The Top-Antitop Threshold – QCD Contributions

Maximilian Stahlhofen

In collaboration with André Hoang

University of Vienna

Outline

- The resonance of $\sigma_{tot}(e^+e^- \rightarrow t\bar{t})$
 - Measurement
 - Theory
- The effective theory vNRQCD
- Renormalization
 - Currents
 - Potentials
- Recent Results
- Summary/Outlook



The resonance of $\sigma_{tot}(e^+e^- \rightarrow t \bar{t})$

EW effects:

• LO: Unstable top $\Rightarrow \Gamma_t \approx 1.5 \text{ GeV} \gg \Lambda_{\text{QCD}}$

$$v_{eff} \equiv \sqrt{\frac{\sqrt{s} - 2m_t}{m_t}} \rightarrow \sqrt{\frac{\sqrt{s} - 2m_t + i\Gamma_t}{m_t}} ;$$

"IR cutoff"

 $|v_{eff}|\gtrsim 0.1$

[Fadin, Khoze]

• Higher orders → Pedro's talk

<u>QCD in the resonance region:</u> $v \sim \alpha_s \ll 1$

3 scales:	$m_t \gg$	$ec{p} \sim m_t v \gg$	$E_{kin} \sim m_t v^2$	$(\sim \Gamma_t \gg \Lambda_{QCD})$	
	"hard"	"soft"	"ultrasoft"		
Problems:	 "Coulomb singularities" Large logs				

Problem of Coulomb singularities:



Solution:

Nonrelativistic effective field theory **vNRQCD**

 \rightarrow summation of $(\alpha_s/v)^n$ terms by means of a Schrödinger Equation !

Constructing vNRQCD:

"Expansion in v $\ll 1$ "

→ Integrate out nonresonant degrees of freedom, e.g.:



Green function:
$$\left[-\frac{\nabla_{\vec{r}}^2}{m} + V(\mathbf{r}) - E\right] G(\vec{r}, \vec{r}', E) = \delta^{(3)}(\vec{r} - \vec{r}')$$
$$G(0, 0, E) \sim \bigotimes + \bigotimes \bigvee + \cdots + \bigotimes \bigvee + \cdots$$

Optical theorem:

$$\sigma_{\rm tot} \sim \int d{\rm PS} \left| \ll + \ll + \ll + \cdots \right|^2 \sim {\rm Im} \left[{\sf G}(0,0,{\sf E}) \right]$$

$$\label{eq:unstabletop: G(0,0,E+i\Gamma_t)} \text{Unstable top: } G(0,0,E+i\Gamma_t) \sim \sum_n \, \frac{|\Psi_n(0)|^2}{E_n-E-i\Gamma_t-i\epsilon} \, + \, \text{continuum}$$

LO Green function:
$$\left[-\frac{\nabla_{\vec{r}}^2}{m} + V_c(r) - E\right] G^1(\vec{r}, \vec{r}', E) = \delta^{(3)}(\vec{r} - \vec{r}')$$

→ Analytic solution for
$$G^1(\vec{r}, \vec{r}', E)$$
 ✓

Higher orders: e.g.
$$G^2(0, 0, E) = -\int d^3 \vec{r} G^1(0, \vec{r}, E) V^2(\vec{r}) G^1(\vec{r}, 0, E)$$

$$\uparrow O(\alpha_s, v)$$

RG running: e.g.
$$\forall V \equiv V(\mu) \Rightarrow G(0, 0, E) \equiv G(0, 0, E, \mu)$$

→ $G(0, 0, E, \mu)$ known up to NNLL (partly N³LO) \checkmark [Hoang, Manohar, Stewart, Teubner; Pineda, Signer] [Beneke, Kiyo, Schuller]

Problem of large logarithms:

