeV)	(GeV)	(GeV)	(GeV)	(GeV)	(GeV)
1	150	150	146.54	154.80	254.13	180.91
2	160	150	154.08	154.80	254.13	217.69
3	100	300	97.69	395.30	254.10	180.68
4	125	250	121.65	254.35	156.76	217.52

BP	$\sigma_{7 Tev} \ (\mathit{fb})$	$L_{7TeV} \ (fb^{-1})$	$\sigma_{14Tev} \ (fb)$	$\begin{array}{c} L_{14TeV} \\ (fb^{-1}) \end{array}$
1	0.91	5.49	4.60	1.09
2	0.41	12.20	1.62	3.09
3	2.81	1.78	20.67	0.24
4	8.78	0.57	42.93	0.12

 $L_{7TeV}, L_{7TeV}$ : luminosities (in  $fb^{-1}$ ) at thse energies required for five signal events

# Some numbers with $m_{\tilde{g}} \simeq 1$ TeV: ss3

Cut	SM	S	Sig(S)
Lepton selection + MET	$7.01 \times 10^{-4}$	2.41	5.9
MET > 50  GeV		2.23	5.6
MET > 100  GeV		1.65	4.7
$m_{eff}^{\ell} > 100~GeV$		2.39	5.8
$m_{eff}^{\ell} > 200~GeV$		1.57	4.6
$m_{eff} > 150 \ GeV$		2.40	5.9
$m_{eff} > 250 \ GeV$		2.10	5.4

Table: Results for the same-sign trilepton search in terms of the effective cross-sections (in fb) and the expected signal significance (in Gaussian  $\sigma$ s). The results are shown for  $7\,\text{TeV}$  LHC and  $1~\text{fb}^{-1}$  of integrated luminosity.

After lepton selection + MET cut, 3 events with zero "real" background, with  $\sim 1.25 fb^{-1}$ 

# Some numbers with $m_{\tilde{g}} \simeq 1$ TeV: SSD

Cut	SM	S	Sig(S)
Lepton selection $+ MET > 30 GeV$	10.7	11.61	3.1
MET > 50 GEV		10.95	
MET > 100  GeV		8.66	
$m_{ ext{ iny eff}}^{\ell} > 100~ ext{ iny GeV}$	6.4	10.59	3.5
$m_{ ext{eff}}^{ ilde{\ell}^{\prime}} > 200 \; GeV$	1.0	5.50	3.7
$m_{\it eff} > 150~\it GeV$	7.4	11.41	3.5
$m_{\rm eff} > 250~{\rm GeV}$	1.8	9.46	4.7

Table: Results for the same-sign dilepton search in terms of the effective cross-sections (in fb) and the expected signal significance (in Gaussian  $\sigma$ s). The results are shown for  $7\,\text{TeV}$  LHC and  $1\,\text{fb}^{-1}$  of integrated luminosity.

After lepton selection + MET cut,  $5\sigma$  discovery in the SSD channel with  $2.6fb^{-1}$ 

## A way of uniquely identifying the dynamics...

General conclusion: same-sign trileptons (SS3 $\ell$ ) with enhanced rate are smoking gun signals of SUSY with L-violation

S. Mukhopadhyay, BM (2010, 2011)

If the lightest neutralino is the LSP, its decays in pair can yield two leptons of identical sign

A third lepton of the same sign can come from the cascade

Similar happening with same-sign four-leptons (SS4 $\ell$ ), too

At the same time, one has mixed-sign tri-and four-lepton events (MS3 $\ell$ , MS4 $\ell$ )



# To identify the dynamics...

Define 
$$x = \sigma_{SS3\ell}/(\sigma_{SS3\ell} + \sigma_{MS3\ell})$$
,

and 
$$y = \sigma_{SS4\ell}/(\sigma_{SS4\ell} + \sigma_{MS4\ell})$$

The dynamics is reflected in x and y

For  $\lambda$ -type coupling,

$$x \simeq 0.12$$

In actual simulations,  $x = 0.11 \pm 0.2$ 



# To identify the dynamics...

With  $\lambda'$ -type coupling,

A  $\chi_1^0$ -LSP can decay into  $I^\pm q \bar q'$  or  $\nu q \bar q$ 

Let 
$$B(\chi_1^0 \longrightarrow l^{\pm}q\bar{q}') = \alpha$$
  
( $\alpha \simeq 0.5$ )

**Then** 

$$x = \alpha^2/4 + 4y(1/\alpha - 1)$$

# To identify the dynamics...

For the bilinear terms  $\epsilon_i L_i H_2$  side by side with  $\mu H_1 H_2$ ,

The  $\epsilon_i$  can be rotated away from the superpotential,

RPV is then driven by sneutrino vev in the scalar potential

Then correct neutrino masses  $\Rightarrow \langle \tilde{\nu} \rangle \simeq 100 \text{keV}$  in that basis

$$\chi_1^0 \longrightarrow \ell W, \nu Z$$

(BR's fixed unless sneutrinos are closely degenerate with the Higgs)

#### Then

$$x = 3.53y + 0.06$$

(Including backgrounds, the relations are satisfied upto 10 - 20 %)

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- Same-sign trileptons: a very clear signal of SUSY with L-violation
- The early run has interesting prospects
- SS3\ell and SS4\ell can differentiate among various R-parity breaking terms

"It is the mark of an educated mind to be able to be able to entertain a thought without accepting it"

—Aristotle