3	scale	es: m _t	$\gg \vec{p} \sim m_t$	$v \ \gg \ E_{kin} \sim m_t v^2$	$(\sim \Gamma_t \gg \Lambda_{QCD})$
		"hard"	′ "soft"	"ultrasoft"	
→	Logs:	$\ln(\frac{m^2}{E^2}), l$	$\ln\left(\frac{\mathbf{m}^2}{\mathbf{p}^2}\right), \ \ln\left(\frac{\mathbf{p}^2}{E^2}\right)$	e.g. $\alpha_{\rm s} \ln \left(\frac{{\rm m}^2}{{\rm E}^2} ight) \sim$	$- \alpha_{\rm s} \ln({\rm v}^4) \sim 1$

Solution:

Two renormalization scales: $\mu_s = m\nu, \ \mu_u = m\nu^2 \longrightarrow "v"NRQCD$

u "subtraction velocity"

→ RGE's resum
$$[\alpha_{s} \ln v]^{n}$$
, $\alpha_{s} [\alpha_{s} \ln v]^{n}$, $\alpha_{s}^{2} [\alpha_{s} \ln v]^{n}$... terms
LL NLL NNLL

Nonresonant dof's integrated out, e.g.:



- Resonant dof's \rightarrow fields in the vNRQCD Lagrangian:
- Systematic expansion in v \Rightarrow consistent power counting in v $\sim \alpha_{\rm s}$

[Luke, Manohar, Rothstein]

$$\mathcal{L}_{vNRQCD} \, = \, \mathcal{L}_{usoft} \, + \, \mathcal{L}_{pot} \, + \, \mathcal{L}_{soft}$$

$$\mathsf{D}^{\mu} = \partial^{\mu} + \mathsf{ig}\mathsf{A}^{\mu}(\mathsf{x})$$

0000

$$\mathcal{L}_{usoft}: \psi^{\dagger}_{\mathbf{p}}(\mathbf{x}) \left[\mathsf{i} \mathsf{D}^{0} - \frac{(\mathbf{p} - \mathsf{i} \mathbf{D})^{2}}{2\mathsf{m}} + \dots \right] \psi_{\mathbf{p}}(\mathbf{x}) + \dots$$









[Luke, Manohar, Rothstein]

$$\mathcal{L}_{vNRQCD} \,=\, \mathcal{L}_{usoft} \,+\, \mathcal{L}_{pot} \,+\, \mathcal{L}_{soft}$$

$$\mathsf{D}^{\mu}=\partial^{\mu}+\mathsf{ig}\mathsf{A}^{\mu}{}_{(\mathsf{x})}$$

$$\mathcal{L}_{usoft}$$
: $\psi^{\dagger}_{\mathbf{p}}(\mathbf{x}) \left[i \mathsf{D}^{0} - \frac{(\mathbf{p} - i\mathbf{D})^{2}}{2\mathsf{m}} + \dots \right] \psi_{\mathbf{p}}(\mathbf{x}) + \dots$

$$\mathcal{L}_{\text{pot}}: - \mathsf{V} \ \psi_{\mathbf{p}'}^{\dagger} \psi_{\mathbf{p}} \chi_{-\mathbf{p}'}^{\dagger} \chi_{-\mathbf{p}} + \dots$$
$$\mathsf{V} \sim \frac{\mathcal{V}_{c}}{\mathbf{k}^{2}} + \frac{\mathcal{V}_{k} \pi^{2}}{\mathbf{m} \mathbf{k}} + \frac{\mathcal{V}_{r} (\mathbf{p}^{2} + \mathbf{p}'^{2})}{2\mathbf{m}^{2} \mathbf{k}^{2}} + \frac{\mathcal{V}_{2}}{\mathbf{m}^{2}} \mathbf{S}^{2} + \dots$$



Production/annihilation current (³S₁):



$$\begin{split} \sigma_{\rm tot} &\sim {\rm Im} \bigg[\underbrace{\otimes} + \underbrace{\otimes} + \underbrace{\otimes} + \underbrace{\otimes} + \underbrace{\otimes} \underbrace{\vee} \underbrace{\vee} + \cdots \bigg] \\ &\sim |\mathsf{c}_1(\nu)|^2 \cdot {\rm Im} \left[-\mathsf{i} \int \mathsf{d}^4 \mathsf{x} \, \, \mathsf{e}^{\mathsf{i}\hat{\mathsf{q}}\mathsf{x}} \, \left\langle \mathsf{0} | \, \mathsf{T} \, \vec{\mathsf{j}}_1^{\, \mathsf{eff}} \, {}^*(\mathsf{x}) \, \vec{\mathsf{j}}_1^{\, \mathsf{eff}}(\mathsf{0}) \, \left| \mathsf{0} \right\rangle \right] + \dots \\ &\sim |\mathsf{c}_1(\nu)|^2 \cdot {\rm Im} \left[\mathsf{G}(\mathsf{0},\mathsf{0},\mathsf{E},\nu) \right] + \dots \quad \mathsf{G}^{\mathsf{NNLL}} \, \mathsf{known} \, \checkmark \end{split}$$





NLL:
$$\nu \frac{\partial}{\partial \nu} \ln[\mathbf{c}_1(\nu)] = -\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{16\pi^2} \left[\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{4} + \mathcal{V}_2(\nu) + \mathcal{V}_{\mathsf{r}}(\nu) + \mathbf{S}^2 \,\mathcal{V}_{\mathsf{s}}(\nu) \right] + \frac{1}{2} \mathcal{V}_{\mathsf{k}}(\nu)$$

NLL running of V:



NLL:
$$\nu \frac{\partial}{\partial \nu} \ln[\mathbf{c_1}(\nu)] = -\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{16\pi^2} \left[\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{4} + \mathcal{V}_{\mathsf{2}}(\nu) + \mathcal{V}_{\mathsf{r}}(\nu) + \mathbf{S}^2 \,\mathcal{V}_{\mathsf{s}}(\nu) \right] + \frac{1}{2} \mathcal{V}_{\mathsf{k}}(\nu)$$

NLL running of V:

• Ultrasoft contributions dominant

$$lpha_{
m s}({
m mv}^2)\simeq 0.27 > \left| lpha_{
m s}({
m mv})\simeq 0.15 \right| ({
m v}\simeq 0.1)$$

• Large ultrasoft contribution to ξ_{nonmix}^{NNLL} !!!



• Soft contributions to V_s known \rightarrow tiny



• Potentials affected by ultrasoft renormalization:

$$\frac{\frac{\mathcal{V}_{\mathsf{k}}\pi^{2}}{\mathsf{m}\mathbf{k}}}{\mathcal{O}(\alpha_{\mathsf{s}}^{2}\mathsf{v})}, \frac{\frac{\mathcal{V}_{\mathsf{r}}(\mathbf{p}^{2}+\mathbf{p'}^{2})}{2\mathsf{m}^{2}\mathbf{k}^{2}}, \frac{\mathcal{V}_{2}}{\mathsf{m}^{2}}}{\mathcal{O}(\alpha_{\mathsf{s}}\mathsf{v}^{2})}$$





NLL:
$$\nu \frac{\partial}{\partial \nu} \ln[\mathbf{c_1}(\nu)] = -\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{16\pi^2} \left[\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{4} + \mathcal{V}_{\mathsf{2}}(\nu) + \mathcal{V}_{\mathsf{r}}(\nu) + \mathbf{S}^2 \,\mathcal{V}_{\mathsf{s}}(\nu) \right] + \frac{1}{2} \mathcal{V}_{\mathsf{k}}(\nu)$$

One usoft loop:

Details:

- Feynman/Coulomb gauge
- $\overline{\mathrm{MS}}$, Dim. Reg.
- A_0 and A couple differently!
- Usoft derivative operator insertions:



$$\mathcal{L}_{usoft}: \quad \psi_{\mathbf{p}}^{\dagger} \Big[\underbrace{(i\partial_{0} - \frac{\mathbf{p}^{2}}{2m})}_{HQ \text{ propagator}} - gA_{0} + \frac{i\mathbf{p}\nabla}{m} + g\frac{\mathbf{p}A}{m} + \frac{\nabla^{2}}{2m} + \dots \Big] \psi_{\mathbf{p}}$$

NLL:
$$\nu \frac{\partial}{\partial \nu} \ln[\mathbf{c_1}(\nu)] = -\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{16\pi^2} \left[\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{4} + \mathcal{V}_{\mathsf{2}}(\nu) + \mathcal{V}_{\mathsf{r}}(\nu) + \mathbf{S}^2 \,\mathcal{V}_{\mathsf{s}}(\nu) \right] + \frac{1}{2} \mathcal{V}_{\mathsf{k}}(\nu)$$

One usoft loop:

Details:

- Feynman/Coulomb gauge
- $\overline{\mathrm{MS}}$, Dim. Reg.
- A_0 and A couple differently!
- Usoft derivative operator insertions:

$$\mathcal{L}_{usoft}: \quad \psi_{\mathbf{p}}^{\dagger} \Big[\underbrace{(i\partial_{0} - \frac{\mathbf{p}^{2}}{2m})}_{HQ \text{ propagator}} - gA_{0} + \frac{i\mathbf{p}\nabla}{m} + g\frac{\mathbf{p}A}{m} + \frac{\nabla^{2}}{2m} + \dots \Big] \psi_{\mathbf{p}}$$

 A_0

NLL:
$$\nu \frac{\partial}{\partial \nu} \ln[\mathbf{c_1}(\nu)] = -\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{16\pi^2} \left[\frac{\mathcal{V}_{\mathsf{c}}(\nu)}{4} + \mathcal{V}_{\mathsf{2}}(\nu) + \mathcal{V}_{\mathsf{r}}(\nu) + \mathbf{S}^2 \,\mathcal{V}_{\mathsf{s}}(\nu) \right] + \frac{1}{2} \mathcal{V}_{\mathsf{k}}(\nu)$$

One usoft loop:

Details:

- Feynman/Coulomb gauge •
- $\overline{\mathrm{MS}}$, Dim. Reg. •
- A_0 and A couple differently! •
- Usoft derivative operator insertions: •

A couple differently!
erivative operator insertions:
$$\mathcal{L}_{usoft}: \quad \psi_{\mathbf{p}}^{\dagger} \Big[\underbrace{(i\partial_0 - \frac{\mathbf{p}^2}{2m})}_{HQ \text{ propagator}} - gA_0 + \frac{i\mathbf{p}\nabla}{m} + g\frac{\mathbf{p}A}{m} + \frac{\nabla^2}{2m} + \dots \Big] \psi_{\mathbf{p}}$$

 A_0





RGE's + matching at hard scale ($\nu = 1$) give: LL NLL $\left[\mathcal{V}_{2}(\nu)\right]_{\text{usoft}}^{\text{NLL}} = 4\pi\alpha_{s}(\mathsf{m}\nu)\left[-\frac{4\pi}{\beta_{0}}\mathsf{A}_{2}\ln\frac{\alpha_{s}(\mathsf{m}\nu^{2})}{\alpha_{s}(\mathsf{m}\nu)} + \left(\frac{\beta_{1}}{\beta_{2}^{2}}\mathsf{A}_{2} - \left[\alpha_{s}(\mathsf{m}\nu^{2}) - \alpha_{s}(\mathsf{m}\nu)\right]\frac{8\pi}{\beta_{0}}\mathsf{B}_{2}\right)\right]$ $\left[\mathcal{V}_{\mathsf{r}}(\nu) \right]_{\mathsf{usoft}}^{\mathsf{NLL}} = 8\pi\alpha_{\mathsf{s}}(\mathsf{m}\nu) \left| -\frac{4\pi}{\beta_0}\mathsf{A}_{\mathsf{r}}\ln\frac{\alpha_{\mathsf{s}}(\mathsf{m}\nu^2)}{\alpha_{\mathsf{s}}(\mathsf{m}\nu)} + \left(\frac{\beta_1}{\beta_0^2}\mathsf{A}_{\mathsf{r}} - \left[\alpha_{\mathsf{s}}(\mathsf{m}\nu^2) - \alpha_{\mathsf{s}}(\mathsf{m}\nu) \right] \frac{8\pi}{\beta_0}\mathsf{B}_{\mathsf{r}} \right) \right|$ $\left[\mathcal{V}_{\mathsf{k}}(\nu)\right]_{\mathsf{usoft}}^{\mathsf{NLL}} = 2\alpha_{\mathsf{s}}^{2}(\mathsf{m}\nu)\left[-\frac{4\pi}{\beta_{0}}\mathsf{A}_{\mathsf{k}}\ln\frac{\alpha_{\mathsf{s}}(\mathsf{m}\nu^{2})}{\alpha_{\mathsf{s}}(\mathsf{m}\nu)} + \left(\frac{\beta_{1}}{\beta_{0}^{2}}\mathsf{A}_{\mathsf{k}} - \left[\alpha_{\mathsf{s}}(\mathsf{m}\nu^{2}) - \alpha_{\mathsf{s}}(\mathsf{m}\nu)\right]\frac{8\pi}{\beta_{0}}\mathsf{B}_{\mathsf{k}}\right)\right]$ $\begin{vmatrix} A_2 \\ B_2 \end{vmatrix} = C_F(C_A - 2C_F) \begin{vmatrix} A \\ B \end{vmatrix}$ $\mathsf{A} = \frac{\mathsf{I}}{3\pi}$ $\begin{vmatrix} A_{r} \\ B_{r} \end{vmatrix} = -C_{A}C_{F} \begin{bmatrix} A \\ B \end{vmatrix}$ $\mathsf{B} = \frac{\mathsf{C}_{\mathsf{A}}(47 + 6\pi^2) - 10\mathsf{n}_{\mathsf{f}}\mathsf{T}}{100}$ [MS, Hoang] $\begin{bmatrix} A_{k} \\ B_{k} \end{bmatrix} = -C_{A}C_{F}(C_{A} - 2C_{F})\begin{bmatrix} A \\ B \end{bmatrix}$ [Pineda]









Detailed analysis → WIP

Summary/Outlook

- $y_t, \alpha_s, \Gamma_t \text{ from } \sigma_{tot}(e^+e^- \rightarrow t \overline{t}) \text{ at threshold}$
- status: $\frac{\delta \sigma_{\rm tot}^{\rm th}}{\sigma_{\rm tot}} \approx 6\%$, needed: $\frac{\delta \sigma_{\rm tot}^{\rm th}}{\sigma_{\rm tot}} \lesssim 3\%$
- $\sigma_{\text{tot}} \sim |\mathbf{c}_1(\nu)|^2 \cdot \text{Im} \left[\mathsf{G}(\mathbf{0}, \mathbf{0}, \mathsf{E}, \nu) \right] + \dots$
- $\mathsf{G}(\mathbf{0},\mathbf{0},\mathsf{E},\nu)\,$ known up to NNLL \checkmark
- New NNLL_{mix,usoft} compensates for large NNLL_{nonmix} contribution to $c_1(\nu)$

$$\rightarrow \frac{\delta \sigma_{\rm tot}^{\rm th}}{\sigma_{\rm tot}}$$
 decreases!

- Outlook:
 - Detailed study of $\sigma_{tot}(e^+e^- \rightarrow t \, \overline{t} \,)$ at threshold
 - Determination of bottom mass from nonrel. Υ sum rules

Backup

vNRQCD label formalism:

momentum space:

